# LIMITING CASE HARDY INEQUALITIES ON THE SPHERE 

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Abstract. We give sharp limiting case Hardy inequalities on the sphere $\mathbb{S}^{2}$ and show that their optimal constants are unattainable by any $f \in H^{1}\left(\mathbb{S}^{2}\right) \backslash\{0\}$. The singularity of the problem is related to the geodesic distance from a point on the sphere.

## 1. Introduction

The classical Hardy inequality

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}|\nabla u|^{2} d x \geqslant \frac{(n-2)^{2}}{4} \int_{\mathbb{R}^{n}} \frac{u^{2}}{|x|^{2}} d x \tag{1}
\end{equation*}
$$

is valid in dimensions $n \geqslant 3$ for all functions $u \in H^{1}\left(\mathbb{R}^{n}\right)([1])$. It obviously fails on $\mathbb{R}^{2}$ as the right hand side of (1) no longer makes sense. In order to obtain a version of (1) in the critical case $n=2$ on bounded domains, a logarithmic weight can be introduced to tame the singularity. In $[2,4-8,10]$, for instance, inequalities of the type

$$
\int_{B}|\nabla u|^{n} d x \geqslant C_{n}(B) \int_{B} \frac{|u|^{n}}{|x|^{n}\left(\log \frac{1}{|x|}\right)^{n}} d x
$$

were analysed for $u \in W_{0}^{1, n}(B)$ where $B$ is the unit ball in $\mathbb{R}^{n}$.
Let $n \geqslant 3$ and $\mathbb{S}^{n}$ be the unit sphere equipped with its Lebesgue surface measure $\sigma_{n}$ in $\mathbb{R}^{n+1}$. Denote by $d(., p): \mathbb{S}^{n} \rightarrow[0, \pi]$ the geodesic distance from $p \in \mathbb{S}^{n}$, and by $\nabla_{\mathbb{S}^{n}}$ the gradient on $\mathbb{S}^{n}$. Recently, Xiao [11] proved that if $f \in C^{\infty}\left(\mathbb{S}^{n}\right)$ then

$$
\begin{equation*}
\bar{c}_{n} \int_{\mathbb{S}^{n}} f^{2} d \sigma_{n}+\int_{\mathbb{S}^{n}}\left|\nabla_{\mathbb{S}^{n}} f\right|^{2} d \sigma_{n} \geqslant c_{n}^{2} \int_{\mathbb{S}^{n}}\left(\frac{f^{2}}{d(x, p)^{2}}+\frac{f^{2}}{(\pi-d(x, p))^{2}}\right) d \sigma_{n} \tag{2}
\end{equation*}
$$

with $\bar{c}_{n}=\left(\frac{4}{3}+\frac{1}{\pi^{2}}\right) c_{n}^{2}+c_{n}, c_{n}=\frac{n-2}{2}$. It was also shown in [11] that the constant $c_{n}$ in (2) is sharp in the sense that

$$
c_{n}^{2}=\inf _{f \in C^{\infty}\left(\mathbb{S}^{n}\right) \backslash\{0\}} \frac{D_{n}(f)}{\int_{\mathbb{S}^{n}} \frac{f^{2}}{d(x, p)^{2}} d \sigma_{n}}=\inf _{f \in C^{\infty}\left(\mathbb{S}^{n}\right) \backslash\{0\}} \frac{D_{n}(f)}{\int_{\mathbb{S}^{n}} \frac{f^{2}}{(\pi-d(x, p))^{2}} d \sigma_{n}}
$$

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where

$$
D_{n}(f):=\bar{c}_{n} \int_{\mathbb{S}^{n}} f^{2} d \sigma_{n}+\int_{\mathbb{S}^{n}}\left|\nabla_{\mathbb{S}^{n}} f\right|^{2} d \sigma_{n}, \quad f \in C^{\infty}\left(\mathbb{S}^{n}\right)
$$

Recently, Xiao's result was extended to the case $p \neq 2$ in [9].
We prove $L^{2}$ Hardy inequalities with optimal constants on the sphere $\mathbb{S}^{2}$ in $\mathbb{R}^{3}$. This is a critical exponent case as the integral $\int_{\mathbb{S}^{2}} \theta^{-1+\lambda} d \sigma_{2}$, where $\theta$ is the polar angle, diverges for $\lambda \leqslant-1$. We also argue the lack of maximizers for our inequalities. Our approach denies the possibility of an equality in Xiao's inequality (2) as well.

## 2. Preliminaries

A point on the sphere $\mathbb{S}^{2}$ will have the standard spherical coordinate parametrization $(\sin \theta \cos \varphi, \sin \theta \sin \varphi, \cos \theta)$ where $\theta \in[0, \pi]$ refers to the polar angle and $\varphi \in$ $[0,2 \pi[$ is the azimuthal angle. Then the surface measure induced by the Lebesgue measure on $\mathbb{R}^{3}$ is $d \sigma_{2}=\sin \theta d \theta d \varphi$, the gradient and the Laplace-Beltrami operator, respectively, are given by

$$
\nabla_{\mathbb{S}^{2}}=\hat{\theta} \frac{\partial}{\partial \theta}+\hat{\varphi} \frac{1}{\sin \theta} \frac{\partial}{\partial \varphi}, \quad \Delta_{\mathbb{S}^{2}}=\frac{1}{\sin \theta} \frac{\partial}{\partial \theta}\left(\sin \theta \frac{\partial}{\partial \theta}\right)+\frac{1}{\sin ^{2} \theta} \frac{\partial^{2}}{\partial \varphi^{2}}
$$

Here $\hat{\theta}, \hat{\varphi}$ denote the orthogonal three-dimensional unit vectors in the direction where $\theta, \varphi$ increase, respectively. The Sobolev space $H^{1}\left(\mathbb{S}^{2}\right)$ is the completion of $C^{\infty}\left(\mathbb{S}^{2}\right)$ in the norm

$$
\|f\|_{H^{1}\left(\mathbb{S}^{2}\right)}:=\left(\|f\|_{L^{2}\left(\mathbb{S}^{2}\right)}^{2}+\|\nabla f\|_{L^{2}\left(\mathbb{S}^{2}\right)}^{2}\right)^{\frac{1}{2}}
$$

In order to find the geodesic distance $d(x, p)$ from a point $x \in \mathbb{S}^{2}$ to a given a point $p \in \mathbb{S}^{2}$, we rotate the axes, if necessary, to put $p$ on the zenith direction then place the great circle passing through $p$ and $x$ in the azimuth reference direction so that we have $d(x, p)=\theta$.

For simplicity, we henceforth denote $d \sigma_{2}, \nabla_{\mathbb{S}^{2}}$ and $\Delta_{\mathbb{S}^{2}}$ by $d \sigma, \nabla$ and $\Delta$, respectively.

## 3. Main results

Let $\phi:] 0, \pi] \rightarrow[1, \infty[$ be defined by $\phi(t):=\log (\pi e / t), \psi:] 0, \pi[\rightarrow[1+\log \pi, \infty[$ be such that $\psi(t):=\phi(\sin t)$, and $\rho_{\phi}(t):=t \phi(t)$. Let $A>0$. Denote by $S, T_{A}$, and $Q(. ; \phi)$ the positive nonlinear functionals on $H^{1}\left(\mathbb{S}^{2}\right)$ given by

$$
\begin{aligned}
S(f) & :=\int_{\mathbb{S}^{2}}|\hat{\theta} \cdot \nabla f|^{2} d \sigma+\frac{1}{2 \pi^{2}} \int_{\mathbb{S}^{2}} f^{2} d \sigma \\
T_{A}(f) & :=\int_{\mathbb{S}^{2}}|\nabla f|^{2} d \sigma_{2}+\frac{A}{4} \int_{\mathbb{S}^{2}} f^{2} d \sigma_{2}, \\
Q(f ; \phi) & :=\frac{1}{4} \int_{\mathbb{S}^{2}}\left(\frac{f^{2}}{\rho_{\phi}^{2}(d(x, p))}+\frac{f^{2}}{\rho_{\phi}^{2}(\pi-d(x, p))}\right) d \sigma_{2} .
\end{aligned}
$$

THEOREM 1. Assume that $f \in H^{1}\left(\mathbb{S}^{2}\right)$. Then there exists constants $A, B>0$, independent of $f$, such that

$$
\begin{align*}
& Q(f ; \phi) \leqslant T_{A}(f)  \tag{3}\\
& Q(f ; \psi) \leqslant T_{B}(f) \tag{4}
\end{align*}
$$

Both inequalities (3) and (4) are optimal, but an equality is impossible in either one:

THEOREM 2.

$$
\begin{array}{r}
\sup _{f \in H^{1}\left(\mathbb{S}^{2}\right) \backslash\{0\}} \frac{Q(f ; \phi)}{T_{A}(f)}=1, \\
\sup _{f \in H^{1}\left(\mathbb{S}^{2}\right) \backslash\{0\}} \frac{Q(f ; \psi)}{T_{B}(f)}=1 . \tag{6}
\end{array}
$$

THEOREM 3. There does not exist $f \in H^{1}\left(\mathbb{S}^{2}\right) \backslash\{0\}$ such that $Q(f ; \phi)=T_{A}(f)$, or $Q(f ; \psi)=T_{B}(f)$.

A variant of the above-mentioned results follows via a different approach:
Theorem 4. Let $f \in H^{1}\left(\mathbb{S}^{2}\right)$. Then

$$
\begin{align*}
& \frac{1}{4} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\rho_{\phi}^{2}(d(x, p))} d \sigma \leqslant S(f)+\frac{1}{2 \pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\pi-d(x, p)} d \sigma  \tag{7}\\
& \frac{1}{4} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\rho_{\phi}^{2}(\pi-d(x, p))} d \sigma \leqslant S(f)+\frac{1}{2 \pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{d(x, p)} d \sigma \tag{8}
\end{align*}
$$

Moreover

$$
\begin{equation*}
\sup _{f \in H^{1}\left(\mathbb{S}^{2}\right) \backslash\{0\}} \frac{\frac{1}{4} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\rho_{\phi^{2}}(d(x, p))} d \sigma}{S(f)+\frac{1}{2 \pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\pi-d(x, p)} d \sigma}=\sup _{f \in H^{1}\left(\mathbb{S}^{2}\right) \backslash\{0\}} \frac{\frac{1}{4} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\rho_{\phi}^{2}(\pi-d(x, p))} d \sigma}{S(f)+\frac{1}{2 \pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{d(x, p)} d \sigma}=1 \tag{9}
\end{equation*}
$$

and the suprema in (9) are not attained in $H^{1}\left(\mathbb{S}^{2}\right) \backslash\{0\}$.

## 4. Proof of Theorem 1

Proof. Let $f \in C^{\infty}\left(\mathbb{S}^{2}\right)$. Notice that $\psi>1$ and write $f(\theta, \varphi)=\sqrt{\psi(\theta)} g(\theta, \varphi)$. We have

$$
\begin{align*}
|\nabla f|^{2} & =\left|\psi^{\frac{1}{2}} \nabla g+g \nabla \psi^{\frac{1}{2}}\right|^{2} \\
& =\psi|\nabla g|^{2}+\left\langle\psi^{\frac{1}{2}} \nabla g, g \psi^{-\frac{1}{2}} \nabla \psi\right\rangle+\left|\frac{1}{2} \psi^{-\frac{1}{2}} \nabla \psi\right|^{2} g^{2} \\
& =\psi|\nabla g|^{2}+\frac{1}{2}\left\langle\nabla \psi, \nabla g^{2}\right\rangle+\frac{1}{4} \frac{1}{\psi}|\nabla \psi|^{2} g^{2} . \tag{10}
\end{align*}
$$

Integrating both sides of (10) over $\mathbb{S}^{2}$ we get

$$
\begin{align*}
\int_{\mathbb{S}^{2}}|\nabla f|^{2} d \sigma & =\int_{\mathbb{S}^{2}}\left(\psi|\nabla g|^{2}+\frac{1}{2}\left\langle\nabla \psi, \nabla g^{2}\right\rangle+\frac{1}{4} \frac{1}{\psi}|\nabla \psi|^{2} g^{2}\right) d \sigma \\
& \geqslant \frac{1}{4} \int_{\mathbb{S}^{2}} \frac{1}{\psi}|\nabla \psi|^{2} g^{2} d \sigma+\frac{1}{2} \int_{\mathbb{S}^{2}}\left\langle\nabla \psi, \nabla g^{2}\right\rangle d \sigma  \tag{11}\\
& =\frac{1}{4} \int_{\mathbb{S}^{2}} \frac{1}{\psi}\left|\psi^{\prime}\right|^{2} g^{2} d \sigma-\frac{1}{2} \int_{\mathbb{S}^{2}} g^{2} \Delta \psi d \sigma \tag{12}
\end{align*}
$$

by partial integration over the closed manifold $\mathbb{S}^{2}$. Calculating, we find

$$
\begin{equation*}
\Delta \psi=\frac{1}{\sin \theta} \frac{\partial}{\partial \theta}\left(\sin \theta \frac{\partial}{\partial \theta} \psi\right)=1 \tag{13}
\end{equation*}
$$

Returning $g$ to $f / \sqrt{\psi}$ and substituting for $\Delta \psi$ from (13) into (12), we obtain

$$
\begin{equation*}
\int_{\mathbb{S}^{2}}|\nabla f|^{2} d \sigma \geqslant \frac{1}{4} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\psi^{2}} \frac{\cos ^{2} \theta}{\sin ^{2} \theta} d \sigma-\frac{1}{2} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\psi} d \sigma \tag{14}
\end{equation*}
$$

Adding the finite integral $\frac{1}{4} \int_{\mathbb{S}^{2}}\left(\frac{1}{\theta^{2} \phi^{2}(\theta)}+\frac{1}{(\pi-\theta)^{2} \phi^{2}(\pi-\theta)}\right) f^{2} d \sigma$ to both sides of (14) transforms it into the inequality

$$
\begin{align*}
\frac{1}{4} \int_{\mathbb{S}^{2}}\left(\frac{1}{\theta^{2} \phi^{2}(\theta)}\right. & \left.+\frac{1}{(\pi-\theta)^{2} \phi^{2}(\pi-\theta)}\right) f^{2} d \sigma \\
& \leqslant \int_{\mathbb{S}^{2}}|\nabla f|^{2} d \sigma+\frac{1}{4} \int_{\mathbb{S}^{2}} F(\theta) f^{2} d \sigma \tag{15}
\end{align*}
$$

where

$$
F(t):=\frac{1}{t^{2} \phi^{2}(t)}+\frac{1}{(\pi-t)^{2} \phi^{2}(\pi-t)}-\frac{\cos ^{2} t}{\sin ^{2} t} \frac{1}{\phi^{2}(\sin t)}+\frac{2}{\phi(\sin t)}
$$

Obviously, $F$ is continuous on $] 0, \pi[$ and, as expected from the facts that $\phi(t) \rightarrow+\infty$ when $t \rightarrow 0^{+}, \sin t=t+o(t)$ as $t \rightarrow 0$, it turns out

$$
\lim _{t \rightarrow 0^{+}} F(t)=\lim _{t \rightarrow \pi^{-}} F(t)=\frac{1}{\pi^{2}}
$$

Hence, $F$ can be extended to a uniformly continuous, consequently a bounded, function on $[0, \pi]$. Noting this in (15) implies (3). Direct computation also shows

$$
A=\sup _{[0, \pi]}|F|=F\left(\frac{\pi}{2}\right)=\frac{2}{1+\log \pi}+\frac{8}{(1+\log 2)^{2}} \frac{1}{\pi^{2}}
$$

To prove (4), we add to both sides of (14) the well-defined integral

$$
\frac{1}{4} \int_{\mathbb{S}^{2}}\left(\frac{1}{\theta^{2}}+\frac{1}{(\pi-\theta)^{2}}\right) \frac{f^{2}}{\psi^{2}(\theta)} d \sigma
$$

We then obtain the following analogue of (15):

$$
\begin{align*}
\frac{1}{4} \int_{\mathbb{S}^{2}}\left(\frac{1}{\theta^{2}}+\right. & \left.\frac{1}{(\pi-\theta)^{2}}\right) \frac{f^{2}}{\psi^{2}(\theta)} d \sigma \\
& \leqslant \int_{\mathbb{S}^{2}}|\nabla f|^{2} d \sigma+\frac{1}{4} \int_{\mathbb{S}^{2}} G(\theta) f^{2} d \sigma \tag{16}
\end{align*}
$$

where

$$
\begin{align*}
G(t) & :=\frac{M(t)}{\psi^{2}(t)}+\frac{2}{\psi(t)} \\
M(t) & :=\frac{1}{t^{2}}+\frac{1}{(\pi-t)^{2}}-\frac{\cos ^{2} t}{\sin ^{2} t} \tag{17}
\end{align*}
$$

Once the boundedness of $G$ is ensured, we see that (16) yields the inequality (4). Evidently, $G$ has the same features as $F$. Since

$$
\begin{equation*}
\lim _{\theta \rightarrow 0} M(\theta)=\lim _{\theta \rightarrow \pi} M(\theta)=\frac{2}{3}+\frac{1}{\pi^{2}}, \quad \lim _{\theta \rightarrow 0^{+}} \psi(t)=\lim _{\theta \rightarrow \pi^{-}} \psi(t)=+\infty \tag{18}
\end{equation*}
$$

then $M \in C[0, \pi]$, and $\lim _{t \rightarrow 0^{+}} G(t)=\lim _{t \rightarrow \pi^{-}} G(t)=0$, which makes $G$ bounded on $[0, \pi]$. Moreover

$$
B=\sup _{[0, \pi]}|G|=G\left(\frac{\pi}{2}\right)=\frac{2}{1+\log \pi}+\frac{8}{(1+\log \pi)^{2}} \frac{1}{\pi^{2}}
$$

## 5. Proof of Theorem 2

Proof. First, we would like to define the weak Laplace-Beltrami gradient of a function $f \in L^{1}\left(\mathbb{S}^{2}\right)$. Suppose $f \in C^{\infty}\left(\mathbb{S}^{2}\right)$ and $v(\theta, \varphi)=v_{\theta}(\theta, \varphi) \hat{\theta}+v_{\varphi}(\theta, \varphi) \hat{\varphi}$ with $v_{\theta}, v_{\varphi} \in C^{\infty}\left(\mathbb{S}^{2}\right)$. Then

$$
\begin{aligned}
& \int_{\mathbb{S}^{2}} \frac{\partial f}{\partial \theta} v_{\theta} d \sigma=\int_{\mathbb{S}^{2}} \nabla f \cdot \hat{\theta} v_{\theta} d \sigma=-\int_{\mathbb{S}^{2}} f \nabla \cdot\left(v_{\theta} \hat{\theta}\right) d \sigma \\
& \int_{\mathbb{S}^{2}} \frac{1}{\sin \theta} \frac{\partial f}{\partial \varphi} v_{\varphi} d \sigma=\int_{\mathbb{S}^{2}} \nabla f \cdot \hat{\varphi} v_{\varphi} d \sigma=-\int_{\mathbb{S}^{2}} f \nabla \cdot\left(v_{\varphi} \hat{\varphi}\right) d \sigma .
\end{aligned}
$$

Adding these identities we get

$$
\begin{equation*}
\int_{\mathbb{S}^{2}} \nabla f \cdot V d \sigma=-\int_{\mathbb{S}^{2}} f \nabla \cdot V d \sigma \tag{19}
\end{equation*}
$$

for any vector field $V \in C^{\infty}\left(\mathbb{S}^{2} \rightarrow T\left(\mathbb{S}^{2}\right)\right)$ where $T\left(\mathbb{S}^{2}\right)$ is the tangent bundle of the smooth manifold $\mathbb{S}^{2}$. Motivated by (19), $f$ is weakly differentiable if there is a vector field $\vartheta_{f} \in L^{1}\left(\mathbb{S}^{2} \rightarrow T\left(\mathbb{S}^{2}\right)\right)$ such that

$$
\begin{equation*}
\int_{\mathbb{S}^{2}} \vartheta_{f} \cdot V d \sigma=-\int_{\mathbb{S}^{2}} f \nabla \cdot V d \sigma, \quad \forall V \in C^{\infty}\left(\mathbb{S}^{2} \rightarrow T\left(\mathbb{S}^{2}\right)\right) \tag{20}
\end{equation*}
$$

This, unique up to a set of zero measure, vector field $\vartheta_{f}$ is the weak surface gradient of $f$. According to ([3], Proposition 3.2., page 15)

$$
H^{1}\left(\mathbb{S}^{2}\right)=W^{1,2}\left(\mathbb{S}^{2}\right):=\left\{f \in L^{2}\left(\mathbb{S}^{2}\right):\left|\vartheta_{f}\right| \in L^{2}\left(\mathbb{S}^{2}\right)\right\}
$$

We start with (5). By Theorem 1, it suffices to prove the existence of a sequence $\left\{f_{n}\right\}_{n \geqslant 1}$ in $H^{1}\left(\mathbb{S}^{2}\right)$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{Q\left(f_{n} ; \phi\right)}{T_{A}\left(f_{n}\right)}=1 \tag{21}
\end{equation*}
$$

Consider the functions

$$
\begin{equation*}
f_{n}(\theta, \varphi):=\phi(\theta)^{\frac{1}{2}-\frac{1}{n}} \tag{22}
\end{equation*}
$$

The functions $f_{n}$ are independent of $\varphi$, hence

$$
\begin{equation*}
\frac{Q\left(f_{n} ; \phi\right)}{T_{A}\left(f_{n}\right)}=\frac{\int_{0}^{\pi} \frac{f_{n}^{2} \sin \theta}{\theta^{2} \phi^{2}(\theta)} d \theta+\int_{0}^{\pi} \frac{f_{n}^{2} \sin \theta}{(\pi-\theta)^{2} \phi^{2}(\pi-\theta)} d \theta}{4 \int_{0}^{\pi}\left(\frac{\partial f_{n}}{\partial \theta}\right)^{2} \sin \theta d \theta+A \int_{0}^{\pi} f_{n}^{2} \sin \theta d \theta} \tag{23}
\end{equation*}
$$

where the derivative $\partial f_{n} / \partial \theta$ is understood in the week sense discussed above. Since $\phi \in L_{\text {loc }}^{1}(\mathbb{R})$ and $\phi \geqslant 1$ on $[0, \pi]$, then

$$
\begin{equation*}
\int_{0}^{\pi} f_{n}^{2} \sin \theta d \theta=\int_{0}^{\pi} \phi(\theta)^{1-\frac{2}{n}} \sin \theta d \theta \leqslant \int_{0}^{\pi} \phi(\theta) d \theta \approx 1 \tag{24}
\end{equation*}
$$

Thus $f_{n} \in L^{2}\left(\mathbb{S}^{2}\right)$ for all $n \geqslant 1$. Notice also that $f_{n}$ is smooth on $[0, \pi] \backslash\{0\}$ and its weak derivative

$$
\begin{equation*}
\frac{\partial f_{n}}{\partial \theta}=\frac{\frac{1}{n}-\frac{1}{2}}{\theta \phi^{\frac{1}{2}+\frac{1}{n}}} \tag{25}
\end{equation*}
$$

Therefore

$$
\int_{0}^{\pi}\left(\frac{\partial f_{n}}{\partial \theta}\right)^{2} \sin \theta d \theta=\frac{a_{n}}{4} \int_{0}^{\pi} \frac{1}{\theta \phi^{1+\frac{2}{n}}} \frac{\sin \theta}{\theta} d \theta, \quad a_{n}:=\left(1-\frac{2}{n}\right)^{2}
$$

And since $\int_{0}^{\pi} \frac{d \theta}{\theta \phi^{1+\frac{2}{n}}}=\frac{n}{2}, \sin \theta \leqslant \theta$, then $\partial f_{n} / \partial \theta \in L^{2}\left(\mathbb{S}^{2}\right)$ for all $n \geqslant 1$. Substituting for $f_{n}$ from (22) and for $\partial f_{n} / \partial \theta$ from (25) into (23) implies

$$
\begin{equation*}
\frac{Q\left(f_{n} ; \phi\right)}{T_{A}\left(f_{n}\right)}=\frac{\alpha_{n}+\beta_{n}}{a_{n} \alpha_{n}+\gamma_{n}}=\frac{1}{a_{n}}\left(1+\frac{\beta_{n}-\gamma_{n} / a_{n}}{\alpha_{n}+\gamma_{n} / a_{n}}\right) \tag{26}
\end{equation*}
$$

where

$$
\begin{aligned}
\alpha_{n} & :=\int_{0}^{\pi} \frac{1}{\theta \phi^{1+\frac{2}{n}}} \frac{\sin \theta}{\theta} d \theta \\
\beta_{n} & :=\int_{0}^{\pi} \frac{\phi^{1-\frac{2}{n}}(\theta) \sin \theta}{(\pi-\theta)^{2} \phi^{2}(\pi-\theta)} d \theta \\
\gamma_{n} & :=A \int_{0}^{\pi} \phi^{1-\frac{2}{n}} \sin \theta d \theta
\end{aligned}
$$

Observe that $\lim _{n \rightarrow+\infty} a_{n}=1$. We shall show that, while $\lim _{n \rightarrow+\infty} \alpha_{n}=+\infty$, the sequences $\left\{\beta_{n}\right\}_{n \geqslant 1}$ and $\left\{\gamma_{n}\right\}_{n \geqslant 1}$ are both convergent. Using this in (26) proves (21).

Exploiting the continuity and positivity of $\sin \theta /\left(\theta^{2} \phi^{1+\frac{2}{n}}\right)$ on $[\pi / 2, \pi]$, then applying the inequality $\sin \theta / \theta \geqslant 2 / \pi$ when $0 \leqslant \theta \leqslant \pi / 2$, we obtain

$$
\begin{align*}
\alpha_{n} & =\int_{0}^{\pi / 2} \frac{1}{\theta \phi^{1+\frac{2}{n}}} \frac{\sin \theta}{\theta} d \theta+\int_{\pi / 2}^{\pi} \frac{\sin \theta}{\theta^{2} \phi^{1+\frac{2}{n}}} d \theta \\
& \geqslant \frac{2}{\pi} \int_{0}^{\pi / 2} \frac{1}{\theta \phi^{1+\frac{2}{n}}} d \theta=\frac{n}{\pi(1+\log (2))^{\frac{2}{n}}} \tag{27}
\end{align*}
$$

This proves the divergence of $\left\{\alpha_{n}\right\}$. Next, by the dominated convergence theorem and (24) we readily find

$$
\lim _{n \rightarrow+\infty} \gamma_{n}=A \lim _{n \rightarrow+\infty} \int_{0}^{\pi} \phi^{1-\frac{2}{n}}(\theta) \sin \theta d \theta=\int_{0}^{\pi} \phi(\theta) \sin \theta d \theta \lesssim 1
$$

Finally, since $\theta \mapsto \sin \theta /\left((\pi-\theta)^{2} \phi^{2}(\pi-\theta)\right) \in C([0, \pi / 2])$, then using the local integrability of $\phi$ and the dominated convergence theorem again implies

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{0}^{\pi / 2} \frac{\phi^{1-\frac{1}{n}}(\theta) \sin \theta}{(\pi-\theta)^{2} \phi^{2}(\pi-\theta)} d \theta=\int_{0}^{\pi / 2} \frac{\phi(\theta) \sin \theta}{(\pi-\theta)^{2} \phi^{2}(\pi-\theta)} d \theta \lesssim 1 \tag{28}
\end{equation*}
$$

Furthermore, since $\phi \in C([\pi / 2, \pi])$, and $\frac{\sin \theta}{\pi-\theta}=\frac{\sin (\pi-\theta)}{\pi-\theta} \leqslant 1$, on $[\pi / 2, \pi]$, then

$$
\begin{equation*}
\int_{\pi / 2}^{\pi} \frac{\phi^{1-\frac{1}{n}}(\theta) \sin \theta}{(\pi-\theta)^{2} \phi^{2}(\pi-\theta)} d \theta \lesssim \int_{\pi / 2}^{\pi} \frac{d \theta}{(\pi-\theta) \phi^{2}(\pi-\theta)} \approx 1 \tag{29}
\end{equation*}
$$

The convergence of $\left\{\beta_{n}\right\}$ follows from (28) together with (29).
The proof of (6) shares the main idea of (5). The functions $g_{n}(\theta, \varphi):=\psi(\theta)^{\frac{1}{2}-\frac{1}{n}} \in$ $L^{2}\left(\mathbb{S}^{2}\right), n \geqslant 1$, and satisfy $\lim _{n \rightarrow \infty} \frac{Q\left(g_{n} ; \psi\right)}{T_{B}\left(g_{n}\right)}=1$. Indeed, we have

$$
\begin{aligned}
\frac{Q\left(g_{n} ; \psi\right)}{T_{B}\left(g_{n}\right)} & =\frac{\int_{0}^{\pi} \frac{g_{n}^{2} \sin \theta}{\theta^{2} \psi^{2}(\theta)} d \theta+\int_{0}^{\pi} \frac{g_{n}^{2} \sin \theta}{(\pi-\theta)^{2} \psi^{2}(\pi-\theta)} d \theta}{4 \int_{0}^{\pi}\left(\frac{\partial g_{n}}{\partial \theta}\right)^{2} \sin \theta d \theta+B \int_{0}^{\pi} g_{n}^{2} \sin \theta d \theta} \\
& =\frac{\tilde{\alpha}_{n}}{a_{n} \tilde{\alpha}_{n}+\tilde{\beta}_{n}}=\frac{1}{a_{n}}\left(1-\frac{\tilde{\beta}_{n} / a_{n}}{\tilde{\alpha}_{n}+\tilde{\beta}_{n} / a_{n}}\right)
\end{aligned}
$$

where

$$
\begin{aligned}
& \tilde{\alpha}_{n}:=\int_{0}^{\pi} \frac{\sin \theta d \theta}{\theta^{2} \psi^{1+\frac{2}{n}}}+\int_{0}^{\pi} \frac{\sin \theta d \theta}{(\pi-\theta)^{2} \psi^{1+\frac{2}{n}}}=2 \int_{0}^{\pi} \frac{\sin \theta d \theta}{\theta^{2} \psi^{1+\frac{2}{n}}} \\
& \tilde{\beta}_{n}:=B \int_{0}^{\pi} \psi^{1-\frac{2}{n}} \sin \theta d \theta-a_{n} \int_{0}^{\pi} M(\theta) \frac{\sin \theta}{\psi^{1+\frac{2}{n}}} d \theta
\end{aligned}
$$

Similarly to (27), we have

$$
\begin{aligned}
\tilde{\alpha}_{n} & =2 \int_{0}^{1} \frac{\sin \theta}{\theta^{2}} \frac{1}{\psi^{1+\frac{2}{n}}} d \theta+2 \int_{1}^{\pi} \frac{\sin \theta}{\theta^{2}} \frac{1}{\psi^{1+\frac{2}{n}}} d \theta \\
& \geqslant 2 \int_{0}^{1} \frac{\sin \theta}{\theta^{2}} \frac{1}{\psi^{1+\frac{2}{n}}} d \theta=2 \int_{0}^{1} \frac{\sin ^{2} \theta}{\theta^{2} \cos \theta} \frac{1}{\psi^{1+\frac{2}{n}}} \frac{\cos \theta}{\sin \theta} d \theta \\
& \geqslant \frac{8}{\pi^{2}} \int_{0}^{1} \frac{1}{\psi^{1+\frac{2}{n}}} \frac{\cos \theta}{\sin \theta} d \theta=\frac{4 n}{\pi^{2}} \frac{1}{(1+\log \pi)^{\frac{2}{n}}} .
\end{aligned}
$$

Hence $\lim _{n \rightarrow \infty} \tilde{\alpha}_{n}=\infty$. Recall from (17) and (18) that $M \in C([0, \pi])$. Also, since $\psi \in L_{\mathrm{loc}}^{1}(\mathbb{R}), \psi>1$ uniformly, then $\lim _{n \rightarrow \infty} \tilde{\beta}_{n}$ exists by the dominated convergence theorem.

## 6. Proof of Theorem 3

Proof. The transition to the inequalities (3) and (4) from their respective stronger versions, (15) and (16), comes from the bounds

$$
\int_{\mathbb{S}^{2}} F(\theta) f^{2} d \sigma \leqslant A \int_{\mathbb{S}^{2}} f^{2} d \sigma, \quad \int_{\mathbb{S}^{2}} G(\theta) f^{2} d \sigma \leqslant B \int_{\mathbb{S}^{2}} f^{2} d \sigma
$$

where the bounded functions $F$ and $G$ are both positive and independent of $f$. Interestingly, as seen in Section 5, the size of $0<A, B<\infty$ played no role in optimising (3) and (4).

Up to the inequality (15) or (16) an equality relation persists except for the only inequality (11). So a sufficient and necessary condition for an equality in (15) or (16) (and a necessary condition for an equality in (3) and (4)) is an equality in (11). But an equality in (11) occurs if and only if

$$
\begin{equation*}
\int_{\mathbb{S}^{2}} \psi|\nabla g|^{2} d \sigma=0 \tag{30}
\end{equation*}
$$

Recalling that $g=f / \sqrt{\psi}$, we compute

$$
\begin{align*}
\psi|\nabla g|^{2} & =\psi\left|\frac{\nabla f}{\sqrt{\psi}}-\frac{1}{2} \frac{f}{\psi^{\frac{3}{2}}} \frac{\partial \psi}{\partial \theta} \hat{\theta}\right|^{2} \\
& =|\nabla f|^{2}-\frac{f}{\psi} \frac{\partial \psi}{\partial \theta} \nabla f \cdot \hat{\theta}+\frac{1}{4} \frac{f^{2}}{\psi^{2}}\left(\frac{\partial \psi}{\partial \theta}\right)^{2} \\
& =|\nabla f|^{2}-\left(\frac{\partial f}{\partial \theta}\right)^{2}+\left(\frac{\partial f}{\partial \theta}\right)^{2}-\frac{f}{\psi} \frac{\partial \psi}{\partial \theta} \frac{\partial f}{\partial \theta}+\frac{1}{4} \frac{f^{2}}{\psi^{2}}\left(\frac{\partial \psi}{\partial \theta}\right)^{2} \\
& =|\nabla f|^{2}-\left(\frac{\partial f}{\partial \theta}\right)^{2}+\left(\frac{\partial f}{\partial \theta}-\frac{1}{2} \frac{f}{\psi} \frac{\partial \psi}{\partial \theta}\right)^{2} \tag{31}
\end{align*}
$$

Since $|\nabla f|^{2}-\left(\frac{\partial f}{\partial \theta}\right)^{2}=\frac{1}{\sin ^{2} \theta}\left(\frac{\partial f}{\partial \varphi}\right)^{2} \geqslant 0$, then, by (31), the equality (30) is equivalent to

$$
\begin{equation*}
\int_{\mathbb{S}^{2}}|\nabla f|^{2}-\left(\frac{\partial f}{\partial \theta}\right)^{2} d \sigma=\int_{\mathbb{S}^{2}}\left(\frac{\partial f}{\partial \theta}-\frac{1}{2} \frac{f}{\psi} \frac{\partial \psi}{\partial \theta}\right)^{2} d \sigma=0 \tag{32}
\end{equation*}
$$

The equalities (32) are, in their turn, equivalent to

$$
\begin{equation*}
\frac{1}{\sin \theta}\left|\frac{\partial f}{\partial \varphi}\right|=\left|\frac{\partial f}{\partial \theta}-\frac{1}{2} \frac{f}{\psi} \frac{\partial \psi}{\partial \theta}\right|=0 \tag{33}
\end{equation*}
$$

Suppose that $f$ is not the zero function. Then (33) are possible if and only if

$$
f=f(\theta), \quad \frac{d f}{f}=\frac{1}{2} \frac{d \psi}{\psi} .
$$

That is $f=c \sqrt{\psi}, c$ is a constant. But such $f \notin H^{1}\left(\mathbb{S}^{2}\right)$ because

$$
\begin{aligned}
\int_{\mathbb{S}^{2}}|\nabla f|^{2} d \sigma & =2 \pi \int_{0}^{\pi}\left(\frac{\partial f}{\partial \theta}\right)^{2} \sin \theta d \theta \gtrsim \int_{0}^{1} \frac{\cos ^{2} \theta}{\sin \theta} \frac{1}{\psi} d \theta \\
& \gtrsim \int_{0}^{1} \frac{d \theta}{\sin \theta \phi(\sin \theta)} \approx \int_{0}^{1} \frac{d \theta}{\theta \phi(\theta)}=+\infty
\end{aligned}
$$

## 7. Proof of Theorem 4

Proof. Write

$$
\frac{1}{\theta} \frac{1}{\phi^{2}(\theta)}=\nabla\left(\frac{1}{\phi(\theta)}\right) \cdot \hat{\theta}
$$

Assume that $f$ is smooth. Then integrating by parts w.r.t. the surface measure $\sigma$ we get

$$
\begin{align*}
\int_{\mathbb{S}^{2}} \frac{f^{2}}{\theta^{2} \phi^{2}(\theta)} d \sigma & =\int_{\mathbb{S}^{2}} \nabla\left(\frac{1}{\phi(\theta)}\right) \cdot \frac{f^{2}}{\theta} \hat{\theta} d \sigma \\
& =-\int_{\mathbb{S}^{2}} \frac{1}{\phi(\theta)} \nabla \cdot\left(\frac{f^{2}}{\theta} \hat{\theta}\right) d \sigma \\
& =-2 \int_{\mathbb{S}^{2}} \frac{f \nabla f \cdot \hat{\theta}}{\theta \phi(\theta)} d \sigma+\int_{\mathbb{S}^{2}} \frac{f^{2}}{\theta^{2} \phi(\theta)} d \sigma-\int_{\mathbb{S}^{2}} \frac{f^{2}}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta} d \sigma \tag{34}
\end{align*}
$$

Observe here that each of the last two integrals on the right hand side of (34) can diverge. They suffer nonintegrable singularities at $\theta=0$. But, when put together, their sum

$$
\begin{equation*}
I:=\int_{\mathbb{S}^{2}} \frac{f^{2}}{\theta^{2} \phi(\theta)} d \sigma-\int_{\mathbb{S}^{2}} \frac{f^{2}}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta} d \sigma=\int_{\mathbb{S}^{2}} \frac{1}{\theta \phi(\theta)}\left(\frac{1}{\theta}-\frac{\cos \theta}{\sin \theta}\right) f^{2} d \sigma \tag{35}
\end{equation*}
$$

is convergent. In fact

$$
\lim _{\theta \rightarrow 0^{+}} \frac{1}{\theta \phi(\theta)}\left(\frac{1}{\theta}-\frac{\cos \theta}{\sin \theta}\right)=0
$$

Also, $\theta \mapsto 1 /\left(\theta^{2} \phi(\theta)\right)$ is continuous on a neighborhood of $\theta=\pi$. Furthermore, if we fix $\delta>0$ and let $D:=\left\{x(\theta, \varphi) \in \mathbb{S}^{2}: 0 \leqslant \theta<\delta\right\}$, then the integral

$$
\int_{\mathbb{S}^{2} \backslash D} \frac{f^{2}}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta} d \sigma
$$

does exist. Unfortunately, we can not control the integral $I$ by $\int_{\mathbb{S}^{2}} f^{2} d \sigma$, up to a constant factor. The reason is

$$
\lim _{\theta \rightarrow \pi^{-}} \frac{1}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta}=-\infty
$$

But since

$$
\lim _{\theta \rightarrow \pi^{-}}\left(\frac{1}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta}+\frac{1}{\pi} \frac{1}{(\pi-\theta)}\right)=0
$$

then, we may introduce the convergent integral $J:=\frac{1}{\pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\pi-\theta} d \sigma$ to the integral $I$ to get

$$
\begin{equation*}
I=I-J+J=\int_{\mathbb{S}^{2}} K(\theta) f^{2} d \sigma+J \tag{36}
\end{equation*}
$$

where

$$
K(\theta):=\frac{1}{\theta \phi(\theta)}\left(\frac{1}{\theta}-\frac{\cos \theta}{\sin \theta}\right)-\frac{1}{\pi} \frac{1}{(\pi-\theta)}
$$

By the continuity of $K$ on $] 0, \pi[$ and since

$$
\lim _{\theta \rightarrow 0^{+}} K(\theta)=-\lim _{\theta \rightarrow \pi^{-}} K(\theta)=-\frac{1}{\pi^{2}}
$$

then $K$ is bounded on $[0, \pi]$. Actually, $K$ is monotonically increasing. Thus

$$
\begin{equation*}
\sup _{[0, \pi]}|K|=\frac{1}{\pi^{2}} \tag{37}
\end{equation*}
$$

Using (37) in (36) we deduce that

$$
\begin{equation*}
I \leqslant \frac{1}{\pi^{2}} \int_{\mathbb{S}^{2}} f^{2} d \sigma+J \tag{38}
\end{equation*}
$$

Returning with (38) to the inequality (34) in the light of (35) we obtain

$$
\begin{equation*}
\int_{\mathbb{S}^{2}} \frac{f^{2}}{\theta^{2} \phi^{2}(\theta)} d \sigma \leqslant-2 \int_{\mathbb{S}^{2}} \frac{f \nabla f . \hat{\theta}}{\theta \phi(\theta)} d \sigma+\frac{1}{\pi^{2}} \int_{\mathbb{S}^{2}} f^{2} d \sigma+\frac{1}{\pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\pi-\theta} d \sigma \tag{39}
\end{equation*}
$$

Applying Cauchy's inequality with an $\varepsilon$ we find

$$
\begin{equation*}
-2 \int_{\mathbb{S}^{2}} \frac{f \nabla f . \hat{\theta}}{\theta \phi(\theta)} d \sigma \leqslant 2 \varepsilon \int_{\mathbb{S}^{2}} \frac{f^{2}}{\theta^{2} \phi^{2}(\theta)} d \sigma+\frac{1}{2 \varepsilon} \int_{\mathbb{S}^{2}}|\hat{\theta} . \nabla f|^{2} d \sigma \tag{40}
\end{equation*}
$$

Therefore, it follows from (39) and (40) that

$$
\begin{align*}
2 \varepsilon(1-2 \varepsilon) \int_{\mathbb{S}^{2}} \frac{f^{2}}{\theta^{2} \phi^{2}(\theta)} d \sigma \leqslant & \int_{\mathbb{S}^{2}}|\hat{\theta} \cdot \nabla f|^{2} d \sigma+\frac{2 \varepsilon}{\pi^{2}} \int_{\mathbb{S}^{2}} f^{2} d \sigma \\
& +\frac{2 \varepsilon}{\pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\pi-\theta} d \sigma, \quad 0<\varepsilon<\frac{1}{2} \tag{41}
\end{align*}
$$

The choice $\varepsilon=1 / 4$ maximizes the factor $2 \varepsilon(1-2 \varepsilon)$ and, consequently, the left hand side of (41). This proves (7). The inequality (8) can be obtained analogously.

In the fashion of the proof of Theorem 2, the sequence $f_{n}=\phi^{\frac{1}{2}-\frac{1}{n}}$ clearly satisfies

$$
\lim _{n \rightarrow \infty} \frac{\frac{1}{4} \int_{0}^{\pi} \frac{f_{n}^{2}}{\rho_{\phi}^{2}(\theta)} \sin \theta d \theta}{U\left(f_{n}\right)+\frac{1}{2 \pi} \int_{0}^{\pi} \frac{f_{n}^{2}}{\pi-\theta} \sin \theta d \theta}=\lim _{n \rightarrow \infty} \frac{\frac{1}{4} \int_{0}^{\pi} \frac{f_{n}^{2}}{\rho_{\phi}^{2}(\pi-\theta)} \sin \theta d \theta}{U\left(f_{n}\right)+\frac{1}{2 \pi} \int_{0}^{\pi} \frac{f_{n}^{2}}{\theta} \sin \theta d \theta}=1
$$

where

$$
U(f)=\int_{0}^{\pi}\left(\frac{\partial f}{\partial \theta}\right)^{2} \sin \theta d \theta+\frac{1}{2 \pi^{2}} \int_{0}^{\pi} f^{2} \sin \theta d \theta
$$

One only needs to inspect the convergence of $\int_{0}^{\pi}\left(\phi^{1-\frac{2}{n}} \sin \theta / \theta\right) d \theta, \int_{0}^{\pi}\left(\phi^{1-\frac{2}{n}} \sin \theta /\right.$ $(\pi-\theta)) d \theta$ as $n \rightarrow \infty$. This is obvious from the bound $\sin \theta \leqslant \min \{\theta, \pi-\theta\}$ on $[0, \pi]$ and the fact $\phi \in L^{1}([0, \pi])$.

Finally, careful review of the proof of (7) above reveals that a necessary condition for a function $f \in H^{1}\left(\mathbb{S}^{2}\right) \backslash\{0\}$ to achieve an equality in (7) is that it yields an equality in (40). This is equivalent to

$$
\begin{equation*}
\nabla f . \hat{\theta}=-\frac{1}{2} \frac{f}{\theta \phi(\theta)} \tag{42}
\end{equation*}
$$

Suppose (42) was true. Then by (34) and (35) we must have

$$
\begin{equation*}
\int_{\mathbb{S}^{2}} \frac{h(\theta) f^{2}}{\theta \phi(\theta)} d \sigma=0 \tag{43}
\end{equation*}
$$

where

$$
h(\theta):=\frac{1}{\theta}-\frac{\cos \theta}{\sin \theta} .
$$

On the other hand

$$
\lim _{\theta \rightarrow 0^{+}} h(\theta)=0, \quad h^{\prime}(\theta)=\frac{\theta^{2}-\sin ^{2} \theta}{\theta^{2} \sin ^{2} \theta}>0, \quad 0<\theta<\pi
$$

This shows $h$ is strictly positive on $] 0, \pi]$ and since $\theta \phi(\theta) \geqslant 0$ then (43) is a contradiction.

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## REFERENCES

[1] A. Balinsky, W. D. Evans, R. T. Lewis, The analysis and geometry of Hardy's inequality, Springer, New York 2015.
[2] J. BYEON, F. TAKAHASHI, Hardy's inequality in a limiting case on general bounded domains, arXiv:1707.04018, 2017.
[3] JÜRgEn Eichhorn, Global analysis on open manifolds, Nova Science Publishers, Inc., New York, 2007.
[4] N. Ioku, M. Ishiwata, A scale invariant form of a critical Hardy inequality, International Mathematics Research Notices 18 (2015), 8830-8846.
[5] N. Ioku, M. Ishiwata and T. Ozawa, Sharp remainder of a critical Hardy inequality, Archiv der Mathematik 106 (2016), 65-71.
[6] S. Machihara, T. Ozawa and H. Wadade, Hardy type inequalities on balls, Tohoku Mathematical Journal 65 (2013), No. 3, 321-330.
[7] M. RuZhansky and D. Suragan, Critical Hardy inequalities, arXiv:1602.04809, 2016.
[8] M. SanO, F. TAKAhashi, Scale invariance structures of the critical and the subcritical Hardy inequalities and their improvements, Calculus of variations and partial differential equations, (2017) 56: 69.
[9] X. Sun, F. Pan, Hardy type inequalities on the sphere, J. Inequalities and Applications, (2017) 2017: 148.
[10] F. TAKAHASHI, A simple proof of Hardy's inequality in a limiting case, Archiv der Mathematik 104 (2015), 1, 77-82.
[11] Y. Xiao, Some Hardy inequalities on the sphere, J. Math. Inequal. 10 (2016), 793-805.

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