LIMITING CASE HARDY INEQUALITIES ON THE SPHERE

AHMED A. ABDELHAKIM

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Abstract. We give sharp limiting case Hardy inequalities on the sphere \mathbb{S}^2 and show that their optimal constants are unattainable by any $f \in H^1(\mathbb{S}^2) \setminus \{0\}$. The singularity of the problem is related to the geodesic distance from a point on the sphere.

1. Introduction

The classical Hardy inequality

$$\int_{\mathbb{R}^n} |\nabla u|^2 dx \ge \frac{(n-2)^2}{4} \int_{\mathbb{R}^n} \frac{u^2}{|x|^2} dx \tag{1}$$

is valid in dimensions $n \ge 3$ for all functions $u \in H^1(\mathbb{R}^n)$ ([1]). It obviously fails on \mathbb{R}^2 as the right hand side of (1) no longer makes sense. In order to obtain a version of (1) in the critical case n = 2 on bounded domains, a logarithmic weight can be introduced to tame the singularity. In [2, 4–8, 10], for instance, inequalities of the type

$$\int_{B} |\nabla u|^{n} dx \ge C_{n}(B) \int_{B} \frac{|u|^{n}}{|x|^{n} \left(\log \frac{1}{|x|}\right)^{n}} dx$$

were analysed for $u \in W_0^{1,n}(B)$ where B is the unit ball in \mathbb{R}^n .

Let $n \ge 3$ and \mathbb{S}^n be the unit sphere equipped with its Lebesgue surface measure σ_n in \mathbb{R}^{n+1} . Denote by $d(.,p): \mathbb{S}^n \to [0,\pi]$ the geodesic distance from $p \in \mathbb{S}^n$, and by $\nabla_{\mathbb{S}^n}$ the gradient on \mathbb{S}^n . Recently, Xiao [11] proved that if $f \in C^{\infty}(\mathbb{S}^n)$ then

$$\overline{c}_n \int_{\mathbb{S}^n} f^2 d\sigma_n + \int_{\mathbb{S}^n} |\nabla_{\mathbb{S}^n} f|^2 d\sigma_n \ge c_n^2 \int_{\mathbb{S}^n} \left(\frac{f^2}{d(x,p)^2} + \frac{f^2}{(\pi - d(x,p))^2} \right) d\sigma_n \quad (2)$$

with $\overline{c}_n = \left(\frac{4}{3} + \frac{1}{\pi^2}\right)c_n^2 + c_n$, $c_n = \frac{n-2}{2}$. It was also shown in [11] that the constant c_n in (2) is sharp in the sense that

$$c_n^2 = \inf_{f \in C^{\infty}(\mathbb{S}^n) \setminus \{0\}} \frac{D_n(f)}{\int_{\mathbb{S}^n} \frac{f^2}{d(x,p)^2} d\sigma_n} = \inf_{f \in C^{\infty}(\mathbb{S}^n) \setminus \{0\}} \frac{D_n(f)}{\int_{\mathbb{S}^n} \frac{f^2}{(\pi - d(x,p))^2} d\sigma_n}$$

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where

$$D_n(f) := \overline{c}_n \int_{\mathbb{S}^n} f^2 d\sigma_n + \int_{\mathbb{S}^n} |\nabla_{\mathbb{S}^n} f|^2 d\sigma_n, \quad f \in C^{\infty}(\mathbb{S}^n).$$

Recently, Xiao's result was extended to the case $p \neq 2$ in [9].

We prove L^2 Hardy inequalities with optimal constants on the sphere \mathbb{S}^2 in \mathbb{R}^3 . This is a critical exponent case as the integral $\int_{\mathbb{S}^2} \theta^{-1+\lambda} d\sigma_2$, where θ is the polar angle, diverges for $\lambda \leq -1$. We also argue the lack of maximizers for our inequalities. Our approach denies the possibility of an equality in Xiao's inequality (2) as well.

2. Preliminaries

A point on the sphere \mathbb{S}^2 will have the standard spherical coordinate parametrization $(\sin\theta\cos\varphi,\sin\theta\sin\varphi,\cos\theta)$ where $\theta \in [0,\pi]$ refers to the polar angle and $\varphi \in [0,2\pi[$ is the azimuthal angle. Then the surface measure induced by the Lebesgue measure on \mathbb{R}^3 is $d\sigma_2 = \sin\theta d\theta d\varphi$, the gradient and the Laplace-Beltrami operator, respectively, are given by

$$\nabla_{\mathbb{S}^2} = \hat{\theta} \frac{\partial}{\partial \theta} + \hat{\varphi} \frac{1}{\sin \theta} \frac{\partial}{\partial \varphi}, \quad \Delta_{\mathbb{S}^2} = \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \varphi^2}$$

Here $\hat{\theta}, \hat{\varphi}$ denote the orthogonal three-dimensional unit vectors in the direction where θ, φ increase, respectively. The Sobolev space $H^1(\mathbb{S}^2)$ is the completion of $C^{\infty}(\mathbb{S}^2)$ in the norm

$$\| f \|_{H^1(\mathbb{S}^2)} := \left(\| f \|_{L^2(\mathbb{S}^2)}^2 + \| \nabla f \|_{L^2(\mathbb{S}^2)}^2 \right)^{\frac{1}{2}}.$$

In order to find the geodesic distance d(x, p) from a point $x \in \mathbb{S}^2$ to a given a point $p \in \mathbb{S}^2$, we rotate the axes, if necessary, to put p on the zenith direction then place the great circle passing through p and x in the azimuth reference direction so that we have $d(x, p) = \theta$.

For simplicity, we henceforth denote $d\sigma_2$, $\nabla_{\mathbb{S}^2}$ and $\Delta_{\mathbb{S}^2}$ by $d\sigma$, ∇ and Δ , respectively.

3. Main results

Let $\phi :]0, \pi] \to [1, \infty]$ be defined by $\phi(t) := \log(\pi e/t), \quad \psi :]0, \pi[\to [1 + \log \pi, \infty]$ be such that $\psi(t) := \phi(\sin t)$, and $\rho_{\phi}(t) := t\phi(t)$. Let A > 0. Denote by S, T_A , and $Q(.;\phi)$ the positive nonlinear functionals on $H^1(\mathbb{S}^2)$ given by

$$\begin{split} S(f) &:= \int_{\mathbb{S}^2} |\hat{\theta} . \nabla f|^2 d\sigma + \frac{1}{2\pi^2} \int_{\mathbb{S}^2} f^2 d\sigma, \\ T_A(f) &:= \int_{\mathbb{S}^2} |\nabla f|^2 d\sigma_2 + \frac{A}{4} \int_{\mathbb{S}^2} f^2 d\sigma_2, \\ Q(f;\phi) &:= \frac{1}{4} \int_{\mathbb{S}^2} \left(\frac{f^2}{\rho_{\phi}{}^2 (d(x,p))} + \frac{f^2}{\rho_{\phi}{}^2 (\pi - d(x,p))} \right) d\sigma_2. \end{split}$$

THEOREM 1. Assume that $f \in H^1(\mathbb{S}^2)$. Then there exists constants A, B > 0, independent of f, such that

$$Q(f;\phi) \leqslant T_A(f), \tag{3}$$

$$Q(f; \psi) \leqslant T_B(f). \tag{4}$$

Both inequalities (3) and (4) are optimal, but an equality is impossible in either one:

THEOREM 2.

$$\sup_{f \in H^1(\mathbb{S}^2) \setminus \{0\}} \frac{Q(f;\phi)}{T_A(f)} = 1,$$
(5)

$$\sup_{f \in H^1(\mathbb{S}^2) \setminus \{0\}} \frac{\mathcal{Q}(f; \psi)}{T_B(f)} = 1.$$
(6)

THEOREM 3. There does not exist $f \in H^1(\mathbb{S}^2) \setminus \{0\}$ such that $Q(f;\phi) = T_A(f)$, or $Q(f;\psi) = T_B(f)$.

A variant of the above-mentioned results follows via a different approach:

THEOREM 4. Let $f \in H^1(\mathbb{S}^2)$. Then

$$\frac{1}{4} \int_{\mathbb{S}^2} \frac{f^2}{\rho_{\phi}^2(d(x,p))} d\sigma \leqslant S(f) + \frac{1}{2\pi} \int_{\mathbb{S}^2} \frac{f^2}{\pi - d(x,p)} d\sigma, \tag{7}$$

$$\frac{1}{4} \int_{\mathbb{S}^2} \frac{f^2}{\rho_{\phi}^2 \left(\pi - d\left(x, p\right)\right)} d\sigma \leqslant S(f) + \frac{1}{2\pi} \int_{\mathbb{S}^2} \frac{f^2}{d\left(x, p\right)} d\sigma.$$
(8)

Moreover

$$\sup_{f \in H^{1}(\mathbb{S}^{2}) \setminus \{0\}} \frac{\frac{1}{4} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\rho_{\phi}^{2}(d(x,p))} d\sigma}{S(f) + \frac{1}{2\pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\pi - d(x,p)} d\sigma} = \sup_{f \in H^{1}(\mathbb{S}^{2}) \setminus \{0\}} \frac{\frac{1}{4} \int_{\mathbb{S}^{2}} \frac{f^{2}}{\rho_{\phi}^{2}(\pi - d(x,p))} d\sigma}{S(f) + \frac{1}{2\pi} \int_{\mathbb{S}^{2}} \frac{f^{2}}{d(x,p)} d\sigma} = 1,$$
(9)

and the suprema in (9) are not attained in $H^1(\mathbb{S}^2) \setminus \{0\}$.

4. Proof of Theorem 1

Proof. Let $f \in C^{\infty}(\mathbb{S}^2)$. Notice that $\psi > 1$ and write $f(\theta, \varphi) = \sqrt{\psi(\theta)}g(\theta, \varphi)$. We have

$$\begin{aligned} |\nabla f|^2 &= |\psi^{\frac{1}{2}} \nabla g + g \nabla \psi^{\frac{1}{2}}|^2 \\ &= \psi |\nabla g|^2 + \langle \psi^{\frac{1}{2}} \nabla g, g \psi^{-\frac{1}{2}} \nabla \psi \rangle + \left| \frac{1}{2} \psi^{-\frac{1}{2}} \nabla \psi \right|^2 g^2 \\ &= \psi |\nabla g|^2 + \frac{1}{2} \langle \nabla \psi, \nabla g^2 \rangle + \frac{1}{4} \frac{1}{\psi} |\nabla \psi|^2 g^2. \end{aligned}$$
(10)

Integrating both sides of (10) over \mathbb{S}^2 we get

$$\int_{\mathbb{S}^2} |\nabla f|^2 d\sigma = \int_{\mathbb{S}^2} \left(\psi |\nabla g|^2 + \frac{1}{2} \langle \nabla \psi, \nabla g^2 \rangle + \frac{1}{4} \frac{1}{\psi} |\nabla \psi|^2 g^2 \right) d\sigma$$
$$\geqslant \frac{1}{4} \int_{\mathbb{S}^2} \frac{1}{\psi} |\nabla \psi|^2 g^2 d\sigma + \frac{1}{2} \int_{\mathbb{S}^2} \langle \nabla \psi, \nabla g^2 \rangle d\sigma \tag{11}$$

$$= \frac{1}{4} \int_{\mathbb{S}^2} \frac{1}{\psi} |\psi'|^2 g^2 d\sigma - \frac{1}{2} \int_{\mathbb{S}^2} g^2 \Delta \psi d\sigma$$
(12)

by partial integration over the closed manifold \mathbb{S}^2 . Calculating, we find

$$\Delta \psi = \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \psi \right) = 1.$$
(13)

Returning g to $f/\sqrt{\psi}$ and substituting for $\Delta \psi$ from (13) into (12), we obtain

$$\int_{\mathbb{S}^2} |\nabla f|^2 d\sigma \ge \frac{1}{4} \int_{\mathbb{S}^2} \frac{f^2}{\psi^2} \frac{\cos^2 \theta}{\sin^2 \theta} d\sigma - \frac{1}{2} \int_{\mathbb{S}^2} \frac{f^2}{\psi} d\sigma.$$
(14)

Adding the finite integral $\frac{1}{4} \int_{\mathbb{S}^2} \left(\frac{1}{\theta^2 \phi^2(\theta)} + \frac{1}{(\pi - \theta)^2 \phi^2(\pi - \theta)} \right) f^2 d\sigma$ to both sides of (14) transforms it into the inequality

$$\frac{1}{4} \int_{\mathbb{S}^2} \left(\frac{1}{\theta^2 \phi^2(\theta)} + \frac{1}{(\pi - \theta)^2 \phi^2(\pi - \theta)} \right) f^2 d\sigma$$
$$\leqslant \int_{\mathbb{S}^2} |\nabla f|^2 d\sigma + \frac{1}{4} \int_{\mathbb{S}^2} F(\theta) f^2 d\sigma, \tag{15}$$

where

$$F(t) := \frac{1}{t^2 \phi^2(t)} + \frac{1}{(\pi - t)^2 \phi^2(\pi - t)} - \frac{\cos^2 t}{\sin^2 t} \frac{1}{\phi^2(\sin t)} + \frac{2}{\phi(\sin t)}.$$

Obviously, *F* is continuous on $]0,\pi[$ and, as expected from the facts that $\phi(t) \to +\infty$ when $t \to 0^+$, $\sin t = t + o(t)$ as $t \to 0$, it turns out

$$\lim_{t \to 0^+} F(t) = \lim_{t \to \pi^-} F(t) = \frac{1}{\pi^2}$$

Hence, *F* can be extended to a uniformly continuous, consequently a bounded, function on $[0, \pi]$. Noting this in (15) implies (3). Direct computation also shows

$$A = \sup_{[0,\pi]} |F| = F\left(\frac{\pi}{2}\right) = \frac{2}{1 + \log \pi} + \frac{8}{\left(1 + \log 2\right)^2} \frac{1}{\pi^2}.$$

To prove (4), we add to both sides of (14) the well-defined integral

$$\frac{1}{4}\int_{\mathbb{S}^2}\left(\frac{1}{\theta^2}+\frac{1}{(\pi-\theta)^2}\right)\frac{f^2}{\psi^2(\theta)}d\sigma.$$

We then obtain the following analogue of (15):

$$\frac{1}{4} \int_{\mathbb{S}^2} \left(\frac{1}{\theta^2} + \frac{1}{(\pi - \theta)^2} \right) \frac{f^2}{\psi^2(\theta)} d\sigma$$
$$\leqslant \int_{\mathbb{S}^2} |\nabla f|^2 d\sigma + \frac{1}{4} \int_{\mathbb{S}^2} G(\theta) f^2 d\sigma, \tag{16}$$

where

$$G(t) := \frac{M(t)}{\psi^2(t)} + \frac{2}{\psi(t)},$$

$$M(t) := \frac{1}{t^2} + \frac{1}{(\pi - t)^2} - \frac{\cos^2 t}{\sin^2 t}.$$
(17)

Once the boundedness of G is ensured, we see that (16) yields the inequality (4). Evidently, G has the same features as F. Since

$$\lim_{\theta \to 0} M(\theta) = \lim_{\theta \to \pi} M(\theta) = \frac{2}{3} + \frac{1}{\pi^2}, \quad \lim_{\theta \to 0^+} \psi(t) = \lim_{\theta \to \pi^-} \psi(t) = +\infty$$
(18)

then $M \in C[0, \pi]$, and $\lim_{t\to 0^+} G(t) = \lim_{t\to \pi^-} G(t) = 0$, which makes G bounded on $[0, \pi]$. Moreover

$$B = \sup_{[0,\pi]} |G| = G\left(\frac{\pi}{2}\right) = \frac{2}{1 + \log \pi} + \frac{8}{\left(1 + \log \pi\right)^2} \frac{1}{\pi^2}. \quad \Box$$

5. Proof of Theorem 2

Proof. First, we would like to define the weak Laplace-Beltrami gradient of a function $f \in L^1(\mathbb{S}^2)$. Suppose $f \in C^{\infty}(\mathbb{S}^2)$ and $v(\theta, \varphi) = v_{\theta}(\theta, \varphi)\hat{\theta} + v_{\varphi}(\theta, \varphi)\hat{\varphi}$ with $v_{\theta}, v_{\varphi} \in C^{\infty}(\mathbb{S}^2)$. Then

$$\int_{\mathbb{S}^2} \frac{\partial f}{\partial \theta} v_{\theta} d\sigma = \int_{\mathbb{S}^2} \nabla f \cdot \hat{\theta} v_{\theta} d\sigma = -\int_{\mathbb{S}^2} f \nabla \cdot (v_{\theta} \hat{\theta}) d\sigma,$$
$$\int_{\mathbb{S}^2} \frac{1}{\sin \theta} \frac{\partial f}{\partial \varphi} v_{\varphi} d\sigma = \int_{\mathbb{S}^2} \nabla f \cdot \hat{\varphi} v_{\varphi} d\sigma = -\int_{\mathbb{S}^2} f \nabla \cdot (v_{\varphi} \hat{\varphi}) d\sigma.$$

Adding these identities we get

$$\int_{\mathbb{S}^2} \nabla f \cdot V \, d\boldsymbol{\sigma} = -\int_{\mathbb{S}^2} f \, \nabla \cdot V \, d\boldsymbol{\sigma} \tag{19}$$

for any vector field $V \in C^{\infty}(\mathbb{S}^2 \to T(\mathbb{S}^2))$ where $T(\mathbb{S}^2)$ is the tangent bundle of the smooth manifold \mathbb{S}^2 . Motivated by (19), f is weakly differentiable if there is a vector field $\vartheta_f \in L^1(\mathbb{S}^2 \to T(\mathbb{S}^2))$ such that

$$\int_{\mathbb{S}^2} \vartheta_f \cdot V \, d\boldsymbol{\sigma} = -\int_{\mathbb{S}^2} f \, \nabla \cdot V \, d\boldsymbol{\sigma}, \quad \forall V \in C^{\infty} \left(\mathbb{S}^2 \to T \left(\mathbb{S}^2 \right) \right). \tag{20}$$

This, unique up to a set of zero measure, vector field ϑ_f is the weak surface gradient of f. According to ([3], Proposition 3.2., page 15)

$$H^{1}(\mathbb{S}^{2}) = W^{1,2}(\mathbb{S}^{2}) := \left\{ f \in L^{2}(\mathbb{S}^{2}) : |\vartheta_{f}| \in L^{2}(\mathbb{S}^{2}) \right\}.$$

We start with (5). By Theorem 1, it suffices to prove the existence of a sequence $\{f_n\}_{n\geq 1}$ in $H^1(\mathbb{S}^2)$ such that

$$\lim_{n \to \infty} \frac{Q(f_n; \phi)}{T_A(f_n)} = 1.$$
(21)

Consider the functions

$$f_n(\theta, \varphi) := \phi(\theta)^{\frac{1}{2} - \frac{1}{n}}.$$
(22)

The functions f_n are independent of φ , hence

$$\frac{\mathcal{Q}(f_n;\phi)}{T_A(f_n)} = \frac{\int_0^{\pi} \frac{f_n^2 \sin\theta}{\theta^2 \phi^2(\theta)} d\theta + \int_0^{\pi} \frac{f_n^2 \sin\theta}{(\pi - \theta)^2 \phi^2(\pi - \theta)} d\theta}{4\int_0^{\pi} \left(\frac{\partial f_n}{\partial \theta}\right)^2 \sin\theta d\theta + A\int_0^{\pi} f_n^2 \sin\theta d\theta}$$
(23)

where the derivative $\partial f_n/\partial \theta$ is understood in the week sense discussed above. Since $\phi \in L^1_{\text{loc}}(\mathbb{R})$ and $\phi \ge 1$ on $[0,\pi]$, then

$$\int_0^{\pi} f_n^2 \sin \theta d\theta = \int_0^{\pi} \phi(\theta)^{1-\frac{2}{n}} \sin \theta d\theta \leqslant \int_0^{\pi} \phi(\theta) d\theta \approx 1.$$
(24)

Thus $f_n \in L^2(\mathbb{S}^2)$ for all $n \ge 1$. Notice also that f_n is smooth on $[0,\pi] \setminus \{0\}$ and its weak derivative

$$\frac{\partial f_n}{\partial \theta} = \frac{\frac{1}{n} - \frac{1}{2}}{\theta \phi^{\frac{1}{2} + \frac{1}{n}}}.$$
(25)

Therefore

$$\int_0^{\pi} \left(\frac{\partial f_n}{\partial \theta}\right)^2 \sin \theta d\theta = \frac{a_n}{4} \int_0^{\pi} \frac{1}{\theta \, \phi^{1+\frac{2}{n}}} \frac{\sin \theta}{\theta} d\theta, \quad a_n := \left(1 - \frac{2}{n}\right)^2.$$

And since $\int_0^{\pi} \frac{d\theta}{\theta \phi^{1+\frac{2}{n}}} = \frac{n}{2}$, $\sin \theta \leq \theta$, then $\partial f_n / \partial \theta \in L^2(\mathbb{S}^2)$ for all $n \ge 1$. Substituting for f_n from (22) and for $\partial f_n / \partial \theta$ from (25) into (23) implies

$$\frac{Q(f_n;\phi)}{T_A(f_n)} = \frac{\alpha_n + \beta_n}{a_n \alpha_n + \gamma_n} = \frac{1}{a_n} \left(1 + \frac{\beta_n - \gamma_n/a_n}{\alpha_n + \gamma_n/a_n} \right)$$
(26)

where

$$\begin{aligned} \alpha_n &:= \int_0^\pi \frac{1}{\theta \, \phi^{1+\frac{2}{n}}} \frac{\sin \theta}{\theta} d\theta, \\ \beta_n &:= \int_0^\pi \frac{\phi^{1-\frac{2}{n}}(\theta) \sin \theta}{(\pi - \theta)^2 \, \phi^2 \, (\pi - \theta)} d\theta, \\ \gamma_n &:= A \int_0^\pi \phi^{1-\frac{2}{n}} \sin \theta d\theta. \end{aligned}$$

Observe that $\lim_{n\to+\infty} a_n = 1$. We shall show that, while $\lim_{n\to+\infty} \alpha_n = +\infty$, the sequences $\{\beta_n\}_{n\geq 1}$ and $\{\gamma_n\}_{n\geq 1}$ are both convergent. Using this in (26) proves (21).

Exploiting the continuity and positivity of $\sin \theta / \left(\theta^2 \phi^{1+\frac{2}{n}}\right)$ on $[\pi/2,\pi]$, then applying the inequality $\sin \theta / \theta \ge 2/\pi$ when $0 \le \theta \le \pi/2$, we obtain

$$\alpha_n = \int_0^{\pi/2} \frac{1}{\theta \phi^{1+\frac{2}{n}}} \frac{\sin \theta}{\theta} d\theta + \int_{\pi/2}^{\pi} \frac{\sin \theta}{\theta^2 \phi^{1+\frac{2}{n}}} d\theta$$
$$\geqslant \frac{2}{\pi} \int_0^{\pi/2} \frac{1}{\theta \phi^{1+\frac{2}{n}}} d\theta = \frac{n}{\pi (1 + \log(2))^{\frac{2}{n}}}.$$
(27)

This proves the divergence of $\{\alpha_n\}$. Next, by the dominated convergence theorem and (24) we readily find

$$\lim_{n \to +\infty} \gamma_n = A \lim_{n \to +\infty} \int_0^{\pi} \phi^{1-\frac{2}{n}}(\theta) \sin \theta d\theta = \int_0^{\pi} \phi(\theta) \sin \theta d\theta \lesssim 1.$$

Finally, since $\theta \mapsto \sin \theta / ((\pi - \theta)^2 \phi^2 (\pi - \theta)) \in C([0, \pi/2])$, then using the local integrability of ϕ and the dominated convergence theorem again implies

$$\lim_{n \to \infty} \int_0^{\pi/2} \frac{\phi^{1-\frac{1}{n}}(\theta) \sin \theta}{(\pi-\theta)^2 \phi^2(\pi-\theta)} d\theta = \int_0^{\pi/2} \frac{\phi(\theta) \sin \theta}{(\pi-\theta)^2 \phi^2(\pi-\theta)} d\theta \lesssim 1.$$
(28)

Furthermore, since $\phi \in C([\pi/2,\pi])$, and $\frac{\sin\theta}{\pi-\theta} = \frac{\sin(\pi-\theta)}{\pi-\theta} \leq 1$, on $[\pi/2,\pi]$, then

$$\int_{\pi/2}^{\pi} \frac{\phi^{1-\frac{1}{n}}(\theta)\sin\theta}{(\pi-\theta)^2 \phi^2(\pi-\theta)} d\theta \lesssim \int_{\pi/2}^{\pi} \frac{d\theta}{(\pi-\theta) \phi^2(\pi-\theta)} \approx 1.$$
(29)

The convergence of $\{\beta_n\}$ follows from (28) together with (29).

The proof of (6) shares the main idea of (5). The functions $g_n(\theta, \varphi) := \psi(\theta)^{\frac{1}{2} - \frac{1}{n}} \in L^2(\mathbb{S}^2)$, $n \ge 1$, and satisfy $\lim_{n \to \infty} \frac{Q(g_n; \psi)}{T_B(g_n)} = 1$. Indeed, we have

$$\frac{Q(g_n; \psi)}{T_B(g_n)} = \frac{\int_0^{\pi} \frac{g_n^2 \sin \theta}{\theta^2 \psi^2(\theta)} d\theta + \int_0^{\pi} \frac{g_n^2 \sin \theta}{(\pi - \theta)^2 \psi^2(\pi - \theta)} d\theta}{4 \int_0^{\pi} \left(\frac{\partial g_n}{\partial \theta}\right)^2 \sin \theta d\theta + B \int_0^{\pi} g_n^2 \sin \theta d\theta}$$
$$= \frac{\tilde{\alpha}_n}{a_n \tilde{\alpha}_n + \tilde{\beta}_n} = \frac{1}{a_n} \left(1 - \frac{\tilde{\beta}_n/a_n}{\tilde{\alpha}_n + \tilde{\beta}_n/a_n}\right)$$

where

$$\tilde{\alpha}_n := \int_0^\pi \frac{\sin \theta \, d\theta}{\theta^2 \, \psi^{1+\frac{2}{n}}} + \int_0^\pi \frac{\sin \theta \, d\theta}{(\pi - \theta)^2 \, \psi^{1+\frac{2}{n}}} = 2 \int_0^\pi \frac{\sin \theta \, d\theta}{\theta^2 \, \psi^{1+\frac{2}{n}}}$$
$$\tilde{\beta}_n := B \int_0^\pi \psi^{1-\frac{2}{n}} \sin \theta \, d\theta - a_n \int_0^\pi M(\theta) \frac{\sin \theta}{\psi^{1+\frac{2}{n}}} \, d\theta.$$

Similarly to (27), we have

$$\tilde{\alpha}_n = 2 \int_0^1 \frac{\sin\theta}{\theta^2} \frac{1}{\psi^{1+\frac{2}{n}}} d\theta + 2 \int_1^\pi \frac{\sin\theta}{\theta^2} \frac{1}{\psi^{1+\frac{2}{n}}} d\theta$$
$$\geqslant 2 \int_0^1 \frac{\sin\theta}{\theta^2} \frac{1}{\psi^{1+\frac{2}{n}}} d\theta = 2 \int_0^1 \frac{\sin^2\theta}{\theta^2 \cos\theta} \frac{1}{\psi^{1+\frac{2}{n}}} \frac{\cos\theta}{\sin\theta} d\theta$$
$$\geqslant \frac{8}{\pi^2} \int_0^1 \frac{1}{\psi^{1+\frac{2}{n}}} \frac{\cos\theta}{\sin\theta} d\theta = \frac{4n}{\pi^2} \frac{1}{(1+\log\pi)^{\frac{2}{n}}}.$$

Hence $\lim_{n\to\infty} \tilde{\alpha}_n = \infty$. Recall from (17) and (18) that $M \in C([0,\pi])$. Also, since $\psi \in L^1_{\text{loc}}(\mathbb{R}), \ \psi > 1$ uniformly, then $\lim_{n\to\infty} \tilde{\beta}_n$ exists by the dominated convergence theorem. \Box

6. Proof of Theorem 3

Proof. The transition to the inequalities (3) and (4) from their respective stronger versions, (15) and (16), comes from the bounds

$$\int_{\mathbb{S}^2} F(\theta) f^2 d\sigma \leqslant A \int_{\mathbb{S}^2} f^2 d\sigma, \quad \int_{\mathbb{S}^2} G(\theta) f^2 d\sigma \leqslant B \int_{\mathbb{S}^2} f^2 d\sigma$$

where the bounded functions F and G are both positive and independent of f. Interestingly, as seen in Section 5, the size of $0 < A, B < \infty$ played no role in optimising (3) and (4).

Up to the inequality (15) or (16) an equality relation persists except for the only inequality (11). So a sufficient and necessary condition for an equality in (15) or (16) (and a necessary condition for an equality in (3) and (4)) is an equality in (11). But an equality in (11) occurs if and only if

$$\int_{\mathbb{S}^2} \psi |\nabla g|^2 d\sigma = 0.$$
(30)

Recalling that $g = f/\sqrt{\psi}$, we compute

$$\begin{split} \psi |\nabla g|^{2} &= \psi \left| \frac{\nabla f}{\sqrt{\psi}} - \frac{1}{2} \frac{f}{\psi^{2}} \frac{\partial \psi}{\partial \theta} \hat{\theta} \right|^{2} \\ &= |\nabla f|^{2} - \frac{f}{\psi} \frac{\partial \psi}{\partial \theta} \nabla f \cdot \hat{\theta} + \frac{1}{4} \frac{f^{2}}{\psi^{2}} \left(\frac{\partial \psi}{\partial \theta} \right)^{2} \\ &= |\nabla f|^{2} - \left(\frac{\partial f}{\partial \theta} \right)^{2} + \left(\frac{\partial f}{\partial \theta} \right)^{2} - \frac{f}{\psi} \frac{\partial \psi}{\partial \theta} \frac{\partial f}{\partial \theta} + \frac{1}{4} \frac{f^{2}}{\psi^{2}} \left(\frac{\partial \psi}{\partial \theta} \right)^{2} \\ &= |\nabla f|^{2} - \left(\frac{\partial f}{\partial \theta} \right)^{2} + \left(\frac{\partial f}{\partial \theta} - \frac{1}{2} \frac{f}{\psi} \frac{\partial \psi}{\partial \theta} \right)^{2}. \end{split}$$
(31)

Since $|\nabla f|^2 - \left(\frac{\partial f}{\partial \theta}\right)^2 = \frac{1}{\sin^2 \theta} \left(\frac{\partial f}{\partial \varphi}\right)^2 \ge 0$, then, by (31), the equality (30) is equivalent to

$$\int_{\mathbb{S}^2} |\nabla f|^2 - \left(\frac{\partial f}{\partial \theta}\right)^2 d\sigma = \int_{\mathbb{S}^2} \left(\frac{\partial f}{\partial \theta} - \frac{1}{2} \frac{f}{\psi} \frac{\partial \psi}{\partial \theta}\right)^2 d\sigma = 0.$$
(32)

The equalities (32) are, in their turn, equivalent to

$$\frac{1}{\sin\theta} \left| \frac{\partial f}{\partial \varphi} \right| = \left| \frac{\partial f}{\partial \theta} - \frac{1}{2} \frac{f}{\psi} \frac{\partial \psi}{\partial \theta} \right| = 0.$$
(33)

Suppose that f is not the zero function. Then (33) are possible if and only if

$$f = f(\theta), \quad \frac{df}{f} = \frac{1}{2} \frac{d\psi}{\psi}.$$

That is $f = c\sqrt{\psi}$, c is a constant. But such $f \notin H^1(\mathbb{S}^2)$ because

$$\begin{split} \int_{\mathbb{S}^2} |\nabla f|^2 d\sigma &= 2\pi \int_0^\pi \left(\frac{\partial f}{\partial \theta}\right)^2 \sin \theta \, d\theta \gtrsim \int_0^1 \frac{\cos^2 \theta}{\sin \theta} \frac{1}{\psi} d\theta \\ &\gtrsim \int_0^1 \frac{d\theta}{\sin \theta \, \phi(\sin \theta)} \approx \int_0^1 \frac{d\theta}{\theta \, \phi(\theta)} = +\infty. \quad \Box \end{split}$$

7. Proof of Theorem 4

Proof. Write

$$\frac{1}{\theta}\frac{1}{\phi^{2}\left(\theta\right)}=\nabla\left(\frac{1}{\phi\left(\theta\right)}\right)\cdot\hat{\theta}$$

Assume that f is smooth. Then integrating by parts w.r.t. the surface measure σ we get

$$\begin{split} \int_{\mathbb{S}^2} \frac{f^2}{\theta^2 \phi^2(\theta)} d\sigma &= \int_{\mathbb{S}^2} \nabla \left(\frac{1}{\phi(\theta)} \right) \cdot \frac{f^2}{\theta} \hat{\theta} d\sigma \\ &= -\int_{\mathbb{S}^2} \frac{1}{\phi(\theta)} \nabla \cdot \left(\frac{f^2}{\theta} \hat{\theta} \right) d\sigma \\ &= -2 \int_{\mathbb{S}^2} \frac{f \nabla f \cdot \hat{\theta}}{\theta \phi(\theta)} d\sigma + \int_{\mathbb{S}^2} \frac{f^2}{\theta^2 \phi(\theta)} d\sigma - \int_{\mathbb{S}^2} \frac{f^2}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta} d\sigma. \end{split}$$
(34)

Observe here that each of the last two integrals on the right hand side of (34) can diverge. They suffer nonintegrable singularities at $\theta = 0$. But, when put together, their sum

$$I := \int_{\mathbb{S}^2} \frac{f^2}{\theta^2 \phi(\theta)} d\sigma - \int_{\mathbb{S}^2} \frac{f^2}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta} d\sigma = \int_{\mathbb{S}^2} \frac{1}{\theta \phi(\theta)} \left(\frac{1}{\theta} - \frac{\cos \theta}{\sin \theta}\right) f^2 d\sigma \quad (35)$$

is convergent. In fact

$$\lim_{\theta \to 0^{+}} \frac{1}{\theta \phi(\theta)} \left(\frac{1}{\theta} - \frac{\cos \theta}{\sin \theta} \right) = 0.$$

Also, $\theta \mapsto 1/(\theta^2 \phi(\theta))$ is continuous on a neighborhood of $\theta = \pi$. Furthermore, if we fix $\delta > 0$ and let $D := \{x(\theta, \varphi) \in \mathbb{S}^2 : 0 \le \theta < \delta\}$, then the integral

$$\int_{\mathbb{S}^2 \setminus D} \frac{f^2}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta} d\sigma$$

does exist. Unfortunately, we can not control the integral I by $\int_{\mathbb{S}^2} f^2 d\sigma$, up to a constant factor. The reason is

$$\lim_{\theta \to \pi^{-}} \frac{1}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta} = -\infty.$$

But since

$$\lim_{\theta \to \pi^{-}} \left(\frac{1}{\theta \phi(\theta)} \frac{\cos \theta}{\sin \theta} + \frac{1}{\pi} \frac{1}{(\pi - \theta)} \right) = 0$$

then, we may introduce the convergent integral $J := \frac{1}{\pi} \int_{\mathbb{S}^2} \frac{f^2}{\pi - \theta} d\sigma$ to the integral I to get

$$I = I - J + J = \int_{\mathbb{S}^2} K(\theta) f^2 d\sigma + J$$
(36)

where

$$K(\theta) := \frac{1}{\theta \phi(\theta)} \left(\frac{1}{\theta} - \frac{\cos \theta}{\sin \theta} \right) - \frac{1}{\pi} \frac{1}{(\pi - \theta)}$$

By the continuity of *K* on $]0,\pi[$ and since

$$\lim_{\theta \to 0^+} K(\theta) = -\lim_{\theta \to \pi^-} K(\theta) = -\frac{1}{\pi^2}$$

then K is bounded on $[0, \pi]$. Actually, K is monotonically increasing. Thus

$$\sup_{[0,\pi]} |K| = \frac{1}{\pi^2}.$$
(37)

Using (37) in (36) we deduce that

$$I \leqslant \frac{1}{\pi^2} \int_{\mathbb{S}^2} f^2 \, d\sigma + J. \tag{38}$$

Returning with (38) to the inequality (34) in the light of (35) we obtain

$$\int_{\mathbb{S}^2} \frac{f^2}{\theta^2 \phi^2(\theta)} d\sigma \leqslant -2 \int_{\mathbb{S}^2} \frac{f \nabla f \cdot \hat{\theta}}{\theta \phi(\theta)} d\sigma + \frac{1}{\pi^2} \int_{\mathbb{S}^2} f^2 d\sigma + \frac{1}{\pi} \int_{\mathbb{S}^2} \frac{f^2}{\pi - \theta} d\sigma.$$
(39)

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Applying Cauchy's inequality with an ε we find

$$-2\int_{\mathbb{S}^2} \frac{f\nabla f.\hat{\theta}}{\theta\phi(\theta)} d\sigma \leq 2\varepsilon \int_{\mathbb{S}^2} \frac{f^2}{\theta^2 \phi^2(\theta)} d\sigma + \frac{1}{2\varepsilon} \int_{\mathbb{S}^2} |\hat{\theta}.\nabla f|^2 d\sigma.$$
(40)

Therefore, it follows from (39) and (40) that

$$2\varepsilon(1-2\varepsilon)\int_{\mathbb{S}^2} \frac{f^2}{\theta^2 \phi^2(\theta)} d\sigma \leqslant \int_{\mathbb{S}^2} |\hat{\theta}.\nabla f|^2 d\sigma + \frac{2\varepsilon}{\pi^2} \int_{\mathbb{S}^2} f^2 d\sigma + \frac{2\varepsilon}{\pi} \int_{\mathbb{S}^2} \frac{f^2}{\pi - \theta} d\sigma, \quad 0 < \varepsilon < \frac{1}{2}.$$
 (41)

The choice $\varepsilon = 1/4$ maximizes the factor $2\varepsilon(1-2\varepsilon)$ and, consequently, the left hand side of (41). This proves (7). The inequality (8) can be obtained analogously.

In the fashion of the proof of Theorem 2, the sequence $f_n = \phi^{\frac{1}{2} - \frac{1}{n}}$ clearly satisfies

$$\lim_{n \to \infty} \frac{\frac{1}{4} \int_0^{\pi} \frac{f_n^2}{\rho_{\phi}^2(\theta)} \sin \theta \, d\theta}{U(f_n) + \frac{1}{2\pi} \int_0^{\pi} \frac{f_n^2}{\pi - \theta} \sin \theta \, d\theta} = \lim_{n \to \infty} \frac{\frac{1}{4} \int_0^{\pi} \frac{f_n^2}{\rho_{\phi}^2(\pi - \theta)} \sin \theta \, d\theta}{U(f_n) + \frac{1}{2\pi} \int_0^{\pi} \frac{f_n^2}{\theta} \sin \theta \, d\theta} = 1$$

where

$$U(f) = \int_0^{\pi} \left(\frac{\partial f}{\partial \theta}\right)^2 \sin \theta \, d\theta + \frac{1}{2\pi^2} \int_0^{\pi} f^2 \sin \theta \, d\theta$$

One only needs to inspect the convergence of $\int_0^{\pi} \left(\phi^{1-\frac{2}{n}} \sin \theta / \theta \right) d\theta$, $\int_0^{\pi} \left(\phi^{1-\frac{2}{n}} \sin \theta / (\pi - \theta) \right) d\theta$ as $n \to \infty$. This is obvious from the bound $\sin \theta \leq \min\{\theta, \pi - \theta\}$ on $[0, \pi]$ and the fact $\phi \in L^1([0, \pi])$.

Finally, careful review of the proof of (7) above reveals that a necessary condition for a function $f \in H^1(\mathbb{S}^2) \setminus \{0\}$ to achieve an equality in (7) is that it yields an equality in (40). This is equivalent to

$$\nabla f.\hat{\theta} = -\frac{1}{2} \frac{f}{\theta \,\phi(\theta)}.\tag{42}$$

Suppose (42) was true. Then by (34) and (35) we must have

$$\int_{\mathbb{S}^2} \frac{h(\theta) f^2}{\theta \phi(\theta)} d\sigma = 0 \tag{43}$$

where

$$h(\theta) := \frac{1}{\theta} - \frac{\cos\theta}{\sin\theta}$$

On the other hand

$$\lim_{\theta \to 0^+} h(\theta) = 0, \quad h'(\theta) = \frac{\theta^2 - \sin^2 \theta}{\theta^2 \sin^2 \theta} > 0, \ 0 < \theta < \pi.$$

This shows *h* is strictly positive on $]0,\pi]$ and since $\theta\phi(\theta) \ge 0$ then (43) is a contradiction. \Box

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A. A. ABDELHAKIM

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Ahmed A. Abdelhakim Mathematics Department, Faculty of Science Assiut University Assiut 71516, Egypt e-mail: ahmed.abdelhakim@aun.edu.eg