# ON A GENERALIZED EGNELL INEQUALITY 

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Abstract. In this paper we prove an inequality which connects the $L^{p}$ norm of the gradient of a function $u$ with its $|x|^{v}$-weighted $L^{\frac{p(N+v)}{N-p}}$ norm and its $L^{p^{*}}$-weak norm. Here $1<p<N$, $-p<v \leqslant 0$ and $p^{*}=\frac{N p}{N-p}$. As a consequence we can provide an alternative proof of the Egnell inequality in $\mathbb{R}^{N}$.

## 1. Introduction

The classical Sobolev inequality in $\mathbb{R}^{N}$ asserts that

$$
\begin{equation*}
\left(\int_{\mathbb{R}^{N}}|\nabla u|^{p} d x\right)^{\frac{1}{p}} \geqslant S(N, p)\left(\left.\int_{\mathbb{R}^{N}}|u|\right|^{p^{*}} d x\right)^{\frac{1}{p^{*}}}, \tag{1}
\end{equation*}
$$

with $1<p<N, p^{*}=\frac{N p}{N-p}$ and $u$ is a real-valued function in $L^{p^{*}}\left(\mathbb{R}^{N}\right)$ such that $|\nabla u| \in L^{p}\left(\mathbb{R}^{N}\right)$, where $\nabla u$ is the distributional gradient of $u$. The value of the sharp constant $S(N, p)$ in (1) is known to be

$$
S(N, p)=\pi^{\frac{1}{2}} 2^{\frac{1}{N}} N^{\frac{1}{p}}(N-p)^{\frac{p-1}{p}}(p-1)^{\frac{1}{N}-\frac{p-1}{p}} p^{-\frac{1}{N}}\left[\frac{\Gamma\left(\frac{N}{p}\right) \Gamma\left(N-\frac{N}{p}\right)}{\Gamma(N) \Gamma\left(\frac{N}{2}\right)}\right]^{\frac{1}{N}},
$$

where $\Gamma$ is the standard Euler function. The equality sign holds when $u$ is of the form

$$
u(x)=\frac{h}{\left[1+k|x|^{\frac{p}{p-1}}\right]^{\frac{N-p}{p}}}, \quad h, k>0
$$

with $h, k$ positive constants (see [6], [8] and [32]).
When $\mathbb{R}^{N}$ is replaced by a bounded domain $\Omega \subset \mathbb{R}^{N}$ and $u \in W_{0}^{1, p}(\Omega)$, the Sobolev inequality (1) still holds with $S(N, p)$ as best constant, but the constant is never achieved. For this reason, several authors studied the problem of improving the inequality (1) for $u \in W_{0}^{1, p}(\Omega)$, by adding a right-hand-side remainder term. The first

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results in this direction are given in [10] and then in [9], where the authors prove several improvements of (1) by adding the norm of $u$ and of $\nabla u$ in suitable $L^{q}$ spaces. Similar results are still true when we consider $u \in W^{1, p}(\Omega)$ vanishing on a fixed part $\Gamma_{0}$ of the boundary $\partial \Omega$ (see [19], [28], [29]).

Analogous questions can be studied in relation to the celebrated Hardy-Sobolev inequality (see [26] and [25])

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}|\nabla u|^{p} d x \geqslant A(N, p) \int_{\mathbb{R}^{N}} \frac{|u|^{p}}{|x|^{p}} d x \tag{2}
\end{equation*}
$$

with $u \in W^{1, p}\left(\mathbb{R}^{N}\right)$ and

$$
\begin{equation*}
A(N, p)=\left(\frac{N-p}{p}\right)^{p} \tag{3}
\end{equation*}
$$

This inequality and its various improvements are used in many contexts, as in the study of stability of solutions to semi-linear elliptic and parabolic equations (see [11], [12], [33]), or in the analysis of the asymptotic behaviour of the heat equation with singular potentials (see [34]). When $\mathbb{R}^{N}$ is replaced by a bounded domain $\Omega$ of $\mathbb{R}^{N}$ and $u \in W_{0}^{1, p}(\Omega)$, the Hardy-Sobolev inequality (2) still holds.

The constant $A(N, p)$ in (3) is the best one in both cases but there is no function $u \in W^{1, p}\left(\mathbb{R}^{N}\right)$ (or $u \in W_{0}^{1, p}(\Omega)$ ) for which it is achieved. For this reason several authors have improved inequality (2) by adding at the right-hand side a non-negative correction term (see e.g. [1], [2], [5], [7], [11], [15], [17], [20], [21], [22], [23], [30]).

The Sobolev and Hardy-Sobolev inequalities described so far represent a special case of a more general inequality, known as Egnell inequality (see [28]).
Let us consider $1<p<N,-p<v \leqslant 0$ and

$$
\begin{equation*}
q=\frac{p(N+v)}{N-p} \tag{4}
\end{equation*}
$$

If we denote by $L^{q}\left(\mathbb{R}^{N},|x|^{v}\right)$ the space of measurable functions $u$ such that

$$
\|u\|_{q,|x|^{v}}:=\left(\int_{\mathbb{R}^{N}}|u|^{q}|x|^{v} d x\right)^{\frac{1}{q}}<\infty
$$

then the Egnell inequality states that

$$
\begin{equation*}
\left(\int_{\mathbb{R}^{N}}|\nabla u|^{p} d x\right)^{\frac{1}{p}} \geqslant C(p, v)\left(\int_{\mathbb{R}^{N}}|u|^{q}|x|^{v} d x\right)^{\frac{1}{q}} \tag{5}
\end{equation*}
$$

The optimal value of $C(p, v)$ is obtained in [18] by means of methods similar to the ones used by Talenti in [32] to get the best constant in Sobolev inequality (see also [27]).

To write down the value of $C(p, v)$, let us consider for any $u \in W^{1, p}\left(\mathbb{R}^{N}\right), u \not \equiv 0$ the functional

$$
F(u):=\frac{\|\mid \nabla u\|_{L^{p}}}{\|u\|_{q,|x|^{v}}}
$$

then the best constant $C(p, v)$ is defined as

$$
C(p, v):=\inf _{\substack{u \in W^{1, p}\left(\mathbb{R}^{N}\right) \\ u \neq 0}} F(u)
$$

In [18] Egnell proves that the value of $C(p, v)$ is given by

$$
\begin{align*}
C(p, v)= & \pi^{\frac{N}{2} \frac{p+v}{p(N+v)}} 2^{\frac{p+v}{p(N+v)}}(N+v)^{\frac{1}{p}}(N-p)^{\frac{p-1}{p}}(p-1)^{-\frac{p-1}{p}+\frac{p+v}{p(N+v)}}  \tag{6}\\
& \times(p+v)^{-\frac{p+v}{p(N+v)}}\left[\frac{\Gamma\left(\frac{(N+v)(p-1)}{p+v}\right) \Gamma\left(\frac{N+v}{p+v}\right)}{\Gamma\left(\frac{p(N+v)}{p+v}\right) \Gamma\left(\frac{N}{2}\right)}\right]^{\frac{p+v}{p(N+v)}}
\end{align*}
$$

and the infimum of $F(u)$ is attained when

$$
\begin{equation*}
u(x)=\frac{h}{\left[1+k|x|^{\frac{p+v}{p-1}}\right]^{\frac{N-p}{p+v}}} \tag{7}
\end{equation*}
$$

with $h$ e $k$ positive constants.
We highlight that the Egnell inequality becomes the Sobolev inequality when $v=0$ and $C(p, 0)=S(N, p)$. In terms of embedding between functional spaces, the Egnell inequality represents the continuous embedding of the Sobolev space $W^{1, p}\left(\mathbb{R}^{N}\right)$ in the weighted Lebesgue space $L^{q}\left(\mathbb{R}^{N},|x|^{v}\right)$, where $q$ is given in (4). We finally remark that the Egnell inequality can be read as a particular case of a class of interpolation inequalities known as Caffarelli-Kohn-Nirenberg inequalities (see [13] and [14]). Improved Caffarelli-Kohn-Nirenberg inequalities are widely studied in the literature.

In this paper we prove an inequality which connects the $L^{p}\left(\mathbb{R}^{N}\right)$ norm of the gradient a function $u$ with the weighted $L^{q}\left(\mathbb{R}^{N},|x|^{v}\right)$ and $L^{p^{*}}$-weak norms of $u$. Namely, we prove the following inequality:

$$
\begin{equation*}
\|u\|_{p^{*}, \infty}^{r p *}\|\mid \nabla u\|_{L^{p}}^{p} \geqslant A(N, p)\|u\|_{q,|x|^{v}}^{q}+B(N, p)\|u\|_{p *, \infty}^{s p^{*}}, \tag{8}
\end{equation*}
$$

where $A(N, p)$ and $q$ are given by (3) and (4) respectively,

$$
\begin{equation*}
r=\frac{p+v}{N}, \quad s=\frac{N+v}{N} \tag{9}
\end{equation*}
$$

and

$$
\begin{equation*}
B(N, p)=2 \pi^{\frac{N}{2}}(N-p)^{p-1}(p-1)^{2-p-\frac{(N-p)(p-1)}{p+v}} p^{\frac{p(N-p)}{p+v}}\left[\frac{\Gamma\left(\frac{(N+v)(p-1)}{p+v}\right) \Gamma\left(\frac{N+v}{p+v}\right)}{\Gamma\left(\frac{p(N+v)}{p+v}\right) \Gamma\left(\frac{N}{2}\right)}\right] \tag{10}
\end{equation*}
$$

Finally, we show that Egnell inequality can be easily deduced by (8) as corollary.
In order to prove the main result, we firstly apply a symmetrization procedure by replacing $u$ with its rearrangement $u^{\#}$, which is spherically symmetric and decreases with respect to $|x|$. Then, the Pólya-Szego principle (see [32]) and Hardy-Littlewood
inequality (see [26]) ensure that the previous assumption is not restrictive. Finally, we apply classical arguments of one-dimensional Calculus of Variations (see [31]) and techniques similar to the ones used in [4] for the Sobolev inequality, which straightforwardly lead to inequality (8).

## 2. Definitions and main result

The main result of this paper is a generalized Egnell inequality which links the $L^{p}\left(\mathbb{R}^{N}\right)$ norm of the gradient a function $u$ with the weighted $L^{q}\left(\mathbb{R}^{N},|x|^{v}\right)$ and $L^{p *}$ weak norm of $u$, where $1<p<N,-p<v \leqslant 0, p^{*}=\frac{N p}{N-p}$ and $q$ is given in (4). The Egnell inequality (5) with its optimal value (6) easily follows as corollary.

We firstly recall the definition of spherically decreasing rearrangement of a function $u$ and some related properties.

DEFINITION 1. Let $\Omega$ be a measurable subset of $\mathbb{R}^{N}$ and $u: \Omega \rightarrow \mathbb{R}$ a measurable function in $\Omega$. The distribution function of $u$ is the decreasing map $\mu$ from $[0,+\infty[$ into $[0,+\infty[$ defined at any point $t \geqslant 0$ as the measure of a level set of $u$, $\{x \in \Omega:|u(x)|>t\}$. The decreasing rearrangement $u^{*}$ of $u$ is the distribution function of $\mu$

$$
u^{*}(s):=\sup \{t \geqslant 0: \mu(t)>s\}, \quad s \in(0,|\Omega|)
$$

The main property of rearrangements is the fact that the distribution of $u^{*}$ is $\mu$, in other words $u$ and $u^{*}$ are equidistribuited.

DEFINITION 2. Let us denote by $\omega_{N}$ the measure of the unit ball of $\mathbb{R}^{N}$ and by $\Omega^{\#}$ the ball of $\mathbb{R}^{N}$ centred at the origin such that $|\Omega|=\left|\Omega^{\#}\right|$. For every $x \in \Omega^{\#}$, the spherically decreasing rearrangement of $u$ is defined as the decreasing rearrangement $u^{*}$ valued in $\omega_{N}|x|^{N}$

$$
u^{\#}(x):=u^{*}\left(\omega_{N}|x|^{N}\right), \quad x \in \Omega^{\#}
$$

Obviously, $u^{\#}$ is decreasing and spherically symmetric; moreover $u$ and $u^{\#}$ are equidistributed, and the level set $\left\{x \in \Omega^{\#}:\left|u^{\#}(x)\right|>t\right\}$ is the ball centred at the origin and whose measure is $\mu(t)$. For an exhaustive treatment of rearrangements see, for example, [16] and [24]. Here we just recall the Hardy-Littlewood inequality

$$
\begin{equation*}
\int_{\Omega}|u(x) v(x)| d x \leqslant \int_{\Omega^{\#}} u^{\#}(x) v^{\#}(x) d x \tag{11}
\end{equation*}
$$

with $u, v$ measurable functions (see [26]), and the Pólya-Szego principle (see [32])

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}\left|\nabla u^{\#}(x)\right|^{p} d x \leqslant \int_{\mathbb{R}^{N}}|\nabla u(x)|^{p} d x \tag{12}
\end{equation*}
$$

where $u \in W_{0}^{1, p}\left(\mathbb{R}^{N}\right), 1<p<N$.

Now we recall the definitions of the Marcinkiewicz $L^{p}$-weak space, of Lorentz space and some of their elementary properties which will be used in the following sections.

Definition 3. Let $\Omega$ be an open subset of $\mathbb{R}^{N}$ and $0<p<\infty$. The Marcinkiewicz $L^{p}$-weak space consists of all measurable functions $u$ such that

$$
\begin{equation*}
\|u\|_{p, \infty}:=\sup _{t>0} \omega_{N}^{\frac{1}{p}}\left[t^{\frac{N}{P}} u^{*}(t)\right] \tag{13}
\end{equation*}
$$

Definition 4. Let $\Omega$ be an open subset of $\mathbb{R}^{N}$ and $0<p, q \leqslant \infty$. The Lorentz space $L(p, q)$ consists of all measurable functions $u$ such that

$$
\|u\|_{p, q}:= \begin{cases}\left(\int_{0}^{+\infty}\left[u^{*}(t) t^{\frac{1}{p}}\right]^{q} \frac{d t}{t}\right)^{\frac{1}{q}} & , 0<q<\infty  \tag{14}\\ \sup _{t>0} u^{*}(t) t^{\frac{1}{p}}, & q=\infty\end{cases}
$$

The Lorentz space $L(p, p)$ coincides with the Lebesgue space $L^{p}$ and

$$
\|u\|_{p, p}=\|u\|_{p}
$$

Moreover, it can be proved that the $L(p, \infty)$ space coincides with the Marcinkiewicz $L^{p}$ - weak space.

Inclusion relations among $L(p, q)$ spaces, with $p$ varying, are like those for the Lebesgue $L^{p}$ spaces, in that they depend on the structure of the underlying measure space. The secondary exponent $q$ is not involved. Thus, if $0<p<r \leqslant \infty$ and $0<$ $q, s \leqslant \infty$ then

$$
\begin{equation*}
L(r, s) \hookrightarrow L(p, q) \tag{15}
\end{equation*}
$$

For what concerns the secondary exponent, we have that if if $0<p \leqslant \infty$ and $0<q<s \leqslant \infty$, then

$$
\begin{equation*}
L(p, q) \hookrightarrow L(p, s) \tag{16}
\end{equation*}
$$

REMARK 1. Thanks to the Sobolev inequality in the Lorentz space $L\left(p^{*}, p\right)$ (see [3]) and Lorentz spaces properties (15) and (16), we get that for $1<p<N$ and $-p<$ $v \leqslant 0$ then

$$
W^{1, p}\left(\mathbb{R}^{N}\right) \hookrightarrow L\left(p^{*}, p\right) \hookrightarrow L\left(p^{*}, \frac{p(N+v)}{N-p}\right) \hookrightarrow L^{\frac{p(N+v)}{N-p}}\left(\mathbb{R}^{N},|x|^{v}\right)
$$

hence we can conclude that the Egnell inequality is more refined than the Sobolev one. The main result of the paper is stated in the following:

THEOREM 1. Let $u \in W^{1, p}\left(\mathbb{R}^{N}\right), 1<p<N$ and $-p<v \leqslant 0$. Then inequality (8) holds:

$$
\|u\|_{p^{*}, \infty}^{r p *}\|\mid \nabla u\|_{L^{p}}^{p} \geqslant A(N, p)\|u\|_{q,|x|^{v}}^{q}+B(N, p)\|u\|_{p^{*}, \infty}^{s p^{*}},
$$

where $A(N, p)$ and $B(N, p)$ are given in (3) and (10), while $q, r, s$ are given in (4) and (9) respectively.

## 3. Proof of Theorem 1

In this section we provide a detailed proof of Theorem 1. As stated in the introduction, Egnell inequality will then be deduced from inequality (8).

Proof. The first step consists in the reduction of the problem to a spherically symmetric one. We replace $u$ with $u^{\#}$ and for the sake of simplicity we keep calling it $u$. By Hardy-Littlewood inequality (11) and Pólya-Szego principle (12), the left hand side of (5) decreases, while the right side increases. This implies that it is enough to prove Theorem 1 only in the radial case. Moreover without loss of generality we can assume that $u \in C_{0}^{1}\left(\mathbb{R}^{N}\right)$, since this assumption can be removed by density.

Let us consider the functional

$$
\begin{equation*}
J(u):=\frac{\omega_{N}}{p} \int_{0}^{\infty}\left|u^{\prime}\right|^{p} r^{N-1} d r-\frac{\omega_{N}}{p} \frac{(N-p)^{p}}{(p-1)^{p-1}} a^{p+v} \int_{0}^{\infty} u^{\frac{p(N+v)}{N-p}} r^{N-1+v} d r \tag{17}
\end{equation*}
$$

and the related Euler equation

$$
\begin{equation*}
-\triangle_{p} u=\left(\frac{N-p}{p-1}\right)^{p-1} a^{p+v}(N+v) u^{\frac{p(N+v)}{N-p}-1} r^{v} \tag{18}
\end{equation*}
$$

It can be proved that the following one-parameter family of functions

$$
\begin{equation*}
u_{\mathcal{\varepsilon}}(r):=u_{\varepsilon}(|x|)=\frac{\varepsilon^{\frac{N-p}{p}}}{\left[1+(a \varepsilon|x|)^{\frac{p+v}{p-1}}\right]^{\frac{N-p}{p+v}}}, \quad \varepsilon>0 \tag{19}
\end{equation*}
$$

satisfy (18). The constant $a>0$ in (19) is a free constant that will be properly chosen in the following.

If we compute the $L^{q}\left(\mathbb{R}^{N},|x|^{v}\right)$-norm of these extremals, we get that

$$
\begin{equation*}
\left\|u_{\varepsilon}\right\|_{q,|x|^{v}}^{q}=2 \pi^{\frac{N}{2}} \frac{p-1}{p+v} a^{-N-v} \frac{\Gamma\left(\frac{(N+v)(p-1)}{p+v}\right) \Gamma\left(\frac{N+v}{p+v}\right)}{\Gamma\left(\frac{p(N+v)}{p+v}\right) \Gamma\left(\frac{N}{2}\right)} \tag{20}
\end{equation*}
$$

so all the functions of the family (19) have the same $L^{q}\left(\mathbb{R}^{N}|x|^{v}\right)$-norm, which is independent of $\varepsilon$.

The curve

$$
\begin{equation*}
y=\gamma_{a}(r)=\frac{(p-1)^{\frac{(p-1)(N-p)}{p(p+v)}}}{p^{\frac{N-p}{p+v}}}(a r)^{-\frac{N-p}{p}}, \quad r>0 \tag{21}
\end{equation*}
$$

is the envelope of the graphs $y=u_{\varepsilon}(r)$; these cover the region $T$ of the first quadrant which lies below the curve (21).

Let $u$ be a sufficiently smooth, compactly supported radial function, and let us consider its norm in the Marcinkiewicz space of the functions weakly $L^{p^{*}}$ :

$$
\|u\|_{p^{*}, \infty}=\sup _{r>0}\left[r^{\frac{N-p}{p}} v(r)\right] .
$$

If we choose

$$
\begin{equation*}
a=\frac{(p-1)^{\frac{p-1}{p+v}}}{p^{\frac{p}{p+v}}}\|u\|_{p^{*}, \infty}^{-\frac{p}{N-p}}, \tag{22}
\end{equation*}
$$

then we get that it is the minimum value such that $u(r) \leqslant \gamma_{a}(r)$ for all $r>0$; the corresponding envelope (21) is given by:

$$
\begin{equation*}
\gamma(r)=\|u\|_{p^{*}, \infty} r^{-\frac{N-p}{p}} \tag{23}
\end{equation*}
$$

For each $\varepsilon>0$ the graph of the extremal $y=u_{\varepsilon}(r)$ defined in (19) touches the envelope $\gamma(r)$ defined in (23) at a point which splits it into two curves, denoted with $C_{1}(\varepsilon)$ and $C_{2}(\varepsilon)$. By varying $\varepsilon, C_{1}(\varepsilon)$ and $C_{2}(\varepsilon)$ define two families of curves that are the trajectories of two different fields of extremals of the functional (17), and are both defined in the same region $T$. Let us denote by $\left(1, q_{1}(r, y)\right)$ the first field and by $\left(1, q_{2}(r, y)\right)$ the second one. We explicitly stress that $q_{1}(r, y)$ represents the slope of the extremal of the first family passing through $(r, y) ; q_{2}(r, y)$ has an analogous meaning. The dashed lines in Figure 1 (and Figure 2) represent some arcs of extremals $C_{1}(\varepsilon)$ (and $C_{2}(\varepsilon)$ ), obtained by varying $\varepsilon$.
Moreover, the envelope in (23) touches the graph of $u$ at least in a point $P=(\alpha, \gamma(\alpha))$, which splits the graph of $u$ itself into two arcs $\Gamma_{1}$ and $\Gamma_{2}$, as in Figure 1 and Figure 2. Finally, we denote by $C_{1}, C_{2}$, respectively, the arcs of the families $C_{1}(\varepsilon)$ and $C_{2}(\varepsilon)$ passing through such $P=(\alpha, \gamma(\alpha))$. In Figure 1 and 2 the graphs of the envelope $y=y(r)$ and the arcs $\Gamma_{1}$ and $\Gamma_{2}, C_{1}$ and $C_{2}$ are sketched in full lines.
At this stage we apply classical arguments of one-dimensional Calculus of Variations (see [31]). Let us denote with

$$
f\left(r, u, u^{\prime}\right):=\frac{\omega_{N}}{p} r^{N-1}\left(\left|u^{\prime}\right|^{p}-\frac{(N-p)^{p}}{(p-1)^{p-1}} a^{p+v} u^{q} r^{v}\right)
$$

then the functional $J(u)$ defined in (17) can be rewritten as

$$
J(u)=\int_{0}^{\alpha} f\left(r, u, u^{\prime}\right) d r+\int_{\alpha}^{\infty} f\left(r, u, u^{\prime}\right) d r=J_{1}(u)+J_{2}(u) .
$$

Our target is to show that the one parameter family of extremals $y=u_{\mathcal{E}}(r)$ minimizes the functional $J(u)$. We begin by estimating $J_{1}(u)$ from below and we embed it in the first field $\left(1, q_{1}(r, y)\right)$.

Since $f$ is convex with respect to the last variable, then the Weierstass condition

$$
\mathscr{E}\left(r, w, \xi, \xi_{1}\right)=f(r, w, \xi)-f\left(r, w, \xi_{1}\right)+\left(\xi_{1}-\xi\right) f_{v^{\prime}}\left(r, w, \xi_{1}\right) \geqslant 0
$$

is satisfied. As a consequence

$$
\begin{equation*}
J_{1}(u) \geqslant \int_{0}^{\alpha} f\left(r, u, q_{1}\right)+\left(u^{\prime}-q_{1}\right) f_{u^{\prime}}\left(r, u, q_{1}\right) d r . \tag{24}
\end{equation*}
$$

Moreover, the differential form

$$
\begin{equation*}
\zeta_{1}=\left[f\left(r, u, q_{1}\right)-q_{1} f_{u^{\prime}}\left(r, u, q_{1}\right)\right] d r+f_{u^{\prime}}\left(r, u, q_{1}\right) d u \tag{25}
\end{equation*}
$$

is exact (see [31]), so its integral along any closed path is equal to zero. We compute the integral of $\zeta_{1}$ along the closed path represented in Figure 1 consisting of the graphs of $C_{1}$ and of $\Gamma_{1}$ between the origin and $\alpha$, and the segment $\tau$ of the vertical axis delimited by the intersection points of $C_{1}$ and $\Gamma_{1}$ with the vertical axis.


Figure 1

As a consequence, the integral of the right-hand side in (24) equals the line integral of (25) along $\tau$, which is null, plus the integral line along the curve $C_{1}$, therefore

$$
\int_{0}^{\alpha} f\left(r, u, q_{1}\right)+\left(u^{\prime}-q_{1}\right) f_{u^{\prime}}\left(r, u, q_{1}\right) d r=\int_{0}^{\alpha} f\left(r, u_{\varepsilon}, u_{\varepsilon}^{\prime}\right)=J_{1}\left(u_{\varepsilon}\right)
$$

This implies that

$$
\begin{equation*}
J_{1}(u) \geqslant J_{1}\left(u_{\varepsilon}\right) \tag{26}
\end{equation*}
$$

In a similar way we get an estimate from below of $J_{2}(u)$ and we embed it in the second field $\left(1, q_{2}(r, y)\right)$. Let us consider the exact differential form

$$
\begin{equation*}
\zeta_{2}=\left[f\left(r, u, q_{2}\right)-q_{2} f_{u^{\prime}}\left(r, u, q_{2}\right)\right] d r+f_{u^{\prime}}\left(r, u, q_{2}\right) d u \tag{27}
\end{equation*}
$$

we integrate $\zeta_{2}$ between $\alpha$ and $\beta$ along the path sketched in Figure 2 delimited by $C_{2}, \Gamma_{2}, S_{\beta}$ and the segment of the horizontal axis between the intersection points of $\Gamma_{2}$ and $S_{\beta}$ with the axis itself.
An asymptotic argument allows us to prove that the line integral of (27) along the


Figure 2
vertical segment $S_{\beta}$ in Figure 2 is infinitesimal when $\beta$ goes to infinity. Therefore

$$
\begin{equation*}
J_{2}(u) \geqslant J_{2}\left(u_{\varepsilon}\right) \tag{28}
\end{equation*}
$$

From (26) and (28) we get

$$
J(u) \geqslant J\left(u_{\varepsilon}\right)
$$

Now we compute $J\left(u_{\varepsilon}\right)$. From Egnell inequality and (20) we deduce

$$
\begin{aligned}
J\left(u_{\mathcal{E}}\right) & =\left\|\mid \nabla u_{\mathcal{E}}\right\|_{L^{p}}^{p}-\frac{(N-p)^{p}}{(p-1)^{p-1}} a^{p+v}\left\|u_{\mathcal{E}}\right\|_{q,|x|^{v}}^{q}= \\
& =C(p, v)\left\|u_{\mathcal{E}}\right\|_{q,|x|^{v}}^{p}-\frac{(N-p)^{p}}{(p-1)^{p-1}} a^{p+v}\left\|u_{\mathcal{E}}\right\|_{q,|x|^{v}}^{q}= \\
& =\frac{2}{p} \pi^{\frac{N}{2}} a^{p-N} \frac{(N-p)^{p-1}}{(p-1)^{p-2}} \frac{\Gamma\left(\frac{(N+v)(p-1)}{p+v}\right) \Gamma\left(\frac{N+v}{p+v}\right)}{\Gamma\left(\frac{p(N+v)}{p+v}\right) \Gamma\left(\frac{N}{2}\right)} .
\end{aligned}
$$

Since $J(u) \geqslant J\left(u_{\varepsilon}\right)$, then

$$
\begin{align*}
\int_{\mathbb{R}^{N}}|\nabla u|^{p} d x \geqslant & \frac{(N-p)^{p}}{(p-1)^{p-1}} a^{p+v}\|u\|_{q,|x|^{v}}^{q}+ \\
& +2 \pi^{\frac{N}{2}} a^{p-N} \frac{(N-p)^{p-1}}{(p-1)^{p-2}} \frac{\Gamma\left(\frac{(N+v)(p-1)}{p+v}\right) \Gamma\left(\frac{N+v}{p+v}\right)}{\Gamma\left(\frac{p(N+v)}{p+v}\right) \Gamma\left(\frac{N}{2}\right)} \tag{29}
\end{align*}
$$

Taking into account the value of $a$ in (22) and using a density argument we get the result (8) with $A(N, p)$ and $B(N, p)$ given by (3) and (10) respectively.

## 4. Conclusions

The inequality (8) can be read as a generalization of the Egnell inequality (5), which can be deduced from it by a minimization argument as follows.

Proof. We start from inequality (29) with $a>0$ free constant and we rewrite the right-hand side in a more concise way setting $x=a^{p+v}$; the inequality becomes:

$$
\int_{\mathbb{R}^{N}}|\nabla u|^{p} d x \geqslant K x+H x^{-\frac{N-p}{p+v}}=\phi(x)
$$

where

$$
K=\frac{(N-p)^{p}}{(p-1)^{p-1}}\|u\|_{q,|x|^{v}}^{q}
$$

and

$$
H=2 \pi^{\frac{N}{2}} \frac{(N-p)^{p-1}}{(p-1)^{p-2}} \frac{\Gamma\left(\frac{(N+v)(p-1)}{p+v}\right) \Gamma\left(\frac{N+v}{p+v}\right)}{\Gamma\left(\frac{p(N+v)}{p+v}\right) \Gamma\left(\frac{N}{2}\right)}
$$

Since $\phi(x)$ reaches its minimum when

$$
x=\left[\frac{H(N-p)}{K(p+v)}\right]^{\frac{p+v}{N+v}}
$$

then we obtain the (5) and the optimal value of the Egnell constant (6).

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