# TRANSFERENCE METHOD FOR CONE-LIKE RESTRICTED SUMMABILITY OF THE TWO-DIMENSIONAL WALSH-LIKE SYSTEMS 

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(Communicated by T. Erdélyi)


#### Abstract

In the present paper we investigate the boundedness of the maximal operator of some $d$-dimensional means, provided that the set of the indeces is inside a cone-like set $L$. Applying some assumptions on the summation kernels $P_{n_{1}, \ldots, n_{d}}$ we state that the cone-like restricted maximal operator $T_{C L R}^{\gamma}$ is bounded from the Hardy space $H_{p}^{\gamma}$ to the Lebesgue space $L^{p}$ for $p>p_{0}$. In the end point $p_{0}$ assuming some natural conditions on one-dimensional kernels we show that the maximal operator $T_{C L R}^{\gamma}$ is not bounded from the Hardy space $H_{p_{0}}^{\gamma}$ to the Lebesgue space $L^{p_{0}}$.


## 1. Definitions and notation

We follow the standard notions of dyadic analysis introduced by Schipp, Simon, Wade and Pál in [20] (see also [1]). Let $\mathbb{N}$ denote the set of natural numbers and let $\mathbb{P}:=\mathbb{N} \backslash\{0\}$. The cyclic group of order 2 will be denoted by $\mathbb{Z}_{2}$. The topology is given by that every subset is open. The Haar measure on $\mathbb{Z}_{2}$ is given such that

$$
\mu(\{0\})=\mu(\{1\})=1 / 2 .
$$

Let $G$ be the complete direct product of countable infinite copies of the compact group $\mathbb{Z}_{2}$. The elements of $G$ are sequences of the form $x=\left(x_{0}, x_{1}, \ldots, x_{k}, \ldots\right)$ with components $x_{k} \in\{0,1\}(k \in \mathbb{N})$. On $G$ the group operation is the component-wise addition, the measure $\mu$ is the product measure and the topology is the product topology. Such compact Abelian group $G$ is called the Walsh group.

A base for the neighbourhoods of $G$ can be given by

$$
\begin{gathered}
I_{0}(x):=G \\
I_{n}(x):=I_{n}\left(x_{0}, \ldots, x_{n-1}\right):=\left\{y \in G: y=\left(x_{0}, \ldots, x_{n-1}, y_{n}, y_{n+1}, \ldots\right)\right\}
\end{gathered}
$$

$(x \in G, n \in \mathbb{N}), I_{n}(x)$ are called dyadic intervals. Let $0=(0: i \in \mathbb{N}) \in G$ denote the null element of $G$, and for the simplicity we write $I_{n}:=I_{n}(0)(n \in \mathbb{N})$. Set $e_{n}:=$ $(0, \ldots, 0,1,0, \ldots) \in G$, the $n$th component of which is 1 and the rest are zeros.

[^0]The Fine's map $\|\cdot\|: G \rightarrow[0,1]$ is defined by

$$
\begin{equation*}
\|x\|:=\sum_{n=0}^{\infty} x_{n} 2^{-(n+1)} \tag{1}
\end{equation*}
$$

Backwards, each $x \in[0,1[$ can be expressed in number system based 2 in the form

$$
x=\sum_{j=0}^{\infty} x_{j} 2^{-j-1}, \quad \text { where } x_{j} \in\{0,1\} \text { for all } j
$$

This expansion is unique except for dyadic rational numbers $x \in\left\{\frac{p}{2^{n}}: p, n \in \mathbb{P}\right\}$. In this case we choose the expansion which terminates in 0 's. For $n \in \mathbb{N}$ let $n_{k}$ be the $k$ th coordinate of $n$ with respect to number system based 2 . That is, we write

$$
n=\sum_{k=0}^{\infty} n_{k} 2^{k}
$$

where $n_{k} \in\{0,1\} \quad k \in \mathbb{N}$. We use the notation $|n|:=\max \left\{j \in \mathbb{N}: n_{j} \neq 0\right\}$, that is $2^{|n|} \leqslant n<2^{|n|+1}$, where $|n|$ is called the order of natural number $n$.

Let $r_{k}$ denote the $k$-th Rademacher function, it is defined by

$$
r_{k}(x):=(-1)^{x_{k}} \quad(k \in \mathbb{N}, x \in G)
$$

The Walsh-Paley system (simply we say Walsh system) is defined as the product system of Rademacher functions. Namely,

$$
w_{n}(x):=\prod_{k=0}^{\infty}\left(r_{k}(x)\right)^{n_{k}}=r_{|n|}(x)(-1)^{\mid n=0} \sum_{k=0}^{|n|-1} n_{k} x_{k} \quad(x \in G, n \in \mathbb{P})
$$

The Walsh-Kaczmarz functions are defined by $\kappa_{0}:=1$ and for $n \geqslant 1$

$$
\kappa_{n}(x):=r_{|n|}(x) \prod_{k=0}^{|n|-1}\left(r_{|n|-1-k}(x)\right)^{n_{k}}=r_{|n|}(x)(-1)^{\sum_{k=0}^{|n|-1} n_{k} x_{|n|-1-k}}
$$

The Walsh-Kaczmarz system was introduced by Šneider [26] in 1948. Some basic result with respect to Walsh-Kaczmarz system can be found in [3, 21, 23, 24, 25, 29], for current results see also [12, 16, 27, 28, 35]. We give more details later.

It is known that the set of Walsh-Kaczmarz functions and the set of Walsh-Paley functions are equal in dyadic blocks. Namely,

$$
\left\{\kappa_{n}: 2^{k} \leqslant n<2^{k+1}\right\}=\left\{w_{n}: 2^{k} \leqslant n<2^{k+1}\right\}
$$

for all $k \in \mathbb{P}$. Moreover, $\kappa_{0}=w_{0}$.
The relation between the Walsh-Paley and Walsh-Kaczmarz system is not a simple relation, it is given by a coordinate transformation (for more details see [25] written by V.A. Skvortsov).

For both system we define the one-dimensional Dirichlet kernels and Cesàro kernels (see [10, 24, 35, 36]) by

$$
D_{n}^{\psi}:=\sum_{k=0}^{n-1} \psi_{k}, \quad K_{n}^{\psi, \alpha}(x):=\frac{1}{A_{n}^{\alpha}} \sum_{k=1}^{n} A_{n-k}^{\alpha-1} D_{k}^{\psi}(x)
$$

$\left(\psi_{n}=w_{n}(n \in \mathbb{N})\right.$ or $\psi_{n}=\kappa_{n}(n \in \mathbb{N})$, and $\left.0<\alpha\right)$, where

$$
A_{j}^{\alpha}:=\binom{j+\alpha}{j}=\frac{(\alpha+1)(\alpha+2) \ldots(\alpha+j)}{j!} \quad(j \in \mathbb{N} ; \alpha \neq-1,-2, \ldots)
$$

Choosing $\alpha=1$ we get back the Fejér kernels

$$
K_{n}^{\psi}=K_{n}^{\psi, 1}=\frac{1}{n} \sum_{k=1}^{n} D_{k}^{\psi}=\sum_{k=0}^{n-1}\left(1-\frac{k}{n}\right) \psi_{k} .
$$

These functions has some good properties, useful in the following investigations. First, we mention a simple result with respect to the Dirichlet kernels, which play a central role in the Walsh-Fourier analysis (see [20]):

$$
D_{2^{n}}^{w}(x)=D_{2^{n}}^{\kappa}(x)=D_{2^{n}}(x)= \begin{cases}0, & \text { if } x \notin I_{n}  \tag{2}\\ 2^{n}, & \text { if } x \in I_{n}\end{cases}
$$

For functions $f, g \in L^{1}(G)$ the dyadic convolution $f * g$ is defined by

$$
(f * g)(x):=\int_{G} f(t) g(x+t) d \mu(t)
$$

For $f \in L^{1}$ let us denote the $n$th partial sum of $f$ by $S_{n}^{\psi}(f)$, the $n$th Fejér means of $f$ by $\sigma_{n}^{\psi}(f)$ and the $n$th $(C, \alpha)$ mean of $f$ by $\sigma_{n}^{\psi, \alpha}(f)$. It is clear that $S_{n}^{\psi}(f)=$ $f * D_{n}^{\psi}, \sigma_{n}^{\psi}(f)=f * K_{n}^{\psi}$ and $\sigma_{n}^{\psi, \alpha}(f)=f * K_{n}^{\psi, \alpha}(n \in \mathbb{N})$. We remark that the Fejér kernels and $(C, \alpha)$ kernels $(0<\alpha)$ are uniformly bounded for both systems, that is,

$$
\sup _{n}\left\|K_{n}^{w, \alpha}\right\|_{1}<\infty \quad \text { and } \quad \sup _{n}\left\|K_{n}^{\kappa, \alpha}\right\|_{1}<\infty
$$

hold (see [3, 20, 22, 23, 24, 25]).
Further, we assume that the summation kernels $P_{n}:=\sum_{k=0}^{n-1} \lambda_{n, k} \psi_{k}$ with real coefficients $\lambda_{n, k}(n, k \in \mathbb{N})$ (here $\left\{\psi_{k}: k \in \mathbb{N}\right\}$ denotes the Walsh-Paley or the WalshKaczmarz system) satisfy the inequality

$$
\begin{equation*}
\sup _{n}\left\|P_{n}\right\|_{1}<\infty . \tag{3}
\end{equation*}
$$

If we consider the maximal operator

$$
\begin{equation*}
T(f):=\sup _{n}\left|f * P_{n}\right| \quad\left(f \in L^{1}(G)\right) \tag{4}
\end{equation*}
$$

then $T: L^{\infty}(G) \rightarrow L^{\infty}(G)$ is evidently bounded.
The $\sigma$-algebra generated by the dyadic intervals $I_{n}(x)(x \in G)$ will be denoted by $\mathbf{F}_{n}(n \in \mathbb{N})$. Let us denote by $f=\left(f^{(n)}, n \in \mathbb{N}\right)$ a martingale with respect to $\left(\mathbf{F}_{n}, n \in \mathbb{N}\right)$ (for details see, e. g. [31]). The maximal function of a martingale $f$ is defined by

$$
f^{*}=\sup _{n \in \mathbb{N}}\left|f^{(n)}\right|
$$

In case $f \in L^{1}$, the maximal function can also be given by

$$
f^{*}(x)=\sup _{n \in \mathbb{N}} \frac{1}{\mu\left(I_{n}(x)\right)}\left|\int_{I_{n}(x)} f(u) d \mu(u)\right|, \quad x \in G
$$

For $0<p \leqslant \infty$ the Hardy martingale space $H_{p}$ consists of all martingales for which

$$
\|f\|_{H_{p}}:=\left\|f^{*}\right\|_{p}<\infty
$$

If $f \in L^{1}$, then it is easily seen that the sequence $\left(S_{2^{n}}(f): n \in \mathbb{N}\right)$ is a martingale. If $f$ is a martingale, that is, $f=\left(f^{(0)}, f^{(1)}, \ldots\right)$ then the Fourier coefficients with respect to both systems Walsh-Paley and Walsh-Kaczmarz, must be defined in a little bit different way:

$$
\widehat{f}^{\psi}(i)=\lim _{k \rightarrow \infty} \int_{G} f^{(k)}(x) \psi_{i}(x) d \mu(x)
$$

The Fourier coefficients of $f \in L^{1}$ are the same as the ones of the martingale $\left(S_{2^{n}}(f)\right.$ : $n \in \mathbb{N}$ ) obtained from $f$. The atomic decomposition is a useful characterization of the Hardy space $H_{p}$. Let $0<p \leqslant 1$. A bounded measurable function $a$ is a $p$-atom, if either $a$ is identically equal to 1 , or there exists a dyadic interval $I$ for which
a) $\operatorname{supp} a \subseteq I$,
b) $\|a\|_{\infty} \leqslant \mu(I)^{-1 / p}$,
c) $\int_{I} a d \mu=0$.

We say that the atom $a$ is supported on the dyadic interval $I$. Then a martingale $f=$ $\left(f_{n}: n \in \mathbb{N}\right)$ belongs to the Hardy space $H_{p}(0<p \leqslant 1)$ if and only if there exists a sequence $\left(\lambda_{k}: k \in \mathbb{N}\right)$ of real numbers such that $\sum_{k=0}^{\infty}\left|\lambda_{k}\right|^{p}<\infty$ and

$$
\begin{equation*}
f=\sum_{k=0}^{\infty} \lambda_{k} a_{k} \tag{5}
\end{equation*}
$$

Moreover, the following equivalence of norms (quasi-norms) holds:

$$
c_{p}\|f\|_{H_{p}} \leqslant \inf \left(\sum_{k=0}^{\infty}\left|\lambda_{k}\right|^{p}\right)^{1 / p} \leqslant C_{p}\|f\|_{H_{p}} \quad\left(f \in H_{p}\right)
$$

where the infimum is taken over all decompositions of $f$ of the form (5). We note that here and later $c_{p}$ and $C_{p}$, denote positive constants depending only on $p$, although not always the same in different occurrences.

In this paper we use the next Lemma of Weisz [31]:

Lemma 1. (Weisz [31]) Suppose that the operator $T$ is $\sigma$-sublinear and $p$-quasilocal for some $0<p<1$. If $T$ is bounded from $L^{\infty}$ to $L^{\infty}$, then

$$
\|T f\|_{p} \leqslant c_{p}\|f\|_{H_{p}} \quad \text { for all } f \in H_{p}
$$

The Kronecker product $\left(\psi_{n}: n \in \mathbb{N}^{d}\right)$ of $d$ Walsh-(Kaczmarz) system is said to be the $d$-dimensional Walsh-(Kaczmarz) system. That is,

$$
\psi_{n}(x)=\psi_{n_{1}}\left(x^{1}\right) \ldots \psi_{n_{d}}\left(x^{d}\right)
$$

where $n=\left(n_{1}, \ldots, n_{d}\right)$ and $x=\left(x^{1}, \ldots, x^{d}\right)$.
If $f \in L^{1}\left(G^{d}\right)$, then the number $\widehat{f}^{\psi}(n):=\int_{G^{d}} f \psi_{n} \quad\left(n \in \mathbb{N}^{d}\right)$ is said to be the $n$th Walsh-(Kaczmarz)-Fourier coefficient of $f$. We can extend this definition to martingales in the usual way (see Weisz [30, 31]).

For $x=\left(x^{1}, \ldots, x^{d}\right) \in G^{d}$ and $n=\left(n_{1}, \ldots, n_{d}\right) \in \mathbb{N}^{d}$ the $d$-dimensional rectangles are defined by $I_{n}(x):=I_{n_{1}}\left(x^{1}\right) \times \ldots \times I_{n_{d}}\left(x^{d}\right)$. For $n \in \mathbb{N}^{d}$ the $\sigma$-algebra generated by the rectangles $\left\{I_{n}(x), x \in G^{d}\right\}$ is denoted by $\mathscr{F}_{n}$. The conditional expectation operators relative to $\mathscr{F}_{n}$ are denoted by $E_{n}$ (with $n=\left(n_{1}, \ldots, n_{d}\right)$ ).

Suppose that the functions $\gamma_{j}:[1,+\infty) \rightarrow[1,+\infty)$ are strictly monotone increasing continuous functions with properties $\lim _{x \rightarrow+\infty} \gamma_{j}(x)=+\infty$ and $\gamma_{j}(1)=1$ for all $j=$ $2, \ldots, d$. Moreover, suppose that for $j=2, \ldots, d$ there exist $\zeta, c_{j, 1}, c_{j, 2}>1$ such that the inequality

$$
\begin{equation*}
c_{j, 1} \gamma_{j}(x) \leqslant \gamma_{j}(\zeta x) \leqslant c_{j, 2} \gamma_{j}(x) \tag{6}
\end{equation*}
$$

holds for each $x \geqslant 1$. In this case the functions $\gamma_{j}$ are called CRF (cone-like restriction functions) [4, 8]. Let us introduce the notion $\gamma:=\left(\gamma_{2}, \ldots, \gamma_{d}\right)$ and set $\beta_{j} \geqslant 1$ be fixed for $j=2, \ldots, d$. We define the $d$-dimensional cone-like set $L$ (with respect to the first dimension) by

$$
L:=\left\{n \in \mathbb{N}^{d}: \gamma_{j}\left(n_{1}\right) / \beta_{j} \leqslant n_{j} \leqslant \beta_{j} \gamma_{j}\left(n_{1}\right), j=2, \ldots, d\right\}
$$

The $d$-dimensional Dirichlet kernels, Fejér kernels and ( $C, \alpha$ ) kernels can be given as the Kronecker product of one-dimensional kernels. That is,

$$
\begin{gathered}
D_{n}^{\psi}(x)=D_{n_{1}}^{\psi}\left(x^{1}\right) \ldots D_{n_{d}}^{\psi}\left(x^{d}\right), \quad K_{n}^{\psi}(x)=K_{n_{1}}^{\psi}\left(x^{1}\right) \ldots K_{n_{d}}^{\psi}\left(x^{d}\right), \\
K_{n}^{\psi, \alpha}(x)=K_{n_{1}}^{\psi, \alpha_{1}}\left(x^{1}\right) \ldots K_{n_{d}}^{\psi, \alpha_{d}}\left(x^{d}\right)
\end{gathered}
$$

where $n=\left(n_{1}, \ldots, n_{d}\right) \in \mathbb{N}^{d}, x=\left(x^{1}, \ldots, x^{d}\right) \in G^{d}$, and $\alpha=\left(\alpha_{1}, \ldots, \alpha_{d}\right)$.
In the present paper we investigate the boundedness of the maximal operator of some $d$-dimensional means, provided that the set of the indeces is inside a cone-like set $L$ and the convergence over a cone-like set $L$. Namely, we consider the $d$-parameter analogue of $T$ (see (4)) as follows. Let $P_{n}=P_{n_{1}, \ldots, n_{d}}\left(n=\left(n_{1}, \ldots, n_{d}\right) \in \mathbb{N}^{d}\right)$ be the Kronecker product of summation kernels $P_{n_{1}}^{1}, \ldots, P_{n_{d}}^{d}$, that is $P_{n_{1}, \ldots, n_{d}}\left(x^{1}, \ldots, x^{d}\right):=$ $P_{n_{1}}^{1}\left(x^{1}\right) \ldots P_{n_{d}}^{d}\left(x^{d}\right)$. For a fixed CRF function $\gamma$ we define the cone-like restricted maximal operator $T_{C L R}^{\gamma}$ by

$$
\begin{equation*}
T_{C L R}^{\gamma}(f):=\sup _{n \in L}\left|f * P_{n}\right| \quad\left(f \in L^{1}\left(G^{d}\right)\right) \tag{7}
\end{equation*}
$$

If $\gamma$ is the identical function (in each coordinate) then we get a $d$-dimensional cone. The cone-like sets were introduced by Gát in [4]. The condition (6) on the function $\gamma$ is natural, because Gát [4] proved that to each cone-like set with respect to the first
dimension there exists a larger cone-like set with respect to the second dimension and reversely, if and only if the inequality (6) holds.

Connecting to the work of Gát [4] Weisz defined a new type martingale Hardy space depending on the function $\gamma$ (see [32]). For a given $n_{1} \in \mathbb{N}$ let us set $n_{j}:=$ $\left|\gamma_{j}\left(2^{n_{1}}\right)\right|$, that is, $n_{j}$ is the order of $\gamma_{j}\left(2^{n_{1}}\right)$ (this means that $\left.2^{n_{j}} \leqslant \gamma_{j}\left(2^{n_{1}}\right)<2^{n_{j}+1}\right)$ for $j=2, \ldots, d$. Let us set $\bar{n}_{1}:=\left(n_{1}, \ldots, n_{d}\right)$. Since, the function $\gamma$ is increasing (for each coordinate function $\gamma_{j}$ ), the sequence ( $\bar{n}_{1}, n_{1} \in \mathbb{N}$ ) is increasing (in each coordinate), as well. A class of one-parameter martingales $f=\left(f_{\bar{n}_{1}}, n_{1} \in \mathbb{N}\right)$ is given with respect to the $\sigma$-algebras $\left(\mathscr{F}_{\overline{n_{1}}}, n_{1} \in \mathbb{N}\right)$. The maximal function of a martingale $f$ is defined by $f^{*}=\sup _{n_{1} \in \mathbb{N}}\left|f_{\bar{n}_{1}}\right|$. For $0<p \leqslant \infty$ the martingale Hardy space $H_{p}^{\gamma}\left(G^{d}\right)$ consists of all martingales for which the $L_{p}$-norm of the maximal function $f^{*}$ is finite, that is, $\|f\|_{H_{p}^{\gamma}}:=\left\|f^{*}\right\|_{p}<\infty$. It is known that $H_{p}^{\gamma} \sim L^{p}$ for $1<p \leqslant \infty$, where $\sim$ denotes the equivalence of the norms and spaces (see [31]).

If $f \in L^{1}\left(G^{d}\right)$ then it is easily shown that the sequence $\left(S_{2^{n_{1}}, \ldots, 2^{n_{d}}}(f): \overline{n_{1}}=\right.$ $\left.\left(n_{1}, \ldots, n_{d}\right), n_{1} \in \mathbb{N}\right)$ is a one-parameter martingale with respect to the $\sigma$-algebras $\left(\mathscr{F}_{\overline{n_{1}}}, n_{1} \in \mathbb{N}\right)$. In this case the maximal function can also be given by

$$
f^{*}(x)=\sup _{n_{1} \in \mathbb{N}} \frac{1}{\operatorname{mes}\left(I_{\overline{n_{1}}}(x)\right)}\left|\int_{I_{\overline{n_{1}}}(x)} f(u) d \mu(u)\right|=\sup _{n_{1} \in \mathbb{N}}\left|S_{2^{n_{1}}, \ldots, 2^{n_{d}}(f, x) \mid}\right|
$$

for $x \in G^{d}$. The Hardy space $H_{p}^{\gamma}$ has atomic structure also. The atoms $a$ are supported on the dyadic rectangles $I$ from the $\sigma$-algebras $\mathscr{F}_{\overline{n_{1}}}$.

If $\gamma$ is the identical function (in each coordinate) then $L$ is a $d$-dimensional cone. In the two-dimensional case, using a two-dimensional cone restriction set, the properties of the maximal operator of Walsh-Paley-Fejér means was discussed by Gát [5] and Weisz [33], separately. Later, the cone restricted maximal operator of two-dimensional Walsh-Kaczmarz-Fejér means was investigated by Simon [22]. Namely, they showed that the maximal operator $\sigma^{*}$ of Fejér means is bounded from the Hardy space $H_{p}$ to the Lebesgue space $L^{p}$ for $p>1 / 2$. It was shown also that the end point $p=1 / 2$ is essential. Further properties in the end point $p=\frac{1}{2}$ was discussed later. Namely, in 2007 Goginava and the first author proved that the cone-restricted maximal operator is not bounded from the Hardy space $H_{1 / 2}$ to the space weak- $L^{1 / 2}[13,14]$.

Connecting to the original paper [4] on trigonometric system, Gát asked the following. "What could we state for other systems for example Walsh-Paley, WalshKaczmarz and Vilenkin systems and for other means for example logarithmic means, Riesz means, (C, $\alpha$ ) means?" Some parts of Gát's question was answered by Weisz [32], Blahota and the authors [2, 17, 18, 19] and naturally in paper [8]. Moreover, there are some results in papers $[5,8]$ about divergence of two-dimensional Fejér means. Namely, if we suppose that $\beta$ is a function and not just a constant, then we have two cases. Either $\beta$ is bounded, then we have a.e. convergence for each integrable function, or $\beta$ is not bounded, then the maximal convergence space is $L \log ^{+} L$.

In 2011, the properties of the maximal operator of the $(C, \alpha)$ and Riesz means of multi-dimensional Vilenkin-Fourier series with cone-like restriction set, was discussed by Weisz [32]. Namely, it was shown that the maximal operator is bounded from dyadic

Hardy space $H_{p}$ to the space $L^{p}$ for $p_{0}<p \leqslant \infty\left(p_{0}:=\max \left\{1 /\left(1+\alpha_{k}\right): k=1, \ldots, d\right\}\right)$ and is of weak type $(1,1)$. Recently, it was shown that the index $p_{0}$ is sharp. Namely, it was proven that the maximal operator is not bounded from the dyadic Hardy space $H_{p_{0}}$ to the space $L^{p_{0}}$ [2] (see also [19]). A detailed list of the reached results for oneand several dimensional Walsh-like systems can be found in [34]. For Walsh-Kaczmarz system the properties of cone-like restricted two-dimensional maximal operator of Fejér and (C, $\alpha$ ) means was discussed in [17, 18]. Namely, it is proven that the maximal operator is bounded from dyadic Hardy space $H_{p}$ to the Lebesgue space $L_{p}$ for $p_{0}<$ $p \leqslant \infty$ (with the same $p_{0}$ as Weisz showed) and is of weak type $(1,1)$. Moreover, at the end point $p=p_{0}$, it is showed that the maximal operator $\sigma_{L}^{\kappa, \alpha, *}$ is not bounded from the Hardy space $H_{p_{0}}^{\gamma}$ to the space $L^{p_{0}}$.

Our work is motivated by the work of the authors [17,18] mentioned above and the paper of Simon [22]. In the last paper Simon improved a so called transference method. Namely, he considered the maximal operator $T$ of a sequence of summations (see (4)) and showed that the $p$-quasi-locality of $T$ implies the same statement for its two-dimensional version $T_{C R}^{I d}$ where $I d$ notes that the CRF function $\gamma$ is the identical function (in each coordinate, see (7), as well). The main aim of this paper is to extend the transference method of Simon for all CRF functions $\gamma$ and cone-like sets $L$ defined by $\gamma$. We mention that our proof is mainly based on the papers [18,22], we generalize the method improved in [18] and in the same time we extend the method of paper [22] for cone-like sets. Applying Lemma 1 of Weisz and some assumptions on the summation kernels $P_{n_{1}, \ldots, n_{d}}$ we state that the maximal operator $T_{C L R}^{\gamma}$ is bounded from the Hardy space $H_{p}^{\gamma}$ to the Lebesgue space $L^{p}$ for $p>p_{0}$. In the end point $p_{0}$ assuming some natural conditions on one-dimensional kernels we show that the maximal operator $T_{C L R}^{\gamma}$ is not bounded from the Hardy space $H_{p_{0}}^{\gamma}$ to the Lebesgue space $L^{p_{0}}$. After proving our main theorems we give an application, we get some new results unknown until the present days.

## 2. Improved transference method

In the one-dimensional case to apply the previous Lemma of Weisz (see Lemma 1) we show usually that the operator $T: L^{\infty}(G) \rightarrow L^{\infty}(G)$ in question is bounded. It immediately follows from inequality (3). The operator $T$ is called $p$-quasi-local if for arbitrary $p$-atom $a$ supported on the dyadic interval $I$ the inequality

$$
\begin{equation*}
\int_{G \backslash I}|T(a)|^{p} d \mu \leqslant C_{p} \tag{8}
\end{equation*}
$$

holds [30]. Hence, $p$-quasi-locality together with $\left(L^{\infty}, L^{\infty}\right)$-boundedness of the $\sigma$ sublinear operator $T$ implies that the operator $T$ is bounded from the Hardy space $H_{p}$ to the Lebesgue space $L^{p}$.

If $I=I_{n}(x)$ is a one-dimensional dyadic interval for some $x \in G$ and $n \in \mathbb{N}$, then for all $r=0,1, \ldots, n$ we define $I^{r}$ by $I^{r}:=I_{n-r}(x)$. Furthermore, the definition of $p$ -quasi-locality of operator $T$ can be modified as follows: there exists $r=0,1, \ldots$ such
that

$$
\begin{equation*}
\int_{G \backslash I^{r}}|T(a)|^{p} d \mu \leqslant C_{p} \tag{9}
\end{equation*}
$$

holds for every $p$-atom $a$ with support $I$. Analogical idea can be applied in multidimensional case. In most of the proofs of the $p$-quasi-locality we could realize the inequality

$$
\begin{equation*}
\int_{G \backslash I_{n}}\left(\sup _{n \geqslant 2^{n}} \int_{I_{n}}\left|P_{n}(x+t)\right| d \mu(t)\right)^{p} d \mu(x) \leqslant C_{p} 2^{-N} \quad(n \in \mathbb{N}) \tag{10}
\end{equation*}
$$

which immediately implies inequalities (8) or (9) (see [22]).
In the next Theorem, we prove that the inequality (10) for one-dimensional kernels $P_{n_{i}}^{i}:=\sum_{k=0}^{n_{i}-1} \lambda_{n_{i}, k}^{i} \psi_{k}(i=1, \ldots, d)$ (that is the $p$-quasi-locality of the one-dimensional operators $T^{1}, \ldots, T^{d}$ defined by the kernels $P_{n_{1}}^{1}, \ldots P_{n_{d}}^{d}$, respectively) implies the same property for cone-like restricted operator $T_{C L R}^{\gamma}$.

THEOREM 1. Let the function $\gamma$ be CRF. Assume that the kernels $P_{n_{i}}^{i}$ satisfy the inequalities (3) and (10) for all $0<p_{i}<p \leqslant 1(i=1, \ldots, d)$. Then the maximal operator $T_{C L R}^{\gamma}$ is bounded from the Hardy space $H_{p}^{\gamma}$ to the Lebesgue space $L^{p}$ for $p_{0}<p \leqslant 1$ (where $\left.p_{0}:=\max \left\{p_{1}, \ldots, p_{d}\right\}\right)$. Moreover, the maximal operator $T_{C L R}^{\gamma}$ is of weak type $(1,1)$.

Proof. The operator $T_{C L R}^{\gamma}$ is bounded from the space $L^{\infty}$ to the space $L^{\infty}$. It immediately follows from the inequality (3). Moreover, it can be seen easily that the operator $T_{C L R}^{\gamma}$ is $\sigma$-sublinear.

Let $a$ be a $p$-atom. Let it be supported on the dyadic rectangle $I$. Without loss of generality we can assume that $I=I_{N_{1}} \times \ldots \times I_{N_{d}}\left(\right.$ with $\left.N_{j}:=\left|\gamma_{j}\left(2^{N_{1}}\right)\right|, j=2, \ldots, d\right)$. The atom $a$ satisfies $\|a\|_{\infty} \leqslant 2^{\left(N_{1}+\ldots+N_{d}\right) / p}$ and $\int_{I} a d \mu=0$. Furthermore, it follows that $a * P_{n}=0\left(n=\left(n_{1}, \ldots, n_{d}\right)\right)$, when $n_{j}<2^{N_{j}}$ for $j=1, \ldots, d$.

In the next steps, we use the following inequality and the monotonicity of CRF functions $\gamma_{j}(j=2, \ldots, d)$.

$$
c_{j, 1}^{l} \gamma_{j}\left(\frac{2^{N_{1}}}{\zeta^{l}}\right) \leqslant \gamma_{j}\left(2^{N_{1}}\right)=\gamma_{j}\left(\frac{2^{N_{1}}}{\zeta^{l}} \zeta^{l}\right) \leqslant c_{j, 2}^{l} \gamma_{j}\left(\frac{2^{N_{1}}}{\zeta^{l}}\right)
$$

holds for all $l \in \mathbb{P}(j=2, \ldots, d)$. That is,

$$
\begin{equation*}
\frac{\gamma_{j}\left(2^{N_{1}}\right)}{c_{j, 2}^{l}} \leqslant \gamma_{j}\left(\frac{2^{N_{1}}}{\zeta^{l}}\right) \leqslant \frac{\gamma_{j}\left(2^{N_{1}}\right)}{c_{j, 1}^{l}} \quad(j=2, \ldots, d) \tag{11}
\end{equation*}
$$

First, we apply the right side of inequality (11) for any positive real number $x$,

$$
\gamma_{j}\left(\frac{2^{N_{1}}}{\zeta^{x}}\right) \leqslant \gamma_{j}\left(\frac{2^{N_{1}}}{\zeta[x]}\right) \leqslant \frac{\gamma_{j}\left(2^{N_{1}}\right)}{c_{j, 1}^{[x]}} \leqslant \frac{c_{j, 1} \gamma_{j}\left(2^{N_{1}}\right)}{c_{j, 1}^{x}} \quad(j=2, \ldots, d)
$$

where $[x]$ denotes the integer part of $x$. Now, let us set $\delta:=\max \left\{\zeta^{\log _{c_{j, 1}} 2 \beta_{j}+1}: j=\right.$ $2, \ldots, d\}$, as we did in paper [18].

If $n_{1} \leqslant 2^{N_{1}} / \delta$, then

$$
\begin{aligned}
n_{j} & \leqslant \beta_{j} \gamma_{j}\left(n_{1}\right) \leqslant \beta_{j} \gamma_{j}\left(2^{N_{1}} \zeta^{-\log _{c_{j, 1}} 2 \beta_{j}-1}\right) \\
& \leqslant \beta_{j} \frac{c_{j, 1}}{\log _{c_{j, 1}} 2 \beta_{j}+1}
\end{aligned} \gamma_{j}\left(2^{N_{1}}\right) \leqslant \frac{\gamma_{j}\left(2^{N_{1}}\right)}{2}<2^{N_{j}} .
$$

$\zeta, c_{j, 1}, c_{j, 2}>1, \beta_{j} \geqslant 1$ imply $n_{1}<2^{N_{1}}$ and $n_{j} \leqslant \gamma_{j}\left(2^{N_{1}}\right) / 2<2^{N_{j}}(j=2, \ldots, d)$. This yields $a * P_{n}=0$ for $n=\left(n_{1}, \ldots, n_{d}\right)$.

That is, we could suppose that $n_{1}>2^{N_{1}} / \delta$. Now, we apply the left side of inequality (11) for any positive real number $x$,

$$
\gamma_{j}\left(\frac{2^{N_{1}}}{\zeta^{x}}\right) \geqslant \gamma_{j}\left(\frac{2^{N_{1}}}{\zeta^{\lceil x\rceil}}\right) \geqslant \frac{\gamma_{j}\left(2^{N_{1}}\right)}{c_{j, 2}^{\lceil x\rceil}} \geqslant \frac{\gamma_{j}\left(2^{N_{1}}\right)}{c_{j, 2}^{x+1}} \quad(j=2, \ldots, d)
$$

where $\lceil x\rceil$ denotes the upper integer part of $x$. This yields that

$$
n_{j} \geqslant \frac{\gamma_{j}\left(n_{1}\right)}{\beta_{j}} \geqslant \frac{\gamma_{j}\left(2^{N_{1}} / \delta\right)}{\beta_{j}} \geqslant \frac{1}{\beta_{j} c_{j, 2}\left\{\log _{c_{j, 1}} 2 \beta_{j}+1: j=2, \ldots, d\right\}+1} \gamma_{j}\left(2^{N_{1}}\right) \geqslant \frac{\gamma_{j}\left(2^{N_{1}}\right)}{\delta_{j}^{\prime}} \geqslant \frac{2^{N_{j}}}{\delta^{\prime}}
$$

with $\delta^{\prime}:=\max _{j=2, \ldots, d} \delta_{j}^{\prime}$ for all $j=2, \ldots, d . \delta^{\prime}>1$ can be assumed.
The proof will be complete, if we show that the maximal operator $T_{C L R}^{\gamma}$ satisfies the modified version of $p$-quasi-locality (9) for $p_{0}<p \leqslant 1$, where $p_{0}:=\max \left\{p_{i}\right.$ : $i=1, \ldots, d\}$. For a $p$-atom $a$ in $H_{p}^{\gamma}$ with support $I=I_{N_{1}} \times \ldots \times I_{N_{d}}$ (with $\overline{N_{1}}=$ $\left(N_{1}, \ldots, N_{d}\right)$ ) the multidimensional version of inequality (9) reads as follows: There exist a constant $c_{p}>0$ and $r=0,1, \ldots$, such that

$$
\begin{equation*}
\int_{\bar{I}}\left|T_{C L R}^{\gamma}(a)\right|^{p} d \mu \leqslant c_{p}<\infty \tag{12}
\end{equation*}
$$

holds for each atom $a$, where $I^{r}:=I_{N_{1}}^{r} \times \ldots \times I_{N_{d}}^{r}:=I_{N_{1}-r} \times \ldots \times I_{N_{d}-r}\left(N_{j}-r \geqslant 0\right.$ for all $j=1, \ldots, d)$. We will give the value of $r$ later.

Let us set $x=\left(x^{1}, \ldots, x^{d}\right) \in \overline{I^{r}}$.

$$
\begin{aligned}
& \left|\left(a * P_{n}\right)(x)\right|=\left|\int_{I} a\left(t^{1}, \ldots, t^{d}\right) P_{n_{1}}^{1}\left(x^{1}+t^{1}\right) \ldots P_{n_{d}}^{d}\left(x^{d}+t^{d}\right) d \mu(t)\right| \\
& \leqslant 2^{\left(N_{1}+\ldots+N_{d}\right) / p} \int_{I_{N_{1}}}\left|P_{n_{1}}^{1}\left(x^{1}+t^{1}\right)\right| d \mu\left(t^{1}\right) \ldots \int_{I_{N_{d}}}\left|P_{n_{d}}^{d}\left(x^{d}+t^{d}\right)\right| d \mu\left(t^{d}\right) \quad\left(n \in \mathbb{N}^{d}\right)
\end{aligned}
$$

Now, we decompose the set $\overline{I^{r}}=\overline{I_{N_{1}}^{r} \times \ldots \times I_{N_{d}}^{r}}$ as the following disjoint union (see [18] or for $d=2$ [22])

$$
\begin{align*}
\overline{I^{r}}= & \left(\overline{I_{N_{1}}^{r}} \times \ldots \times \overline{I_{N_{d}}^{r}}\right) \cup \\
& \cup\left(I_{N_{1}}^{r} \times \overline{I_{N_{2}}^{r}} \times \ldots \times \overline{I_{N_{d}}^{r}}\right) \cup \ldots \cup\left(\overline{I_{N_{1}}^{r}} \times \ldots \times \overline{I_{N_{d-1}}^{r}} \times I_{N_{d}}^{r}\right) \cup \\
& \vdots  \tag{13}\\
& \cup\left(\overline{I_{N_{1}}^{r}} \times I_{N_{2}}^{r} \times \ldots \times I_{N_{d}}^{r}\right) \cup \ldots \cup\left(I_{N_{1}}^{r} \times \ldots \times I_{N_{d-1}}^{r} \times \overline{I_{N_{d}}^{r}}\right) .
\end{align*}
$$

Let us set $\delta^{\prime \prime}:=\max \left\{\delta, \delta^{\prime}\right\}$ and set $r \in \mathbb{P}$ such that $2^{-r} \leqslant 1 / \delta^{\prime \prime}<2^{-r+1}$. Moreover,
we set $L^{r, l}:=I_{N_{1}}^{r} \times \ldots \times I_{N_{l}}^{r} \times \overline{I_{N_{l+1}}^{r}} \times \ldots \times \overline{I_{N_{d}}^{r}}$ for $l=0, \ldots, d-1$. Let us define

$$
\begin{gathered}
J_{i}:=\int_{I_{N_{i}}^{r}}\left(\sup _{n_{i} \geqslant 2^{N_{i}} / \delta^{\prime \prime}} \int_{I_{N_{i}}}\left|P_{n_{i}}^{i}\left(x^{i}+t^{i}\right)\right| d \mu\left(t^{i}\right)\right)^{p} \mu\left(x^{i}\right), \quad i=1, \ldots, l, \\
\overline{J_{j}}:=\int_{\overline{I_{N_{j}}}}\left(\sup _{n_{j} \geqslant 2^{N_{j}} / \delta^{\prime \prime}} \int_{I_{N_{j}}}\left|P_{n_{j}}^{j}\left(x^{j}+t^{j}\right)\right| d \mu\left(t^{j}\right)\right)^{p} \mu\left(x^{j}\right), \quad j=l+1, \ldots, d .
\end{gathered}
$$

We immediately get

$$
\begin{equation*}
\int_{L^{, r l}}\left|T_{C L R}^{\gamma}(a)\right|^{p} d \mu \leqslant 2^{N_{1}+\ldots+N_{d}} J_{1} \cdot \ldots \cdot J_{l} \cdot \overline{J_{l+1}} \cdot \ldots \cdot \overline{J_{d}} \tag{14}
\end{equation*}
$$

First, we discuss the integrals $J_{i}(i=1, \ldots, l)$. Inequality (3) and the definitions of $\delta^{\prime \prime}, r$ immediately yield that

$$
\begin{equation*}
J_{i} \leqslant 2^{-\left(N_{i}-r\right)}\left(\sup _{n_{1} \in \mathbb{N}}\left\|P_{n_{i}}^{i}\right\|_{1}\right)^{p} \leqslant c_{p} 2^{-N_{i}} \quad(i=1, \ldots, l) \tag{15}
\end{equation*}
$$

Second, we discuss the integrals $\overline{J_{j}}(j=l+1, \ldots, d)$. Inequality (10) implies

$$
\begin{equation*}
\overline{J_{j}} \leqslant \int_{\overline{I_{N_{j}}}}\left(\sup _{n_{j} \geqslant 2^{N_{j}-r}} \int_{I_{N_{j}}^{r}}\left|P_{n_{j}}^{j}\left(x^{j}+t^{j}\right)\right| d \mu\left(t^{j}\right)\right)^{p} \mu\left(x^{j}\right) \leqslant c_{p} 2^{-N_{j}}, \text { if } p>p_{j} \tag{16}
\end{equation*}
$$

$(j=l+1, \ldots, d)$. Inequalities (14)-(16) yield

$$
\int_{L^{r, l}}\left|T_{C L R}^{\gamma}(a)\right|^{p} d \mu \leqslant c_{p} \quad \text { if } p>p_{0}
$$

for all $l=0, \ldots, d-1$. The decomposition (13) of $\overline{I^{r}}$ written above gives

$$
\int_{\bar{I}^{r}}\left|T_{C L R}^{\gamma}(a)\right|^{p} d \mu \leqslant c_{p} 2^{d} \quad \text { if } p>p_{0}
$$

Lemma 1 completes the proof of Theorem 1. The maximal operator $T_{C L R}^{\gamma}$ is of weak type $(1,1)$ follows by Marcinkiewicz interpolation theorem.

It follows that the a.e. convergence holds for polynomials, they form a dense subset of $L^{1}$. Thus, by standard argument the next Corollary holds.

Corollary 1. Let $\gamma$ be CRF and $L$ be a cone-like set. Let $f \in L^{1}\left(G^{d}\right)$. Assume that the kernels $P_{n_{i}}^{i}$ satisfy the inequalities (3) and (10) for all $0<p_{i}<p \leqslant 1$ ( $i=$ $1, \ldots, d)$. Then

$$
\lim _{\substack{\wedge n \rightarrow \infty \\ n \in L}} T_{n}(f)=f
$$

holds almost everywhere.
Fejér and $(C, \alpha)$ means were investigated with respect to Walsh-Paley sysetem in papers [8, 32] and Walsh-Kaczmarz system in [18].

In many cases the following question arises. Is the border point $p_{0}$ sharp or not? In the next Theorem, we prove that if we could construct a one-dimensional counterex-
ample martingale sequence in $H_{p_{0}}(G)$ which shows that the one-dimensional maximal operator $T$ is not bounded from the Hardy space $H_{p_{0}}$ to the space $L^{p_{0}}$. Then this enable us to construct a new $d$-dimensional counterexample martingale sequence in $H_{p_{0}}^{\gamma}$, as well, which has the same property from the Hardy space $H_{p_{0}}^{\gamma}$ to the space $L^{p_{0}}$.

In many papers the one-dimensional martingale

$$
\begin{equation*}
f_{A}(x):=D_{2^{A+1}}(x)-D_{2^{A}}(x) \tag{17}
\end{equation*}
$$

is applied as a counterexample martingale [9, 22]. Since,

$$
\hat{f}_{A}^{\psi}(k)= \begin{cases}1, & \text { if } k=2^{A}, \ldots, 2^{A+1}-1 \\ 0, & \text { otherwise }\end{cases}
$$

and

$$
S_{j}^{\psi}\left(f_{A}\right)(x)= \begin{cases}D_{j}^{\psi}(x)-D_{2^{A}}(x), & \text { if } j=2^{A}+1, \ldots, 2^{A+1}-1 \\ f_{A}(x), & \text { if } j \geqslant 2^{A+1} \\ 0, & \text { otherwise }\end{cases}
$$

It could be concluded that

$$
\left\|f_{A}\right\|_{H_{p}}=\left\|f_{A}^{*}\right\|_{p}=\left\|D_{2^{A}}\right\|_{p}=2^{A(1-1 / p)}
$$

and $f \in H_{p}$. Moreover, $f_{A}$ is a $p$-atom and satisfies

$$
\begin{equation*}
S_{2^{A+1}}\left(f_{A}\right)=f_{A} \tag{18}
\end{equation*}
$$

The next inequality usually is proved

$$
\begin{equation*}
\frac{\left\|T\left(f_{n}\right)\right\|_{p_{0}}}{\left\|f_{n}\right\|_{H_{p_{0}}}} \geqslant a_{n} \tag{19}
\end{equation*}
$$

where $T$ is defined by one-dimensional summation kernels $P_{n}$ (see (4)) and $\left(a_{n}\right)$ is a positive real valued sequence, which tends to $+\infty$ monotone increasingly. Inequality (19) and $\lim _{n \rightarrow+\infty} a_{n}=+\infty$ yield that the maximal operator $T$ is not bounded from the Hardy space $H_{p_{0}}$ to the space $L^{p_{0}}$. In the end point $p_{0}$ assuming some natural conditions for one-dimensional kernels in the next Theorem, we show that the maximal operator $T_{C L R}^{\gamma}$ is not bounded from the Hardy space $H_{p_{0}}^{\gamma}$ to the Lebesgue space $L^{p_{0}}$. Although, we have to require conditions not only for the kernels $P_{n_{1}}^{1}$, but for kernels $P_{n_{l}}^{l}(l=2, \ldots, d)$, as well. Namely, there exist positive constants $c_{l}^{*}$ such that

$$
\begin{equation*}
\left|P_{n_{l}}^{l} * \psi_{2^{k}-1}\right|=\left|\lambda_{n_{l}, 2^{k}-1}^{l}\right| \geqslant c_{l}^{*} \tag{20}
\end{equation*}
$$

hold for all $n_{l} \geqslant 2^{k+1}(l=2, \ldots, d)$. Inequality (20) hold automatically, if the kernel functions $P_{n_{l}}^{l}$ are the Fejér kernels $K_{n_{l}}^{\psi}$ or the $(C, \alpha)$ kernels $K_{n_{l}}^{\psi, \alpha}$. In the first case $c_{l}^{*}=\frac{1}{2}$ and in the second case it is easily seen that there exist positive constants $c_{l}^{*}$ such that

$$
\frac{1}{A_{n_{l}}^{\alpha}} \sum_{i=2^{k}}^{n_{l}-1} A_{n_{l}-i}^{\alpha-1} \geqslant c_{l}^{*}>0 \text { hold for all } n_{l} \geqslant 2^{k+1}
$$

So, inequality (20) is a natural condition for many type of kernels $P_{n_{l}}^{l}(l=2, \ldots, d)$.

THEOREM 2. Let $\gamma$ be CRF and $p_{0}=p_{1} \geqslant p_{i}(i=1, \ldots, d)$. Let $f_{A}$ be a $p_{0}$-atom in $H_{p_{0}}(G)$ with support $I_{A}$ which satisfies inequalities (18) and (19) with a positive real valued sequence $\left(a_{n}\right)$ which tends to $+\infty$ monotone increasingly. Moreover, the kernels $P_{n_{l}}^{l}$ satisfy inequality (20) for $l=2, \ldots, d$.

Then the maximal operator $T_{C L R}^{\gamma}$ is not bounded from the Hardy space $H_{p_{0}}^{\gamma}\left(G^{d}\right)$ to the space $L^{p_{0}}\left(G^{d}\right)$.

We note that the fact that $f_{A}$ is a $p_{0}$-atom in $H_{p_{0}}(G)$ with support $I_{A}$ and satisfies equality (18) determine $f_{A}$ having the form of (17) multiplied by a constant $c \neq 0$, where $|c| \leqslant 1$. This yields that

$$
\left\|f_{A}\right\|_{H_{p_{0}}}=\left\|f_{A}^{*}\right\|_{p_{0}}=\left\|f_{A}\right\|_{p_{0}}
$$

Proof. We note that some idea of this proof is coming from papers [11, 2]. Let $f_{A} \in H_{p_{0}}(G)$ such that it satisfies the conditions of our theorem. We define a $d$ dimensional martingale $F_{\overline{n_{1}}}$ in $H_{p_{0}}^{\gamma}\left(G^{d}\right)$ by

$$
F_{\overline{n_{1}}}(x):=f_{n_{1}}\left(x^{1}\right) \prod_{j=2}^{d} \psi_{2^{n_{j}-1}-1}\left(x^{j}\right)
$$

where $n_{2}, \ldots, n_{d}$ is defined to $n_{1}$, earlier, that is, $\overline{n_{1}}=\left(n_{1}, \ldots, n_{d}\right)$ and $x=\left(x^{1}, \ldots, x^{d}\right) \in$ $G^{d}$.

We have

$$
\hat{F}_{n_{1}}^{\psi}(k)= \begin{cases}\hat{f}_{n_{1}}^{\psi}\left(k_{1}\right), & \text { if } k_{j}=2^{n_{j}-1}-1 \text { for all } j=2, \ldots, d \\ 0, & \text { otherwise }\end{cases}
$$

for $k=\left(k_{1}, \ldots, k_{d}\right)$. Now, we calculate $S_{j}^{\psi}\left(F_{\overline{n_{1}}}\right)$. Using equality (18) and the Fourier coefficients of $F_{\overline{n_{1}}}$ we may write
$S_{j}^{\psi}\left(F_{\overline{n_{1}}}, x\right)= \begin{cases}S_{j_{1}}^{\psi}\left(f_{n_{1}}, x^{1}\right) \prod_{l=2}^{d} \psi_{2^{n_{l}-1}-1}\left(x^{l}\right), & \text { if } j_{l} \geqslant 2^{n_{l}-1} \text { for all } l=2, \ldots, d ; j_{1}<2^{n_{1}+1} \\ F_{\overline{n_{1}}}(x), & \text { if } j_{1} \geqslant 2^{n_{1}+1} \text { and } j_{l} \geqslant 2^{n_{l}-1} \\ 0, & \text { for all } l=2, \ldots, d ; \\ 0, & \text { otherwise. }\end{cases}$
We immediately have that

$$
F_{\overline{n_{1}}}^{*}(x)=\sup _{m_{1} \in \mathbb{N}}\left|S_{2^{m_{1}}, \ldots, 2^{m_{d}}}\left(F_{\overline{n_{1}}}, x\right)\right|=\left|F_{\overline{n_{1}}}(x)\right|=\left|f_{n_{1}}\left(x^{1}\right)\right|
$$

where $\overline{m_{1}}=\left(m_{1}, \ldots, m_{d}\right)$. Moreover,

$$
\begin{equation*}
\left\|F_{\overline{n_{1}}}\right\|_{H_{p_{0}}^{\gamma}\left(G^{d}\right)}=\left\|F_{\overline{n_{1}}}^{*}\right\|_{p_{0}}=\left\|f_{n_{1}}^{*}\right\|_{p_{0}}=\left\|f_{n_{1}}\right\|_{H_{p_{0}}(G)}<\infty . \tag{22}
\end{equation*}
$$

That is, $F_{\overline{n_{1}}} \in H_{p_{0}}^{\gamma}\left(G^{d}\right)$.
First, we set $L_{1}^{N}:=2^{n_{1}}+N$, where $0<N$ and $L_{j}^{N}:=\left[\gamma_{j}\left(2^{n_{1}}+N\right)\right]$ for $j=2, \ldots, d$ (where $[x]$ denotes the integer part of $x$ ). In this case $L^{N}:=\left(L_{1}^{N}, \ldots, L_{d}^{N}\right) \in L$. Let us
calculate $F_{\overline{n_{1}}} * P_{L^{N}}$.

$$
\left(F_{\overline{n_{1}}} * P_{L^{N}}\right)(x)=\left(f_{n_{1}} * P_{L_{1}^{N}}^{1}\right)\left(x^{1}\right) \prod_{l=2}^{d} \lambda_{L_{l}^{N}, 2^{n_{l}-1}-1}^{l} \psi_{2^{n_{l}-1}-1}\left(x^{l}\right)
$$

From this and inequality (20) we write

$$
\begin{align*}
\left|\left(F_{\overline{n_{1}}} * P_{L^{N}}\right)(x)\right| & =\left|\left(f_{n_{1}} * P_{L_{1}^{N}}^{1}\right)\left(x^{1}\right) \prod_{l=2}^{d} \lambda_{L_{l}^{N}, 2^{n_{l}-1}-1}^{l}\right| \\
& \geqslant\left|\left(f_{n_{1}} * P_{L_{1}^{N}}^{1}\right)\left(x^{1}\right)\right| \prod_{l=2}^{d} c_{l}^{*} \geqslant c^{*}\left|\left(f_{n_{1}} * P_{L_{1}^{N}}^{1}\right)\left(x^{1}\right)\right| \tag{23}
\end{align*}
$$

with a positive constant $c^{*}$. For the maximal operator $T_{C L R}^{\gamma}$ we get

$$
\begin{equation*}
T_{C L R}^{\gamma} F_{\overline{n_{1}}}(x)=\sup _{n \in L}\left|\left(F_{\overline{n_{1}}} * P_{n}\right)(x)\right| \geqslant \sup _{L_{1}^{N}}\left|\left(F_{\overline{n_{1}}} * P_{L^{N}}\right)(x)\right| \geqslant c^{*} \sup _{0<N}\left|\left(f_{n_{1}} * P_{L_{1}^{N}}^{1}\right)\left(x^{1}\right)\right| \tag{24}
\end{equation*}
$$

Since, $f_{n_{1}} \in H_{p_{0}}(G)$ is a $p_{0}$-atom with support $I_{n_{1}}$, we have that

$$
f_{n_{1}} * P_{m}=0 \quad \text { for } m \leqslant 2^{n_{1}}
$$

This together with inequality (24) yield that

$$
\begin{equation*}
T_{C L R}^{\gamma}\left(F_{\overline{n_{1}}}\right)(x) \geqslant c^{*} \sup _{m}\left|\left(f_{n_{1}} * P_{m}^{1}\right)\left(x^{1}\right)\right|=c^{*} T^{1}\left(f_{n_{1}}\right)\left(x^{1}\right) \tag{25}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{\left\|T_{C L R}^{\gamma}\left(F_{\overline{n_{1}}}\right)\right\|_{p_{0}}}{\left\|F_{\overline{n_{1}}}\right\|_{H_{p_{0}}^{\gamma}\left(G^{d}\right)}^{\gamma}} \geqslant \frac{c^{*}\left\|T^{1}\left(f_{n_{1}}\right)\right\|_{p_{0}}}{\left\|f_{n_{1}}\right\|_{H_{p_{0}}(G)}} \geqslant c^{*} a_{n_{1}} \tag{26}
\end{equation*}
$$

$n_{1} \rightarrow \infty$ and inequality (19) complete the proof of Theorem 2.
At last, we note that condition (20) can be weakened a little bit. Applying inequality (20) more precisely in the form

$$
\begin{equation*}
\left|P_{n_{l}}^{l} * \psi_{2^{k}-1}\right|=\left|\lambda_{n_{l}, 2^{k}-1}^{l}\right| \geqslant c_{l}^{k} \tag{27}
\end{equation*}
$$

in inequality (23). That is, instead of $\prod_{l=2}^{d} c_{l}^{*}$ we could write $\prod_{l=2}^{d} c_{l}^{n_{l}-1}$. Thus, on the right side of inequality (26) the expression $a_{n_{1}}^{*}:=\prod_{l=2}^{d} c_{l}^{n_{l}-1} a_{n_{1}}$ appears (where $n_{2}, \ldots, n_{d}$ defined to $n_{1}$ earlier). If $a_{n_{1}}^{*} \rightarrow+\infty$, while $n_{1} \rightarrow \infty$ our Theorem remain valid. It means that the absolute value of some coefficients in inequality (20) is not necessarily bounded from below by a positive constant.

## 3. Application

It is well-known (see Fine's map in the first section, equation (1)) that there is a direct connection between the Walsh group $G$ and the interval $I:=[0,1[$. The coordinate wise addition in $G$ defines a so-called dyadic addition denoted by $\dot{+}$ in the interval [ $0,1[$. The characters of the Walsh group $G$ are the Walsh-Paley functions. If we change the operation $\dot{+}$ on the interval $I:=[0,1[$ by the usual arithmetical sum denoted by + , then we get the so-called group of 2 -adic integers. It is denoted by $(I,+)$. The character system belongs to $(I,+)$ will change, as well. Namely, the 2 -adic (or arithmetic)
sum $a+b:=\sum_{n=0}^{\infty} r_{n} 2^{-(n+1)}(a, b \in I)$, where bits $q_{n}, r_{n} \in\{0,1\}(n \in \mathbb{N})$ are defined recursively as follows : $q_{-1}:=0, a_{n}+b_{n}+q_{n-1}=2 q_{n}+r_{n}$ for $n \in \mathbb{N}$. (Since $q_{n}, r_{n}$ take on only the values 0,1 , these equations uniquely determine the coefficients $q_{n}$ and $r_{n}$.) Set

$$
\varepsilon(t):=\exp (2 \pi l t) \quad(t \in \mathbb{R}), \quad\left(\imath=(-1)^{\frac{1}{2}}\right)
$$

and

$$
v_{2^{n}}(x):=\varepsilon\left(\frac{x_{n}}{2}+\ldots+\frac{x_{0}}{2^{n+1}}\right) \quad(x \in I, n \in \mathbb{N})
$$

We define the product system by

$$
v_{n}:=\prod_{n=0}^{\infty} v_{2 j}^{n_{j}},
$$

where $n_{j}$ is the $j$ th coordinate of natural number $n$, for more details see the notion of product system in the first section. It is known [15] that the system $\left(v_{n}, n \in \mathbb{N}\right)$ is the character system of $(I,+)$. In the present section, the system $\left\{\psi_{k}: k \in \mathbb{N}\right\}$ is defined by $\psi_{k}:=v_{k}$ for all $k \in \mathbb{N}$. The Fourier coefficients, the Dirichlet and the Fejér kernels and $(C, \alpha)$ kernels are defined in the same way as we did in the first part of this paper. The partial sums, Fejér means and $(C, \alpha)$ means of an integrable function $f$ are defined by the convolution of $f$ with the kernel functions. For system $\left(v_{n}, n \in \mathbb{N}\right)$ equality (2) remain valid. In the most cases the proving methods derived from the dyadic case.

In one-dimensional case the behaviour of the Cesàro means of the Fourier series on the group of 2-adic integers was discussed, the a.e. convergence and (H.L) issue were treated by Gát [6]. In paper [9] Gát and the first author investigated the maximal operator $\sigma^{*}:=\sup _{n}\left|\sigma_{n}^{1}\right|$ of the Fejér means with respect to the character system of 2 -adic integers. Among others, they proved that this operator is bounded from the Hardy space $H_{p}$ to the Lebesgue space $L^{p}$ if and only if $1 / 2<p<\infty$. They showed inequality (10) holds for the Fejér kernels for all $1 / 2<p<1$, as well (see inequality (4) in [9, page 75]). Inequality (3) for Fejér kernels follows from paper [7]. Applying our transference method (Theorem 1), we immediately get the following Theorem.

THEOREM 3. Let $\gamma$ be CRF. Set $\beta>1$. The maximal operator $\sigma_{C L R}^{*}$ is bounded from the Hardy space $H_{p}^{\gamma}\left(I^{d}\right)$ to the space $L^{p}\left(I^{d}\right)$ for all $\frac{1}{2}<p \leqslant 1$. Moreover, the maximal operator is of weak type $(1,1)$.

According to Corollary 1 we could state that the $d$-dimensional Fejér means $\sigma_{n}(f)$ of an integrable function $f \in L^{1}\left(I^{d}\right)$ converge almost everywhere to the function $f$ itself provided that the indeces are inside a cone-like set $L$.

In the end point case $p_{0}=\frac{1}{2}$ the one-dimensional martingale $f_{A}(x)=D_{2^{A+1}}(x)-$ $D_{2^{A}}(x)$ (see equality (17)) is applied to show that the maximal operator $\sigma^{*}$ of Fejér means is not bounded from Hardy space $H_{1 / 2}(I)$ to the Lebesque space $L_{1 / 2}(I)$ [9, Theorem 2.2]. During the proof of this result it was proven, that

$$
\frac{\left\|\sigma^{*}\left(f_{n}\right)\right\|_{1 / 2}}{\left\|f_{n}\right\|_{H_{1 / 2}}} \geqslant c n^{2}
$$

(see inequality (19), as well). For Fejér kernel functions

$$
K_{n}^{\psi} * \varphi_{2^{k}-1} \geqslant \frac{1}{2} \quad \text { hold for all } n \geqslant 2^{k+1}
$$

(see inequality (20), as well). Applying Theorem 2 we immediately have the following Theorem.

THEOREM 4. The cone-like restricted maximal operator $\sigma_{C L R}^{*}$ is not bounded from the Hardy space $H_{1 / 2}^{\gamma}\left(I^{d}\right)$ to the space $L^{1 / 2}\left(I^{d}\right)$.

## REFERENCES

[1] G. N. Agajev, N. Ya. Vilenkin, G. M. Dzhafarli and A. I. Rubinstein, Multiplicative systems of functions and harmonic analysis on O-dimensional groups, "ELM" Baku, USSR, 1981 (Russian).
[2] I. Blahota and K. Nagy, On the restricted summability of the multi-dimensional Vilenkin-Cesàro means, J. Math. Ineq., 11, (4) (2017), 997-1006.
[3] G. GÁt, On $(C, 1)$ summability of integrable functions with respect to the Walsh-Kaczmarz system, Studia Math., 130, (2) (1998), 135-148.
[4] G. GÁT, Pointwise convergence of cone-like restricted two-dimensional ( $C, 1$ ) means of trigonometric Fourier series, J. Approx. Theory, 149, (2007), 74-102.
[5] G. GÁt, Pointwise convergence of the Cesàro means of double Walsh series, Annales Univ. Sci. Budapest., Sect. Comp., 16, (1996), 173-184.
[6] G. GÁt, Almost everywhere convergence of Cesàro means of Fourier series on the group of 2-adic integers, Acta Math. Hungar., 116 (3), (2007), 209-221.
[7] G. GÁt, On the almost everywhere convergence of Fejér means of functions on the group of 2-adic integers, J. Approx. Theory, 90, (1997), 88-96.
[8] G. GÁt and K. Nagy, Pointwise convergence of cone-like restricted two-dimensional Fejér means of Walsh-Fourier series, Acta Math. Sinica, English Series, 26, (12) (2010), 2295-2304.
[9] G. GÁt And K. NAGY, On the maximal operators of Fejér means with respect to the character system of the group of 2-adic integers in Hardy spaces, Mathematical Notes, 98, (1-2) (2015), 68-77.
[10] U. Goginava, Maximal operators of $(C, \alpha)$-means of cubical partial sums of $d$-dimensional WalshFourier series, Analysis Math., 33, (2007), 263-286.
[11] U. Goginava, The maximal operator of the ( $C, \alpha$ ) means of the Walsh-Fourier series, Ann. Univ. Sci. Budapest. Sect. Comput., 26, (2006), 127-135.
[12] U. Goginava, The maximal operator of the Fejér means of the character system of the p-series field in the Kaczmarz rearrangement, Publ. Math. (Debrecen), 71, (1-2) (2007), 43-55.
[13] U. GoginaVa, Maximal operators of Fejér means of double Walsh-Fourier series, Acta Math. Hungar., 115, (4) (2007), 333-340.
[14] U. Goginava and K. Nagy, On the Fejér means of double Fourier series with respect to the WalshKaczmarz system, Periodica Math. Hungar., 55, (1) (2007), 39-45.
[15] E. Hewitt and K. Ross, Abstract Harmonic Analysis, Vol. I, Vol. II, Springer-Verlag, Heidelberg, 1963.
[16] N. MEMIC, Almost everywhere convergence of Fejér means of some subsequences of Fourier series for integrable functions with respect to the Kaczmarz system, Advances in Mathematics: Scientific Journal, 4, (1) (2015), 65-77.
[17] K. NAGY, On the restricted summability of Walsh-Kaczmarz-Fejér means, Georgian Math. J., 22, (1) (2015), 131-140.
[18] K. Nagy and M. Salim, Restricted summability of multi-dimensional Cesàro means of Walsh-Kaczmarz-Fourier series, Publ. Math. Debrecen, 94, (3-4) (2019), 381-394.
[19] K. NAGY, On the restricted summability of two-dimensional Walsh-Fejér means, Publ. Math. Debrecen, 85, (1-2) (2014), 113-122.
[20] F. Schipp, W. R. Wade, P. Simon, and J. PáL, Walsh Series. An Introduction to Dyadic Harmonic Analysis, Adam Hilger Bristol-New York, 1990.
[21] F. SCHIPP, Pointwise convergence of expansions with respect to certain product systems, Anal. Math., 2, (1976), 63-75.
[22] P. Simon, Cesàro summability with respect to two-parameter Walsh system, Monatsh. Math., 131, (2000), 321-334.
[23] P. Simon, On the Cesàro summability with respect to the Walsh-Kaczmarz system, J. Approx. Theory, 106, (2000), 249-261.
[24] P. Simon, ( $C, \alpha)$-summability of Walsh-Kaczmarz-Fourier series, J. Approx. Theory, 127, (2004), 39-60.
[25] V. A. Skvortsov, On Fourier series with respect to the Walsh-Kaczmarz system, Anal. Math., 7, (1981), 141-150.
[26] A. A. Šneider, On series with respect to the Walsh functions with monotone coefficients, Izv. Akad. Nauk SSSR Ser. Math., 12, (1948), 179-192.
[27] G. Tephnadze, On the maximal operators of Walsh-Kaczmarz-Fejér means, Period. Math. Hungar., 67, (1) (2013), 33-45.
[28] G. Tephnadze, Approximation by Walsh-Kaczmarz-Fejér means on the Hardy space, Acta Math. Sci. Ser. B Engl. Ed., 34, (5) (2014), 1593-1602.
[29] W. S. Young, On the a.e converence of Walsh-Kaczmarz-Fourier series, Proc. Amer. Math. Soc., 44, (1974), 353-358.
[30] F. Weisz, Martingale Hardy spaces and their applications in Fourier analysis, Springer-Verlang, Berlin, 1994.
[31] F. WeIsz, Summability of multi-dimensional Fourier series and Hardy space, Kluwer Academic, Dordrecht, 2002.
[32] F. Weisz, Restricted summability of multi-dimensional Vilenkin-Fourier series, Annales Univ. Sci. Budapest., Sect. Comp., 35, (2011), 305-317.
[33] F. Weisz, Cesàro summability of two-dimensional Walsh-Fourier series, Trans. Amer. Math. Soc., 348, (1996), 2169-2181.
[34] F. Weisz, Convergence of trigonometric and Walsh-Fourier series, Acta Math. Acad. Paed. Nyíregyháziensis, 32, (2) (2016), 277-301.
[35] X. Zhang and C. Zhang, $(C, \alpha)$-summability of character system of p-series field in Kaczmarz rearrangement, Acta Math. Sci. Ser. A, Chin. Ed., 29, (6) (2009), 1623-1633.
[36] A. ZyGMund, Trigonometric series, Cambridge Press, London, 3rd edition, 2002.
(Received May 22, 2020)
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[^0]:    Mathematics subject classification (2010): 42C10, 43A75, 42B08, 42B30.
    Keywords and phrases: Walsh-Paley system, Walsh-Kaczmarz system, maximal operator, multidimensional system, restricted summability, almost everywhere convergence, Cesàro means, group of 2-adic integers, Hardy space.

    Research supported by project UAEU UPAR 2017 Grant G00002599

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