# WHEN DOES THE MOORE-PENROSE INVERSE FLIP? 

R. E. Hartwig and P. Patrício

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#### Abstract

In this paper, we give necessary and sufficient conditions for the matrix $\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]$, over a *-regular ring, to have a Moore-Penrose inverse of four different types, corresponding to the four cases where the zero element can stand. In particular, we study the case where the MoorePenrose inverse of the matrix flips.


## 1. Introduction

Let $R$ be a regular *-ring with 1 , that is, for all $a \in R$ there exist $a^{-}$such that $a a^{-} a=a$, and with an involutory anti-isomorphism $(\cdot)^{*}$ on $R$, such that $\left(a^{*}\right)^{*}=$ $a,(a+b)^{*}=a^{*}+b^{*}$ and $(a b)^{*}=b^{*} a^{*}$.

It is well known [9, Lemma 4], that if the involution on $R$ satisfies the one term star-cancellation law

$$
\begin{equation*}
S C_{1}: a^{*} a=0 \Rightarrow a=0 \tag{1}
\end{equation*}
$$

then the Moore-Penrose inverse $a^{\dagger}$ can be defined. It is the unique solution to the four equations

$$
\begin{equation*}
\text { (i) } a x a=a, \text { (ii) } x a x=x, \text { (iii) }(a x)^{*}=a x, \text { (iv) }(x a)^{*}=x a \tag{2}
\end{equation*}
$$

We say $x$ is a 1-3 inverse of $a$ if it satisfies equations (i) and (iii) above, and $y$ is a 1-4 inverse of $a$ if it satisfies equations (i) and (iv) above. From the well known result due to Urquhart (cf. [1, page 48]), if $x$ and $y$ are a 1-3 and 1-4 inverse of $a$, respectively, then $a^{\dagger}=y a x$.

We note that regular rings that satisfy $S C_{1}$ are exactly those for which all of its elements are Moore-Penrose invertible. Such a ring is said to be a $*$-regular ring. We use $R_{2 \times 2}$ to denote the ring of $2 \times 2$ matrices over $R$.

A matrix $M=\left[\begin{array}{ll}a & c \\ b & d\end{array}\right]$ with coefficients in $R$ is said to be of $(i, j, 0)$ type if the $(i, j)$ entry $(M)_{i j}$ of $M$ is zero.

[^0]In this note we will be interested in the questions of when the matrix $\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]$ has a Moore-Penrose inverse of $(i, j, 0)$ type, for $i, j \in\{1,2\}$. In particular, we will address to the case when this inverse has the "flipped" form $\left[\begin{array}{l}* * \\ 0 *\end{array}\right]$. We will repeatedly use Cline's results ([3] and [4]) in order to express the Moore-Penrose inverse of a semiorthogonal sum and of a column matrix. The expressions derived are simpler when compared with [7].

We only consider the special involution on $R_{2 \times 2}$ of the form $\left[\begin{array}{ll}a & c \\ b & d\end{array}\right]^{*}=\left[\begin{array}{ll}a^{*} & b^{*} \\ c^{*} & d^{*}\end{array}\right]$.

## 2. Existence of the Moore-Penrose inverse

Consider the matrix $M=\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]$. In order to guarantee the existence and to be able to give a formula of $M^{\dagger}$, we assume the following extra conditions on the regular $R$ :

1. $S C_{2}: a^{*} a+b^{*} b=0 \Rightarrow a=0=b$ (two term star-cancellation)
2. for each $a \in R$, there is $c \in R$ such that $1+a^{*} a=c^{*} c=c c^{*}$ (square root axiom)

We note the following consequences:
(i) $1+a^{*} a$ is a unit for all $a \in R$, that is, $R$ has the symmetry property (see [2, page 9]). Indeed, if $R$ is regular and satisfies $S C_{2}$ then it also satisfies $S C_{1}$, which in turn implies all its elements are Moore-Penrose invertible. Let $u=1+a^{*} a$. If $u x=0$, then $x^{*} x+(a x)^{*}(a x)=0$ and hence, using condition $S C_{2}, x=0$. Thus $u$ is not a divisor of 0 . But $u\left(1-u^{\dagger} u\right)=0$ and hence $1-u^{\dagger} u=0$. Likewise $1-u u^{\dagger}=0$ and $u$ is a unit.
(ii) Since $1+a^{*} a=c c^{*}=c^{*} c$ is a unit, then the square root $c$ must be a unit as well.
(iii) $1+a^{*} a+b^{*} b=c^{*} c+b^{*} b=c^{*}\left[1+\left(b c^{-1}\right)^{*}\left(b c^{-1}\right)\right] c$, which is again a unit.
(iv) If $R$ satisfies $S C_{2}$ and is regular, then every $2 \times 2$ matrix over $R$ is MoorePenrose invertible. This follows from the facts that
(a) $S C_{2}$ holds in $R$ if and only if $S C_{1}$ holds in $R_{2 \times 2}$.
(b) $R$ is regular if and only if the ring $R_{2 \times 2}$ is regular.
(v) The previous item shows that the regularity of the involutory ring $R$ together with $S C_{2}$ is sufficient to garantee the existence of $A^{\dagger}$, for any $2 \times 2$ matrix $A$ over $R$, with respect to the special involution in $R_{2 \times 2}$ induced by the involution on $R$. In the remainder of this paper we will give an expression for the Moore-Penrose inverse of a $2 \times 2$ matrix over $R$, and for this we will need the symmetry of $R_{2 \times 2}$.

We note that symmetry of $R_{2 \times 2}$ does not follow from $R$ being regular and satisfying $S C_{2}$. Indeed, set $R=\mathbb{Z}_{7}$ which is a field and thus regular. The involution we take is the identity map. The squares are $\{0,1,2,4\}$. It is clear that $x^{2}+y^{2}=0 \Rightarrow x=0=y$. That is, $S C_{2}$ holds. Now, let $M=\left[\begin{array}{ll}2 & 0 \\ 3 & 0\end{array}\right]$. Then $M^{*} M=M^{T} M=\left[\begin{array}{ll}2 & 3 \\ 0 & 0\end{array}\right]\left[\begin{array}{ll}2 & 0 \\ 3 & 0\end{array}\right]=\left[\begin{array}{ll}6 & 0 \\ 0 & 0\end{array}\right]$. Hence $I_{2}+M^{*} M=\left[\begin{array}{ll}0 & 0 \\ 0 & 1\end{array}\right]$, which is not invertible.
(vi) In a regular symmetric ring, idempotents $e$ have a Moore-Penrose inverse via $e^{\dagger}=e^{*}\left[1+\left(e-e^{*}\right)\left(e^{*}-e\right)\right]^{-1}$. Indeed, setting $u=1+\left(e-e^{*}\right)\left(e^{*}-e\right)$, then $u$ and $e e^{*}$ commute, and so do $u^{-1}$ and $e e^{*}, u$ and $e^{*} e$ commute, and so do $u^{-1}$ and $e^{*} e$, and also $u^{-1} e e^{*} e e=e=e e^{*} e u^{-1}$. Since $e\left(e^{*} u^{-1}\right)$ and $\left(u^{-1} e^{*}\right) e$ are symmetric, and $e\left(e^{*} u^{-1}\right) e=u^{-1} e e^{*} e=e=e e^{*} e u^{-1}=e\left(u^{-1} e^{*}\right) e$, then $e^{*} u^{-1}$ is a 1-3 inverse of $e$ and $u^{-1} e^{*}$ is a 1-4 inverse of $e$, which lead to $e^{\dagger}=$ $\left(u^{-1} e^{*}\right) e\left(e^{*} u^{-1}\right)=u^{-1} e^{*} e e^{*} u^{-1}=e^{*} u^{-1}$.
As such the orthogonal projections $P_{a R}$ and $P_{R a}$ can be defined as $p=\left(a a^{-}\right)\left(a a^{-}\right)^{\dagger}$ and $q=\left(a^{-} a\right)^{\dagger}\left(a^{-} a\right)$. It then follows that the Moore-Penrose inverse $a^{\dagger}=q a^{-} p$ exists and the $S C_{1}$ property follows.

### 2.1. The Moore-Penrose inverse of a sum

We recall that if $c a^{*}=0$, then $a+c$ has a Moore-Penrose inverse, which takes the form

$$
\begin{equation*}
(a+c)^{\dagger}=\left(1+y^{*}\right)\left(1+y y^{*}\right)^{-1} s+u^{\dagger} \tag{3}
\end{equation*}
$$

where

$$
\begin{aligned}
u & =\left(1-a a^{\dagger}\right) c \\
s & =a^{\dagger}\left(1-c u^{\dagger}\right) \\
y & =a^{\dagger} c\left(1-u^{\dagger} u\right)=s c .
\end{aligned}
$$

Indeed, and since $1-y^{*} y\left(1+y^{*} y\right)^{-1}=\left(1+y^{*} y\right)^{-1},\left(1+y^{*} y\right)^{-1}$ and $1-u^{\dagger} u$ commute, $y\left(1-u^{\dagger} u\right)=y$, and $\left(1+y y^{*}\right)^{-1}=1-y\left(1+y^{*} y\right)^{-1} y^{*}$, then, using [4, Theorem 2],

$$
\begin{aligned}
(a+c)^{\dagger}= & a^{\dagger}-a^{\dagger} c u^{\dagger}-a^{\dagger} c\left(1-u^{\dagger} u\right)\left(1+y^{*} y\right)^{-1} c^{*} a^{\dagger^{*}} a^{\dagger}\left(1-c u^{\dagger}\right)+u^{\dagger}+ \\
& +\left(1-u^{\dagger} u\right)\left(1+y^{*} y\right)^{-1} c^{*} a^{\dagger^{*}} a^{\dagger}\left(1-c u^{\dagger}\right) \\
= & s-y\left(1+y^{*} y\right)^{-1} c^{*} a^{\dagger^{*}} s+u^{\dagger}+\left(1-u^{\dagger} u\right)\left(1+y^{*} y\right)^{-1} c^{*} a^{\dagger^{*}} s \\
= & s-y\left(1+y^{*} y\right)^{-1}\left(1-u^{\dagger} u\right) c^{*} a^{\dagger^{*}} s+u^{\dagger}+\left(1-u^{\dagger} u\right)\left(1-y^{*} y\left(1+y^{*} y\right)^{-1}\right) c^{*} a^{\dagger^{*}} s \\
= & s-y\left(1+y^{*} y\right)^{-1} y^{*} s+u^{\dagger}+y^{*} s+y^{*} y\left(1+y^{*} y\right)^{-1} c^{*} a^{\dagger^{*}} s \\
= & u^{\dagger}+\left(1-y\left(1+y^{*} y\right)^{-1} y^{*}\right) s+y^{*}\left(1-y\left(1+y^{*} y\right)^{-1} y^{*}\right) s \\
= & u^{\dagger}+\left(1+y y^{*}\right)^{-1} s+y^{*}\left(1+y y^{*}\right)^{-1} s \\
= & \left(1+y^{*}\right)\left(1+y y^{*}\right)^{-1} s+u^{\dagger}
\end{aligned}
$$

Moreover, we also have, from [3, Theorem 2] (also from [7, Lemma 2]),

$$
\left[\begin{array}{l}
a  \tag{4}\\
b
\end{array}\right]^{\dagger}=\left[\xi a^{*}, \xi b^{*}\right] \text { and }\left[\begin{array}{l}
a \\
b
\end{array}\right]^{\dagger}\left[\begin{array}{l}
a \\
b
\end{array}\right]=a^{\dagger} a+v^{\dagger} v
$$

where $\xi=\left(a^{*} a+b^{*} b\right)^{\dagger}$ and $v=b\left(1-a^{\dagger} a\right)$. We may re-express the former element as

$$
\begin{equation*}
\xi=t \mu^{-1} t^{*}+\left(v^{*} v\right)^{\dagger} \tag{5}
\end{equation*}
$$

in which

$$
\begin{equation*}
t=\left(1-v^{\dagger} b\right) a^{\dagger}, \quad x=\left(1-v v^{\dagger}\right) b a^{\dagger}=b t, \quad \mu=1+x^{*} x \tag{6}
\end{equation*}
$$

Indeed, from [4, Theorem 1],

$$
\xi=\left(a^{*} a+b^{*} b\right)^{\dagger}=t \ell t^{*}+v^{\dagger}\left(v^{*}\right)^{\dagger}
$$

where

$$
\ell=1-\left(\left(1-v v^{\dagger}\right) b a^{\dagger}\right)^{*} k\left(b a^{\dagger}\right)
$$

and

$$
k=\left(1+\left(1-v v^{\dagger}\right) b a^{\dagger}\left(\left(1-v v^{\dagger}\right) b a^{\dagger}\right)^{*}\right)^{-1}=\left(1+x x^{*}\right)^{-1}
$$

Since $\left(1-v v^{\dagger}\right) k=k\left(1-v v^{\dagger}\right)=\left(1-v v^{\dagger}\right) k\left(1-v v^{\dagger}\right)$,

$$
\begin{aligned}
\ell & =1-\left(b a^{\dagger}\right)^{*}\left(1-v v^{\dagger}\right) k\left(1-v v^{\dagger}\right) b a^{\dagger} \\
& =1-\left(\left(1-v v^{\dagger}\right) b a^{\dagger}\right)^{*} k\left(1-v v^{\dagger}\right) b a^{\dagger} \\
& =1-x^{*}\left(1+x x^{*}\right)^{-1} x \\
& =\left(1+x^{*} x\right)^{-1}=\mu^{-1}
\end{aligned}
$$

Lastly, $v^{\dagger}\left(v^{*}\right)^{\dagger}=\left(v^{*} v\right)^{\dagger}$ by [5, Lemma 5], or simply by checking the Penrose equations (2).

### 2.2. The lower triangular case

Consider the $2 \times 2$ triangular matrix $M=\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]$. We may split $M$ as

$$
M=\left[\begin{array}{ll}
0 & 0 \\
0 & d
\end{array}\right]+\left[\begin{array}{ll}
a & 0 \\
b & 0
\end{array}\right]=\mathscr{A}+\mathscr{C},
$$

where $\mathscr{C} \mathscr{A}^{*}=0$. In order to apply (3) to this semi-orthogonal splitting, we need to show that $I+A^{*} A$ is invertible for any matrix $A \in R_{2 \times 2}$. This we now undertake.

The key fact is the following factorization. If $\alpha$ is a unit then

$$
\left[\begin{array}{ll}
\alpha & \beta^{*}  \tag{7}\\
\beta & \delta
\end{array}\right]=\left[\begin{array}{cc}
1 & 0 \\
\beta \alpha^{-1} & 1
\end{array}\right]\left[\begin{array}{ll}
\alpha & 0 \\
0 & z
\end{array}\right]\left[\begin{array}{lc}
1 & \alpha^{-1} \beta^{*} \\
0 & 1
\end{array}\right]
$$

where $z$ is the Schur complement $z=\delta-\beta \alpha^{-1} \beta^{*}$. Now consider the matrix $A=$ $[\mathbf{a}, \mathbf{b}]=\left[\begin{array}{ll}a_{1} & b_{1} \\ a_{2} & b_{2}\end{array}\right]$. Then

$$
I+A^{*} A=\left[\begin{array}{cc}
1+\mathbf{a}^{*} \mathbf{a} & \mathbf{a}^{*} \mathbf{b}  \tag{8}\\
\mathbf{b}^{*} \mathbf{a} & 1+\mathbf{b}^{*} \mathbf{b}
\end{array}\right]
$$

and its Schur complement becomes

$$
\begin{aligned}
z & =1+\mathbf{b}^{*} \mathbf{b}-\left(\mathbf{b}^{*} \mathbf{a}\right)\left(1+\mathbf{a}^{*} \mathbf{a}\right)^{-1} \mathbf{a}^{*} \mathbf{b} \\
& =1+\mathbf{b}^{*}\left[I_{2}-\mathbf{a}\left(1+\mathbf{a}^{*} \mathbf{a}\right)^{-1} \mathbf{a}^{*}\right] \mathbf{b} \\
& =1+\mathbf{b}^{*}\left[I_{2}+\mathbf{a} \mathbf{a}^{*}\right]^{-1} \mathbf{b}
\end{aligned}
$$

since $\left(I_{2}+\mathbf{a} \mathbf{a}^{*}\right)^{-1}=I_{2}-\mathbf{a}\left(1+\mathbf{a}^{*} \mathbf{a}\right)^{-1} \mathbf{a}^{*}$.
We now turn to the matrix

$$
\begin{aligned}
G & =I+\mathbf{a a}^{*} \\
& =\left[\begin{array}{cc}
1+a_{1} a_{1}^{*} & a_{1} a_{2}^{*} \\
a_{2} a_{1}^{*} & 1+a_{2} a_{2}^{*}
\end{array}\right] \\
& =\left[\begin{array}{cc}
1 & 0 \\
\beta \alpha^{-1} & 1
\end{array}\right]\left[\begin{array}{ll}
\alpha & 0 \\
0 & \zeta
\end{array}\right]\left[\begin{array}{cc}
1 & \alpha^{-1} \beta^{*} \\
0 & 1
\end{array}\right]
\end{aligned}
$$

where $\alpha=1+a_{1} a_{1}^{*}$ is a unit, $\beta=a_{2} a_{1}^{*}$ and the Schur complement $\zeta$ takes the form

$$
\begin{aligned}
\zeta & =1+a_{2} a_{2}^{*}-a_{2} a_{1}^{*}\left(1+a_{1} a_{1}^{*}\right)^{-1} a_{1} a_{2}^{*} \\
& =1+a_{2}\left(1-a_{1}^{*}\left(1+a_{1} a_{1}^{*}\right)^{-1} a_{1}\right) a_{2}^{*} \\
& =1+a_{2}\left(1+a_{1}^{*} a_{1}\right)^{-1} a_{2}^{*},
\end{aligned}
$$

since $\left(1+a_{1}^{*} a_{1}\right)^{-1}=1-a_{1}^{*}\left(1+a_{1} a_{1}^{*}\right)^{-1} a_{1}$.
By using the square root axiom, we may set $1+a_{1} a_{1}^{*}=e e^{*}$ and therefore $e$ is a unit. Consequentely, there exists $f$ such that $\left(1+a_{1} a_{1}^{*}\right)^{-1}=f f^{*}$ and hence $\zeta=$ $1+\left(a_{2} f\right)\left(a_{2} f\right)^{*}$. Again $\zeta$ is a unit, and by the square root axiom, $\zeta=h h^{*}$, which leads to $\zeta^{-1}=g g^{*}$, for some $g$.

Substituting into $z$ now gives

$$
\begin{aligned}
z & =1+\mathbf{b}^{*}\left(I+\mathbf{a a}^{*}\right)^{-1} \mathbf{b} \\
& =1+\mathbf{b}^{*}\left[\begin{array}{cc}
1-\alpha^{-1} \beta^{*} \\
0 & 1
\end{array}\right]\left[\begin{array}{cc}
\alpha^{-1} & 0 \\
0 & \zeta^{-1}
\end{array}\right]\left[\begin{array}{cc}
1 & 0 \\
-\beta \alpha^{-1} & 1
\end{array}\right] \mathbf{b} \\
& =1+\left[b_{1}^{*} w^{*}\right]\left[\begin{array}{cc}
f f^{*} & 0 \\
0 & g g^{*}
\end{array}\right]\left[\begin{array}{c}
b_{1} \\
w
\end{array}\right] \\
& =1+b_{1}^{*} f f^{*} b_{1}+w^{*} g g^{*} w,
\end{aligned}
$$

where $w=b_{2}-\beta \alpha^{-1} b_{1}$, and therefore $z$ is a unit. Thus $R_{2 \times 2}$ is again symmetric.

We now may apply (3) to our matrix $M$, giving

$$
\begin{equation*}
M^{\dagger}=\mathscr{U}^{\dagger}+\left(I+\mathscr{Y}^{*}\right)\left(I+\mathscr{Y} \mathscr{Y}^{*}\right)^{-1} S \tag{9}
\end{equation*}
$$

where

$$
\mathscr{U}=\left(I-\mathscr{A}_{\mathscr{A}^{\dagger}}\right) \mathscr{C}=\left[\begin{array}{ll}
a & 0 \\
B & 0
\end{array}\right],
$$

with $B=\left(1-d d^{\dagger}\right) b$, and

$$
\mathscr{Y}=\mathscr{A}^{\dagger} \mathscr{C}\left(I-\mathscr{U}^{\dagger} \mathscr{U}\right)
$$

We next compute $\mathscr{U}^{\dagger}=\left[\begin{array}{cc}\xi a^{*} & \xi B^{*} \\ 0 & 0\end{array}\right]$ in which

$$
\begin{aligned}
\xi & =\left(a^{*} a+B^{*} B\right)^{\dagger}=t \mu^{-1} t^{*}+\left(v^{*} v\right)^{\dagger} \\
v & =B\left(1-a^{\dagger} a\right), t=\left(1-v^{\dagger} B\right) a^{\dagger} \\
\mu & =I+x^{*} x, \quad \text { and } \\
x & =\left(I-v v^{\dagger}\right) B a^{\dagger}=B t
\end{aligned}
$$

By combining these, and by using the equalities in (4), we arrive at

$$
\mathscr{U}^{\dagger} \mathscr{U}=\left[\begin{array}{cc}
a^{\dagger} a+v^{\dagger} v & 0 \\
0 & 0
\end{array}\right]
$$

and

$$
\mathscr{Y}=\left[\begin{array}{ll}
0 & 0 \\
0 & d^{\dagger}
\end{array}\right]\left[\begin{array}{ll}
a & 0 \\
b & 0
\end{array}\right]\left[\begin{array}{ccc}
1-a^{\dagger} a-v^{\dagger} & 0 \\
0 & 1
\end{array}\right]=\left[\begin{array}{ll}
0 & 0 \\
f & 0
\end{array}\right],
$$

where

$$
f=d^{\dagger} b\left(1-a^{\dagger} a-v^{\dagger} v\right)
$$

Likewise,

$$
\begin{aligned}
\mathscr{S} & =\mathscr{A}^{\dagger}-\mathscr{A}^{\dagger} \mathscr{C} \mathscr{U}^{\dagger} \\
& =\left[\begin{array}{cc}
0 & 0 \\
0 & d^{\dagger}
\end{array}\right]-\left[\begin{array}{ll}
0 & 0 \\
0 & d^{\dagger}
\end{array}\right]\left[\begin{array}{ll}
a & 0 \\
b & 0
\end{array}\right]\left[\begin{array}{cc}
\xi a^{*} & \xi B^{*} \\
0 & 0
\end{array}\right] \\
& =\left[\begin{array}{cc}
0 \\
-d^{\dagger} b \xi a^{*} & d^{\dagger}-d^{\dagger} b \xi B^{*}
\end{array}\right] .
\end{aligned}
$$

We then compute

$$
\left(I+\mathscr{Y}^{*}\right)\left(I+\mathscr{Y} \mathscr{Y}^{*}\right)^{-1}=\left[\begin{array}{c}
1 f^{*}\left(1+f f^{*}\right)^{-1} \\
0\left(1+f f^{*}\right)^{-1}
\end{array}\right]
$$

followed by

$$
\left(I+\mathscr{Y}^{*}\right)\left(I+\mathscr{Y} \mathscr{Y}^{*}\right)^{-1} \mathscr{S}=\left[\begin{array}{cc}
1 & f^{*}\left(1+f f^{*}\right)^{-1} \\
0 & \left(1+f f^{*}\right)^{-1}
\end{array}\right]\left[\begin{array}{cc}
0 & 0 \\
-d^{\dagger} b \xi a^{*} d^{\dagger}-d^{\dagger} b \xi B^{*}
\end{array}\right] .
$$

This then gives, using equation (9),

$$
\left[\begin{array}{ll}
a & 0  \tag{10}\\
b & d
\end{array}\right]^{\dagger}=\left(I+\mathscr{Y}^{*}\right)\left(I+\mathscr{Y} \mathscr{Y}^{*}\right)^{-1} \mathscr{S}+\mathscr{U}^{\dagger}=\left[\begin{array}{cc}
p & q \\
s & r
\end{array}\right]
$$

where

$$
\begin{align*}
p & =\xi a^{*}-\left(1+f^{*} f\right)^{-1} f^{*} d^{\dagger} b \xi a^{*}  \tag{11}\\
s & =-\left(1+f f^{*}\right)^{-1} d^{\dagger} b \xi a^{*}  \tag{12}\\
q & =\xi b^{*}\left(1-d d^{\dagger}\right)+\left(1+f^{*} f\right)^{-1} f^{*} d^{\dagger}\left[1-b \xi b^{*}\left(1-d d^{\dagger}\right)\right]  \tag{13}\\
r & =\left(1+f f^{*}\right)^{-1} d^{\dagger}\left[1-b \xi^{*} b^{*}\left(1-d d^{\dagger}\right)\right] \tag{14}
\end{align*}
$$

in which

$$
\begin{align*}
\xi & =\left[a^{*} a+b^{*}\left(1-d d^{\dagger}\right) b\right]^{\dagger}=\xi^{*}=t\left(1+x^{*} x\right)^{-1} t^{*}+\left(v^{*} v\right)^{\dagger}  \tag{15}\\
x & =\left(1-v v^{\dagger}\right)\left(1-d d^{\dagger}\right) b a^{\dagger}  \tag{16}\\
t & =\left[1-v^{\dagger}\left(1-d d^{\dagger}\right) b\right] a^{\dagger}  \tag{17}\\
f & =d^{\dagger} b\left(1-a^{\dagger} a-v^{\dagger} v\right) \text { and }  \tag{18}\\
v & =\left(1-d d^{\dagger}\right) b\left(1-a^{\dagger} a\right) \text { (corner stone). } \tag{19}
\end{align*}
$$

We have presented an alternative expression to main theorem of [7] for the MoorePenrose inverse of a $2 \times 2$ lower triangular matrix.

For later use, we observe that
(a) $v a^{*}=0,\left(v^{*} v\right)^{\dagger} a^{*}=0$.
(b) $\xi a^{*}=t \mu^{-1} t^{*} a^{*}$,where $\mu=1+x^{*} x$.
(c) $t=\left[1-v^{\dagger}\left(1-d d^{\dagger}\right) b\right] a^{\dagger}=\left(1-v^{\dagger} b\right) a^{\dagger}=a^{\dagger}-v^{\dagger} b a^{\dagger}$.
(d) $t a a^{\dagger}=t$ and $a t=a a^{\dagger}-a v^{\dagger} b a^{\dagger}=a a^{\dagger}=(a t)^{*}$
(e) $x a a^{\dagger}=x$ and so $a a^{\dagger} x^{*}=x^{*}$, and
(f) $\mu a a^{\dagger}=a a^{\dagger} \mu$ and $\mu^{-1} a a^{\dagger}=a a^{\dagger} \mu^{-1}$.

From the above,

$$
\xi a^{*}=t \mu^{-1}(a t)^{*}=t \mu^{-1} a a^{\dagger}=t a a^{\dagger} \mu^{-1}=t \mu^{-1}
$$

The equality

$$
\begin{equation*}
\xi a^{*}=t \mu^{-1} \tag{20}
\end{equation*}
$$

will be used later in this document.

## 3. The four "faces" of $M^{\dagger}$

We now examine the four cases where the block lower triangular matrix $M=$ $\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]$ has a Moore-Penrose inverse of the form:
(i) $M^{\dagger}=\left[\begin{array}{l}* 0 \\ * *\end{array}\right]$ the $(1,2,0)$ case (unflipped),
(ii) $M^{\dagger}=\left[\begin{array}{l}* * \\ 0 *\end{array}\right]$ the $(2,1,0)$ case (flipped),
(iii) $M^{\dagger}=\left[\begin{array}{l}0 * \\ * *\end{array}\right]$ the $(1,1,0)$ case,
(iv) $M^{\dagger}=\left[\begin{array}{l}* * \\ * 0\end{array}\right]$ the $(2,2,0)$ case.

### 3.1. The $(1,2,0)$ case (unflipped)

The Moore-Penrose inverse of the block lower triangular matrix $M=\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]$ is again of $(1,2,0)$ type if and only if $b=d d^{\dagger} b=b a^{\dagger} a$ (see [12]).

We may also use the general triangular case (10) to rederive this consistency. Indeed this occurs precisely when

$$
0=q=\xi b^{*}\left(1-d d^{\dagger}\right)+\left(1+f^{*} f\right)^{-1} f^{*} d^{\dagger}\left[1-b \xi b^{*}\left(1-d d^{\dagger}\right)\right]
$$

By post-multiplying by $d d^{\dagger}$ gives $\left(1+f^{*} f\right)^{-1} f^{*} d^{\dagger}=0$ which reduces to $d f=0$. By substituting this back into $q$, then shows that also $\xi b^{*}\left(1-d d^{\dagger}\right)=0$. Thus $M^{\dagger}$ has the desired lower triangular form if and only if

$$
\begin{equation*}
d f=0 \text { and } \xi b^{*}\left(1-d d^{\dagger}\right)=0 \tag{21}
\end{equation*}
$$

Now recall that if $B=\left(1-d d^{\dagger}\right) b$ then $\xi=\left(a^{*} a+B^{*} B\right)^{\dagger}$. Hence the second consistency condition becomes $\left(a^{*} a+B^{*} B\right)^{\dagger} B^{*}=0$, which is equivalent to $\left(a^{*} a+\right.$ $\left.B^{*} B\right) B^{*}=0$. This implies that $B\left(a^{*} a+B^{*} B\right) B^{*}=0$ and hence by star-cancellation, $B B^{*}=0$ and thus $B=0$. This says that $\mathrm{b}=d d^{\dagger} b$ and hence $v=0$.

By substituting in $0=d f=d d^{\dagger} b\left[1-a^{\dagger} a-v^{\dagger} v\right]$ then yields $0=b\left(1-a^{\dagger} a\right)$, and we recover the necessary condition $b=d d^{\dagger} b a^{\dagger} a$, which is also sufficient. We have proved

THEOREM 3.1. Given $M=\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]$, the following conditions are equivalent:

1. $M^{\dagger}$ is of $(1,2,0)$ type.
2. $b \in d R a$.
3. $b=d d^{\dagger} b a^{\dagger} a$.
4. $d d^{\dagger} b=b=b a^{\dagger} a$.

In this case, $M^{\dagger}=\left[\begin{array}{cc}a^{\dagger} & 0 \\ -d^{\dagger} b a^{\dagger} & d^{\dagger}\end{array}\right]$.
This can be extended to the $n \times n$ case (as in [6]).

### 3.2. The $(2,1,0)$ case (flipped)

Next we examine the case here the Moore-Penrose inverse of the lower triangular matrix $M$ "flips" and takes the form $M^{\dagger}=\left[\begin{array}{cc}p & q \\ 0 & r\end{array}\right]$ for some $p, q$, and $r$. We will give necessary and sufficient conditions for this to happen, in terms of the blocks $a, b$ and $d$.

From (12) we see that a necessary and sufficient condition for $M^{\dagger}$ to have the flipped form $\left[\begin{array}{cc}p & q \\ 0 & r\end{array}\right]$ is that $d^{\dagger} b \xi a^{*}=0$.

We now observe from Equation (20), that the consistency condition collapses to $0=d^{*} b t=d^{*} b\left(1-v^{\dagger} b\right) a^{\dagger}$, which yields

$$
\begin{equation*}
d^{*} b a^{*}=d^{*} b v^{\dagger} b a^{*} \tag{22}
\end{equation*}
$$

or equivalently

$$
d d^{\dagger} b\left(b^{\dagger}-v^{\dagger}\right) b a^{\dagger} a=0
$$

We thus have
THEOREM 3.2. Given $M=\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]$, then $M^{\dagger}$ is of $(2,1,0)$ type if and only if

$$
d d^{\dagger} b\left(b^{\dagger}-v^{\dagger}\right) b a^{\dagger} a=0
$$

in which case

$$
M^{\dagger}=\left[\begin{array}{cc}
\xi a^{*} \xi b^{*}\left(1-d d^{\dagger}\right)+\left(1+f^{*} f\right)^{-1} f^{*} d^{\dagger}\left[1-b \xi b^{*}\left(1-d d^{\dagger}\right)\right] \\
0 & \left(1+f f^{*}\right)^{-1} d^{\dagger}\left[1-b \xi^{*} b^{*}\left(1-d d^{\dagger}\right)\right]
\end{array}\right]
$$

where $\xi, f$ are as above.

If we set $e=a^{\dagger} a$ and $f=d d^{\dagger}$, then the consistency condition can be written as

$$
\zeta=f b e-f b[(1-f) b(1-e)]^{\dagger} b e=0
$$

which is the $(2,2)$ Schur complement in $\left[\begin{array}{cc}f b e & b e \\ f b & (1-f) b(1-e)\end{array}\right]$. It only involves $b$, $e$ and $f$. It is not clear how to simplify this condition. All we have is that $v v^{\dagger}=$ $\left(1-d d^{\dagger}\right) b v^{\dagger}$.

### 3.3. The $(2,2,0)$ case

From (10) we see that $M^{\dagger}$ is of $(2,2,0)$ type if and only if $r=0$, which is equivalent to

$$
d^{\dagger}=d^{\dagger} b \xi^{*} b^{*}\left(1-d d^{\dagger}\right)
$$

Right multiplication by $d d^{\dagger}$ shows that necessarily $d^{\dagger}=0$, that is, $d=0$. The sufficiency is clear. We may thus state the following result:

THEOREM 3.3. Given $M=\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right], M^{\dagger}$ is of $(2,2,0)$ type if and only if $d=0$, in which case $\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right]^{\dagger}=\left[\begin{array}{cc}\left(a^{*} a+b^{*} b\right)^{\dagger} a^{*}\left(a^{*} a+b^{*} b\right)^{\dagger} b^{*} \\ 0 & 0\end{array}\right]$.

### 3.4. The $(1,1,0)$ case

Lastly, we analyze the case where $M^{\dagger}$ is of $(1,1,0)$ type. This corresponds to

$$
p=\xi a^{*}-\left(1+f^{*} f\right)^{-1} f^{*} d^{\dagger} b \xi a^{*}=0
$$

with $\xi=\left(a^{*} a+B^{*} B\right)^{\dagger}, B=\left(1-d d^{\dagger}\right) b, f=d^{\dagger} b\left(1-a^{\dagger} a-v^{\dagger} v\right)$ and $v=\left(1-d d^{\dagger}\right) b(1-$ $\left.a^{\dagger} a\right)$.

Now recall, from equation (20), that $\xi a^{*}=t \mu^{-1}$, where $\mu=1+x^{*} x=(1-$ $\left.v v^{\dagger}\right) B a^{\dagger}$ and $t=a^{\dagger}-v^{\dagger} b a^{\dagger}$. Thus $p=0$ is equivalent to $t \mu^{-1}=\left(1+f^{*} f\right)^{-1} f^{*} d^{\dagger} b t \mu^{-1}$, i.e. to

$$
\begin{equation*}
\left(1+f^{*} f\right) t=f^{*} d^{\dagger} b t \tag{23}
\end{equation*}
$$

Since $v a^{\dagger}=0=a v^{\dagger}$ we know that $f a^{\dagger}=0=f v^{\dagger}$ and consequently $f t=0$.
The equality (23) now reduces to $t=f^{*} d^{\dagger} b t$. Lastly, left multiplication by $a^{\dagger} a$ shows that necessarily $a^{\dagger}=0$, that is, $a=0$. This is, trivially, sufficient for $p=0$. We may thus conclude that

THEOREM 3.4. Given $M=\left[\begin{array}{ll}a & 0 \\ b & d\end{array}\right], M^{\dagger}$ is of $(1,1,0)$ type if and only if $a=0$, in which case $\left[\begin{array}{ll}0 & 0 \\ b & d\end{array}\right]^{\dagger}=\left[\begin{array}{ll}0 & b^{*}\left(b b^{*}+d d^{*}\right)^{\dagger} \\ 0 & d^{*}\left(b b^{*}+d d^{*}\right)^{\dagger}\end{array}\right]$.

It is easily seen that these reduce to Cline's result.

## 4. Questions and remarks

1. The consistency condition for the Moore-Penrose inverse to flip involves the corner matrix $v=\left(1-d d^{\dagger}\right) b\left(1-a^{\dagger} a\right)$. Its Moore-Penrose inverse is a perturbation of $b^{\dagger}$.
2. Can we use the theory of Schur complements or partial orders, to simplify the consistency condition $f b e=f b[(1-f) b(1-e)]^{\dagger} b e$ ?
3. No further simplification of the Condition (22) seems possible.
4. The unflipped case can be, inductively, generalized to the $n \times n$ case. What can be said for the flipped case for $n \times n$ matrices?
5. To ensure the symmetry of $R_{2 \times 2}$ with $R$ regular and symmetric, we may replace the square-root axiom on $R$ by the condition $S C_{4}$.
6. $S C_{2}(R)$ does not imply $S C_{n}(R)$, nor implies the square root property, as remarked in Examples 2 and 3 in [13, page 215].
7. We have not used any of the other conditions that relate $p, q$, and $r$ to $a, b$ and $d$.

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R.E. Hartwig

Mathematics Department, N.C.S.U. Raleigh, NC 27695-8205 U.S.A.
e-mail: hartwig@unity.ncsu.edu
P. Patrício

Departamento de Matemática e Aplicações
Universidade do Minho 4710-057 Braga

Portugal
e-mail: pedro@math.uminho.pt


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