# COMPLEX SYMMETRIC WEIGHTED COMPOSITION-DIFFERENTIATION OPERATORS OF ORDER $n$ ON THE WEIGHTED BERGMAN SPACES 

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Abstract. We study the complex symmetry of weighted composition-differentiation operators of order $n$ on the weighted Bergman spaces $A_{\alpha}^{2}$. Several concrete examples are provided.

## 1. Preliminaries

Let $\mathbb{D}$ denote the open disk in the complex plane $\mathbb{C}$. For $\alpha>-1$, the weighted Bergman space $A_{\alpha}^{2}$ is the Hilbert space consisting of all analytic functions $f(z)=$ $\sum_{j=0}^{\infty} a_{j} z^{j}$ on $\mathbb{D}$ such that $\|f\|^{2}=\sum_{j=0}^{\infty}\left|a_{j}\right|^{2} \beta(j)^{2}<\infty$, where

$$
\beta(j)=\left\|z^{j}\right\|=\sqrt{\frac{j!\Gamma(\alpha+2)}{\Gamma(j+\alpha+2)}}
$$

for each non-negative integer $j$. The inner product of two functions in this space is given by the rule

$$
\left\langle\sum_{j=0}^{\infty} a_{j} z^{j}, \sum_{j=0}^{\infty} b_{j} z^{j}\right\rangle=\sum_{j=0}^{\infty} a_{j} \overline{b_{j}} \beta(j)^{2}
$$

It is well known that this space is a reproducing kernel Hilbert space; for any $w$ in $\mathbb{D}$ and any non-negative integer $m$, there is a kernel function $K_{w}^{[m]}$ such that $\left\langle f, K_{w}^{[m]}\right\rangle=$ $f^{(m)}(w)$ for each $f$ in $A_{\alpha}^{2}$. To simplify notation, we write $K_{w}$ to denote $K_{w}^{[0]}$. In particular,

$$
K_{w}(z)=\frac{1}{(1-\bar{w} z)^{\alpha+2}}=\sum_{j=0}^{\infty} \frac{\bar{w}^{j} z^{j}}{\beta(j)^{2}}
$$

and

$$
K_{w}^{[m]}(z)=\frac{(\alpha+2) \ldots(\alpha+m+1) z^{m}}{(1-\bar{w} z)^{m+\alpha+2}}=\frac{m!z^{m}}{\beta(m)^{2}(1-\bar{w} z)^{m+\alpha+2}}
$$

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for $m \geqslant 1$ (note that from [2, p. 20], we can see that $K_{w}^{[m]}(z)=\frac{d^{m} K}{d \bar{w}^{m}}$, where $K(z, \bar{w})=$ $K_{w}(z)$ for each $z, w \in \mathbb{D}$ and $\left.\beta(m)^{2}=\frac{m!}{(\alpha+2) \ldots(\alpha+m+1)}\right)$. Moreover, for each nonnegative integer $m$, we have

$$
\left\|K_{w}^{[m]}\right\|^{2}=\sum_{j=m}^{\infty} \frac{\left(|w|^{2}\right)^{j-m}}{\beta(j)^{2}}\left(\frac{j!}{(j-m)!}\right)^{2}
$$

(note that $K_{w}^{[m]}(z)=\sum_{j=m}^{\infty} \frac{j!}{\beta(j)^{2}(j-m)!} \bar{w}^{j-m} z^{j}$ by [2, Theorem 2.16]). Recall that $H^{\infty}$ is the Banach space consisting of all bounded analytic functions defined on $\mathbb{D}$, with supremum norm $\|f\|_{\infty}=\sup _{z \in \mathbb{D}}|f(z)|$. Let $P_{\alpha}$ denote the projection of $L^{2}\left(\mathbb{D}, d A_{\alpha}\right)$ onto $A_{\alpha}^{2}$. Given a function $h$ in $L^{\infty}(\mathbb{D})$, the Toeplitz operator $T_{h}$ on $A_{\alpha}^{2}$ is defined by the rule

$$
T_{h}(f)=P_{\alpha}(h f)
$$

for $f$ in $A_{\alpha}^{2}$. If $h$ belongs to $H^{\infty}$, it is easy to see that $T_{h}(f)=h \cdot f$. For an analytic self-map $\varphi: \mathbb{D} \rightarrow \mathbb{D}$, the composition operator $C_{\varphi}$ is defined by the rule

$$
C_{\varphi}(f)=f \circ \varphi
$$

for $f$ in $A_{\alpha}^{2}$. All Toeplitz operators and all composition operators are bounded on $A_{\alpha}^{2}$. As a natural generalization of both of these classes, consider the operator $C_{\psi, \varphi}$ that takes $f$ to $\psi \cdot(f \circ \varphi)$, where $\varphi: \mathbb{D} \rightarrow \mathbb{D}$ and $\psi: \mathbb{D} \rightarrow \mathbb{C}$ are both analytic on $\mathbb{D}$. Such an operator is called a weighted composition operator.

For a positive integer $n$, we define the differentiation operator of order $n$ on $A_{\alpha}^{2}$ by $D^{(n)}(f)=f^{(n)}$. None of these operators is bounded on $A_{\alpha}^{2}$. Nevertheless, for many analytic self-maps $\varphi$, the operator $C_{\varphi} D^{(n)}$ is bounded on $A_{\alpha}^{2}$. This class of operators was initially considered by Hibschweiler and Portnoy [9] and by Ohno [11], and has been studied further by other researchers (see [4], [5], and [12]). Ohno [11] characterized the boundedness and compactness of $C_{\varphi} D^{(1)}$ on the Hardy space; Stević [12] obtained analogous results for $C_{\varphi} D^{(n)}$ on the weighted Bergman spaces. We will write $D_{\varphi, n}$ to denote $C_{\varphi} D^{(n)}$, particularly when such an operator is bounded on $A_{\alpha}^{2}$, referring to it as a composition-differentiation operator of order $n$. For an analytic function $\psi: \mathbb{D} \rightarrow \mathbb{C}$, the weighted composition-differentiation operator of order $n$ on $A_{\alpha}^{2}$ is defined by the rule

$$
D_{\psi, \varphi, n}(f)=\psi \cdot\left(f^{(n)} \circ \varphi\right)
$$

Note that $D_{\psi, \varphi, n}$ is actually the product of the Toeplitz operator $T_{\psi}$ and $D_{\varphi, n}$, whenever $\psi$ belongs to $H^{\infty}$ and $D_{\varphi, n}$ is bounded. To avoid trivial situations, we will assume throughout this paper that $\varphi$ is not constant and that $\psi$ is not identically 0 .

A bounded linear operator $T$ is called complex symmetric on a complex Hilbert space $\mathscr{H}$ if there exits a conjugation $C$ (i.e., an antilinear isometric involution) such that $C T^{*} C=T$; for a particular conjugation $C$, we say that $T$ is $C$-symmetric. Garcia and Putinar initiated the study of complex symmetric operators on Hilbert spaces of analytic functions (see [7] and [8]). Complex symmetric weighted composition operators have been considered in [3], [6], [10], and [13]. In this paper, we use the symbol $J$ to denote the specific conjugation $(J f)(z)=\overline{f(\bar{z})}$.

Any complex number $z$ can be represented $z=|z| e^{i \theta}$, where $0 \leqslant \theta<2 \pi$. We write $\operatorname{Arg}(z)$ to denote this value of $\theta$, taking $\operatorname{Arg}(0)=0$.

## 2. Complex symmetric operators $D_{\psi, \varphi, n}$

For an analytic $\varphi: \mathbb{D} \rightarrow \mathbb{D}$ and $\alpha>-1$, the generalized Nevanlinna counting function $N_{\varphi, \alpha+2}$ is defined by the rule

$$
N_{\varphi, \alpha+2}(w)=\sum_{\varphi(z)=w}(\ln (1 /|z|))^{\alpha+2}
$$

where $w$ belongs to $\mathbb{D} \backslash\{\varphi(0)\}$. The next proposition provides necessary and sufficient conditions for $D_{\varphi, n}$ to be bounded and compact.

Proposition 2.1. [12, Theorem 9] Let $\varphi$ be an analytic self-map of $\mathbb{D}$, with $n$ in $\mathbb{N}$ and $\alpha>-1$.
a) An operator $D_{\varphi, n}: A_{\alpha}^{2} \rightarrow A_{\alpha}^{2}$ is bounded if and only if

$$
N_{\varphi, \alpha+2}(w)=O\left((\ln (1 /|w|))^{\alpha+2+2 n}\right) .
$$

b) An operator $D_{\varphi, n}: A_{\alpha}^{2} \rightarrow A_{\alpha}^{2}$ is compact if and only if

$$
N_{\varphi, \alpha+2}(w)=o\left((\ln (1 /|w|))^{\alpha+2+2 n}\right), \quad \text { as }|w| \rightarrow 1^{-} .
$$

Since $\ln (1 /|w|)$ is comparable to $1-|w|$ as $|w| \rightarrow 1^{-}$, the following characterization holds in the case where $\varphi$ is univalent on $\mathbb{D}$.

Corollary 2.2. Let $\varphi$ be a univalent self-map of $\mathbb{D}$, with $n$ in $\mathbb{N}$ and $\alpha>-1$.
a) An operator $D_{\varphi, n}$ is bounded on $A_{\alpha}^{2}$ if and only if

$$
\sup _{w \in \mathbb{D}} \frac{(1-|w|)^{\alpha+2}}{(1-|\varphi(w)|)^{\alpha+2+2 n}}<\infty .
$$

b) An operator $D_{\varphi, n}$ is compact on $A_{\alpha}^{2}$ if and only if

$$
\lim _{|w| \rightarrow 1} \frac{(1-|w|)^{\alpha+2}}{(1-|\varphi(w)|)^{\alpha+2+2 n}}=0
$$

Note that Corollary 2.2 shows that if $D_{\varphi, n}$ is bounded, then $\varphi$ does not have finite angular derivative at any point on $\partial \mathbb{D}$ (see [2, Theorem 2.44]). Moreover, we infer from Corollary 2.2 that an operator $D_{\varphi, n}$ is bounded if $\|\varphi\|_{\infty}<1$ and so $D_{\psi, \varphi, n}$ is bounded on $A_{\alpha}^{2}$ whenever $\psi$ belongs to $H^{\infty}$. We will employ the following lemma frequently.

Lemma 2.3. If an operator $D_{\psi, \varphi, n}$ is bounded on $A_{\alpha}^{2}$, then

$$
D_{\psi, \varphi, n}^{*}\left(K_{w}\right)=\overline{\psi(w)} K_{\varphi(w)}^{[n]}
$$

Proof. Observe that

$$
\left\langle f, D_{\psi, \varphi, n}^{*}\left(K_{w}\right)\right\rangle=\left\langle D_{\psi, \varphi, n} f, K_{w}\right\rangle=\psi(w) f^{(n)}(\varphi(w))=\left\langle f, \overline{\psi(w)} K_{\varphi(w)}^{[n]}\right\rangle
$$

for any $f$ in $A_{\alpha}^{2}$. Our result follows from the fact that the span of the kernel functions $K_{w}$ is dense in $A_{\alpha}^{2}$.

Throughout this paper, we set $t=(\alpha+2)(\alpha+3) \ldots(\alpha+n+1)$, which will be appeared several times in this paper. We will now make a few observations about $J$ symmetric operators $D_{\psi, \varphi, n}$, which will be used in the proof of Theorem 2.7.

Proposition 2.4. If an operator $D_{\psi, \varphi, n}$ is $J$-symmetric on $A_{\alpha}^{2}$, the following conditions hold:
(i) $\psi^{(m)}(0)=0$ for each $0 \leqslant m<n$;
(ii) $\psi^{(n)}(0) \neq 0$;
(iii) $\psi(w) \neq 0$ for any $w$ in $\mathbb{D} \backslash\{0\}$;
(iv) the map $\varphi$ is univalent.

Proof. Suppose that $D_{\psi, \varphi, n}$ is $J$-symmetric. Observe that

$$
\begin{equation*}
J D_{\psi, \varphi, n}\left(K_{0}\right)=0 \tag{2.1}
\end{equation*}
$$

Lemma 2.3 shows that

$$
\begin{equation*}
D_{\psi, \varphi, n}^{*} J\left(K_{0}\right)=\overline{\psi(0)} K_{\varphi(0)}^{[n]} \tag{2.2}
\end{equation*}
$$

Since $D_{\psi, \varphi, n}$ is $J$-symmetric, it follows from (2.1) and (2.2) that $\psi(0)=0$. Assume that $\psi^{(m)}(0)=0$ for $m<n-1$. One can see that

$$
\begin{equation*}
J D_{\psi, \varphi, n} K_{0}^{[m+1]}=0 \tag{2.3}
\end{equation*}
$$

On the other hand, for any $f$ in $A_{\alpha}^{2}$, we obtain

$$
\begin{aligned}
\left\langle f, D_{\psi, \varphi, n}^{*} J K_{0}^{[m+1]}\right\rangle & =\left\langle f, D_{\psi, \varphi, n}^{*} K_{0}^{[m+1]}\right\rangle \\
& =\left\langle D_{\psi, \varphi, n} f, K_{0}^{[m+1]}\right\rangle \\
& =\left(\psi \cdot\left(f^{(n)} \circ \varphi\right)\right)^{(m+1)}(0) \\
& =\sum_{j=0}^{m+1}\binom{m+1}{j} \psi^{(m+1-j)}(0)\left(f^{(n)} \circ \varphi\right)^{(j)}(0)
\end{aligned}
$$

$$
\begin{align*}
& =\psi^{(m+1)}(0) f^{(n)}(\varphi(0)) \\
& \quad \quad+\sum_{j=1}^{m+1}\binom{m+1}{j} \psi^{(m+1-j)}(0)\left(f^{(n)} \circ \varphi\right)^{(j)}(0) \\
& =\psi^{(m+1)}(0) f^{(n)}(\varphi(0)) \\
& =\left\langle f, \overline{\psi^{(m+1)}(0)} K_{\varphi(0)}^{[n]}\right\rangle \tag{2.4}
\end{align*}
$$

so

$$
\begin{equation*}
D_{\psi, \varphi, n}^{*} J K_{0}^{[m+1]}=D_{\psi, \varphi, n}^{*} K_{0}^{[m+1]}=\overline{\psi^{(m+1)}(0)} K_{\varphi(0)}^{[n]} \tag{2.5}
\end{equation*}
$$

If $D_{\psi, \varphi, n}$ is $J$-symmetric, then (2.3) and (2.5) imply that $\psi^{(m+1)}(0)=0$. By the same idea as in (2.4), we have

$$
\begin{equation*}
D_{\psi, \varphi, n}^{*} J K_{0}^{[n]}=D_{\psi, \varphi, n}^{*} K_{0}^{[n]}=\overline{\psi^{(n)}(0)} K_{\varphi(0)}^{[n]} \tag{2.6}
\end{equation*}
$$

since $\psi^{(m)}(0)=0$ for any $m<n$. Because

$$
\begin{equation*}
J D_{\psi, \varphi, n} K_{0}^{[n]}=\operatorname{tn}!J(\psi) \tag{2.7}
\end{equation*}
$$

and $\psi$ is not identically 0 , it follows from (2.6) and (2.7) that $\psi^{(n)}(0) \neq 0$. Now suppose that $\psi(w)=0$ for some $w$ in $\mathbb{D}$. Lemma 2.3 shows that $D_{\psi, \varphi, n}^{*} J\left(K_{\bar{w}}\right)=0$. Moreover,

$$
J D_{\psi, \varphi, n}\left(K_{\bar{w}}\right)=\frac{t \overline{w^{n}} J(\psi)}{(1-\bar{w} J(\varphi))^{n+\alpha+2}} .
$$

Since $D_{\psi, \varphi, n}$ is $J$-symmetric and $\psi$ is not identically zero, we observe that $w=0$.
Now assume that $D_{\psi, \varphi, n}$ is $J$-symmetric and that there exist distinct points $w_{1}$ and $w_{2}$ in $\mathbb{D}$ with $\varphi\left(w_{1}\right)=\varphi\left(w_{2}\right)$. (If either $w_{1}$ or $w_{2}$ is zero, the open mapping theorem allows us to find a pair of distinct nonzero points $w_{3}$ and $w_{4}$ in $\mathbb{D}$ with $\varphi\left(w_{3}\right)=\varphi\left(w_{4}\right)$. Hence we may assume that $w_{1}$ and $w_{2}$ are both nonzero.) One can easily see that the kernel of $D_{\psi, \varphi, n}$ consists of the set of all polynomials with degree less than $n$. Lemma 2.3 implies that

$$
\begin{aligned}
D_{\psi, \varphi, n}^{*} J\left(\psi\left(w_{2}\right) K_{\overline{w_{1}}}-\psi\left(w_{1}\right) K_{\overline{w_{2}}}\right) & =D_{\psi, \varphi, n}^{*}\left(\overline{\psi\left(w_{2}\right)} K_{w_{1}}-\overline{\psi\left(w_{1}\right)} K_{w_{2}}\right) \\
& =\overline{\psi\left(w_{1}\right) \psi\left(w_{2}\right)} K_{\varphi\left(w_{1}\right)}^{[n]}-\overline{\psi\left(w_{1}\right) \psi\left(w_{2}\right)} K_{\varphi\left(w_{2}\right)}^{[n]}=0 .
\end{aligned}
$$

Since $D_{\psi, \varphi, n}$ is $J$-symmetric, it follows that $\psi\left(w_{2}\right) K_{\overline{W_{1}}}-\psi\left(w_{1}\right) K_{\overline{w_{2}}}$ is a polynomial of degree less than $n$. Therefore

$$
\psi\left(w_{2}\right) \sum_{j=n}^{\infty} \frac{\Gamma(j+2+\alpha)\left(w_{1}\right)^{j} z^{j}}{j!\Gamma(\alpha+2)}-\psi\left(w_{1}\right) \sum_{j=n}^{\infty} \frac{\Gamma(j+2+\alpha)\left(w_{2}\right)^{j} z^{j}}{j!\Gamma(\alpha+2)}=0 .
$$

Thus $\psi\left(w_{2}\right) w_{1}{ }^{m}=\psi\left(w_{1}\right) w_{2}{ }^{m}$ for each $m \geqslant n$. We observe that

$$
\psi\left(w_{1}\right) w_{2}^{n+1}=\psi\left(w_{2}\right) w_{1}^{n+1}=\psi\left(w_{2}\right) w_{1}^{n} w_{1}=\psi\left(w_{1}\right) w_{2}^{n} w_{1},
$$

so $w_{1}=w_{2}$. Consequently $\varphi$ must be univalent.

REmARK 2.5. We can follow the outline of the proof of Proposition 2.4 to see that an analogue of Proposition 2.4 holds for any normal operator $D_{\psi, \varphi, n}$.

If

$$
\varphi(z)=\frac{a z+b}{c z+d}
$$

is a nonconstant linear fractional self-map of $\mathbb{D}$, then the map

$$
\sigma(z)=\frac{\bar{a} z-\bar{c}}{-\bar{b} z+\bar{d}}
$$

also takes $\mathbb{D}$ into itself (see [1, Lemma 1]). Recall that $\|\sigma\|_{\infty}<1$ whenever $\|\varphi\|_{\infty}<1$, in which case both $D_{\varphi, n}$ and $D_{\sigma, n}$ are bounded operators on $A_{\alpha}^{2}$. Cowen [1] determined the adjoint of $C_{\varphi}$ acting on the Hardy space $H^{2}$. Similarly, the second and third authors investigated the adjoints of certain weighted composition-differentiation operators $D_{\psi, \varphi, 1}$ on $H^{2}$ (see [4, Theorem 1]). Our next result shows that an analogue of [4, Theorem 1] holds in the context of the weighted Bergman spaces $A_{\alpha}^{2}$. Recall that $t=(\alpha+2)(\alpha+3) \ldots(\alpha+n+1)$.

PROPOSITION 2.6. For the linear fractional self-maps $\varphi$ and $\sigma$ described above, it follows that

$$
D_{K_{\sigma(0)}^{[n]}, \varphi, n}^{*}=D_{K_{\varphi(0)}^{[n]}, \sigma, n}
$$

Proof. We know that

$$
K_{\varphi(0)}^{[n]}(z)=\frac{t z^{n}}{(1-\overline{(b / d)} z)^{n+\alpha+2}}=\frac{t \overline{d^{n+\alpha+2}} z^{n}}{(\bar{d}-\bar{b} z)^{n+\alpha+2}}
$$

and

$$
K_{\sigma(0)}^{[n]}(z)=\frac{t z^{n}}{(1+(c / d) z)^{n+\alpha+2}}=\frac{t d^{n+\alpha+2} z^{n}}{(c z+d)^{n+\alpha+2}}
$$

We see that

$$
\begin{align*}
D_{K_{\varphi(0)}^{[n]}, \sigma, n}\left(K_{w}\right)(z) & =T_{K_{\varphi(0)}^{[n]}}\left(\frac{t \overline{w^{n}}}{(1-\bar{w} \sigma(z))^{n+\alpha+2}}\right) \\
& =\frac{t^{2} \overline{d^{n+\alpha+2} w^{n}} z^{n}}{(-\bar{b} z+\bar{d}-\overline{w a} z+\overline{w c})^{n+\alpha+2}} \tag{2.8}
\end{align*}
$$

(note that $D^{(n)}\left(K_{w}\right)=\frac{d^{n} K_{w}}{d z^{n}}=\frac{t \overline{w^{n}}}{(1-\bar{w} z)^{n+\alpha+2}}$ ). By Lemma 2.3, we obtain

$$
\begin{align*}
D_{K_{\sigma(0)}^{[n]}, \varphi, n}^{*}\left(K_{w}\right)(z) & =\frac{t \overline{d^{n+\alpha+2} w^{n}}}{(\overline{c w}+\bar{d})^{n+\alpha+2}} K_{\varphi(w)}^{[n]}(z) \\
& =\frac{t^{2} \overline{d^{n+\alpha+2} w^{n}} z^{n}}{(\overline{c w}+\bar{d}-(\overline{a w}+\bar{b}) z)^{n+\alpha+2}} \tag{2.9}
\end{align*}
$$

Since the span of the reproducing kernel functions $K_{w}$ is dense in $A_{\alpha}^{2}$, the result follows from (2.8) and (2.9).

Now we give an example for Proposition 2.6.

Example 1. Suppose that $\varphi(z)=\frac{i}{2} z+\frac{1}{3}$. We can see that $\sigma(z)=\frac{\frac{-i}{2} z}{\frac{1}{3} z+1}, \varphi(0)=$ $1 / 3$, and $\sigma(0)=0$. Then by Proposition 2.6 , we can see that $\frac{n!}{\beta(n)^{2}} D_{z^{n}, \varphi, n}^{*}=D_{K_{\frac{1}{3}}^{[n]}, \sigma, n}$.

Some more examples for Proposition 2.6 will be seen in the proofs of Theorem 2.7 and Propositions 3.1 and 3.2.

Our next theorem completely describes the $J$-symmetric operators $D_{\psi, \varphi, n}$.

THEOREM 2.7. A bounded operator $D_{\psi, \varphi, n}$ is $J$-symmetric on $A_{\alpha}^{2}$ if and only if

$$
\psi(z)=\frac{a}{t n!} K_{\bar{c}}^{[n]}(z)=\frac{a z^{n}}{n!(1-c z)^{n+\alpha+2}}
$$

and

$$
\varphi(z)=c+\frac{b z}{1-c z}
$$

where $a=\psi^{(n)}(0)$ and $b=\varphi^{\prime}(0)$ are both nonzero complex number and $c=\varphi(0)$ belongs to $\mathbb{D}$.

Proof. Suppose that $D_{\underline{\psi, \varphi, n}}$ is $J$-symmetric. By (2.6), (2.7), and Proposition 2.4, we conclude that $J(\psi)=\frac{\overline{\psi^{(n)}(0)}}{t n!} K_{\varphi(0)}^{[n]}$ and so $\psi=\frac{\psi^{(n)}(0)}{t n!} K_{\varphi}^{[n]}=\frac{\psi^{(n)}(0) z^{n}}{n!(1-\varphi(0) z)^{n+\alpha+2}}$, where $\psi^{(n)}(0) \neq 0$. By the general Leibniz rule, we can see that

$$
\psi^{(n+1)}(z)=\frac{\psi^{(n)}(0)}{n!} \sum_{k=0}^{n+1}\binom{n+1}{k}\left(z^{n}\right)^{(n+1-k)}\left(\frac{1}{(1-\varphi(0) z)^{n+\alpha+2}}\right)^{(k)}
$$

Since $\left(z^{n}\right)^{(n+1-k)}(0)=0$, when $0 \leqslant k \leqslant n+1$ and $k \neq 1$, we obtain

$$
\begin{equation*}
\psi^{(n+1)}(0)=(n+1)(n+\alpha+2) \varphi(0) \psi^{(n)}(0) \tag{2.10}
\end{equation*}
$$

Observe that

$$
\begin{align*}
J D_{\psi, \varphi, n}\left(K_{0}^{[n+1]}\right)(z) & =t(n+1)!(\alpha+n+2) J(\psi)(z) J(\varphi)(z) \\
& =\frac{t(n+1)(n+\alpha+2) \overline{\psi^{(n)}(0)} z^{n}}{(1-\overline{\varphi(0)} z)^{n+\alpha+2}} J(\varphi)(z) \tag{2.11}
\end{align*}
$$

By the proof of (2.4), we can see that for any $f \in A_{\alpha}^{2}$, we obtain

$$
\begin{align*}
\left\langle f, D_{\psi, \varphi, n}^{*} K_{0}^{[n+1]}\right\rangle= & \psi^{(n+1)}(0) f^{(n)}(\varphi(0)) \\
& +\sum_{j=1}^{n+1}\binom{n+1}{j} \psi^{(n+1-j)}(0)\left(f^{(n)} \circ \varphi\right)^{(j)}(0) \\
= & \psi^{(n+1)}(0) f^{(n)}(\varphi(0))+(n+1) \psi^{(n)}(0)\left(f^{(n)} \circ \varphi\right)^{(1)}(0) \\
& +\sum_{j=2}^{n+1}\binom{n+1}{j} \psi^{(n+1-j)}(0)\left(f^{(n)} \circ \varphi\right)^{(j)}(0) \\
= & \psi^{(n+1)}(0) f^{(n)}(\varphi(0))+(n+1) \psi^{(n)}(0) \varphi^{\prime}(0) f^{(n+1)}(\varphi(0)) \tag{2.12}
\end{align*}
$$

(note that Proposition 2.4(i) implies that $\psi^{(n+1-j)}(0)=0$ for each $2 \leqslant j \leqslant n+1$ ). Hence by (2.12), we have

$$
\begin{equation*}
D_{\psi, \varphi, n}^{*}\left(K_{0}^{[n+1]}\right)(z)=\overline{\psi^{(n+1)}(0)} K_{\varphi(0)}^{[n]}(z)+(n+1) \overline{\psi^{(n)}(0) \varphi^{\prime}(0)} K_{\varphi(0)}^{[n+1]}(z) \tag{2.13}
\end{equation*}
$$

Therefore by (2.10) and (2.13), we observe that

$$
\begin{align*}
D_{\psi, \varphi, n}^{*} J\left(K_{0}^{[n+1]}\right)(z)= & D_{\psi, \varphi, n}^{*}\left(K_{0}^{[n+1]}\right)(z) \\
= & \overline{\psi^{(n+1)}(0)} K_{\varphi(0)}^{[n]}(z)+(n+1) \overline{\psi^{(n)}(0) \varphi^{\prime}(0)} K_{\varphi(0)}^{[n+1]}(z) \\
= & \frac{t(n+1)(n+\alpha+2) \overline{\varphi(0) \psi^{(n)}(0)} z^{n}}{(1-\overline{\varphi(0)} z)^{n+\alpha+2}} \\
& \quad+\frac{t(n+1)(n+\alpha+2) \overline{\psi^{(n)}(0) \varphi^{\prime}(0)} z^{n+1}}{(1-\overline{\varphi(0)} z)^{n+\alpha+3}} \tag{2.14}
\end{align*}
$$

Because $D_{\psi, \varphi, n}$ is $J$-symmetric, it follows from (2.11) and (2.14) that

$$
J(\varphi)(z)=\overline{\varphi(0)}+\frac{\overline{\varphi^{\prime}(0)} z}{1-\overline{\varphi(0)} z}
$$

and so

$$
\varphi(z)=\varphi(0)+\frac{\varphi^{\prime}(0) z}{1-\varphi(0) z}
$$

with $\varphi^{\prime}(0) \neq 0$ because $\varphi$ is nonconstant.
Conversely, take $\psi$ and $\varphi$ as in the statement of the theorem. For each $f$ in $A_{\alpha}^{2}$, we have

$$
\begin{equation*}
J D_{\psi, \varphi, n}(f)(z)=J(\psi)(z) J\left(f^{(n)}(\varphi(z))\right)=J(\psi)(z) \overline{f^{(n)}(\varphi(\bar{z}))} \tag{2.15}
\end{equation*}
$$

On the other hand, by Proposition 2.6, we see that

$$
D_{\psi, \varphi, n}^{*} J=\frac{\bar{a}}{n!t} D_{K_{\sigma(0)}^{[n]}, \varphi, n}^{*} J=\frac{\bar{a}}{n!t} D_{K_{\varphi(0)}^{[n]}, \sigma, n} J .
$$

Thus

$$
\begin{equation*}
D_{\psi, \varphi, n}^{*} J(f)(z)=\frac{\bar{a}}{n!t} K_{\varphi(0)}^{[n]}(z) \overline{f^{(n)}(\overline{\sigma(z)})}=J(\psi)(z) \overline{f^{(n)}(\varphi(\bar{z}))} \tag{2.16}
\end{equation*}
$$

Therefore, by (2.15) and (2.16), the operator $D_{\psi, \varphi, n}$ is $J$-symmetric.
We infer from the paragraph after Corollary 2.2, from [10, Lemma 4.8], and from the proof of [10, Theorem 4.10] that an operator $D_{\psi, \varphi, n}$ from Theorem 2.7 is bounded on $A_{\alpha}^{2}$ whenever $2\left|c+\bar{c}\left(b-c^{2}\right)\right|<1-\left|b-c^{2}\right|^{2}$.

By an idea similar to one stated in the proof of [3, Proposition 2.1] (see also [13, Theorem 4.1]), we remark that $C_{\psi, \varphi}$ is unitary and $J$-symmetric on $A_{\alpha}^{2}$ if and only if either

$$
\begin{equation*}
\psi(z)=\frac{\alpha\left(1-|p|^{2}\right)^{\frac{\alpha+2}{2}}}{(1-\bar{p} z)^{\alpha+2}} \tag{2.17}
\end{equation*}
$$

and

$$
\begin{equation*}
\varphi(z)=\frac{\bar{p}}{p} \frac{p-z}{1-\bar{p} z} \tag{2.18}
\end{equation*}
$$

where $p$ belongs to $\mathbb{D} \backslash\{0\}$ and $|\alpha|=1$, or $\psi \equiv \mu$ and $\varphi(z)=\lambda z$, where $|\mu|=|\lambda|=$ 1. In the case that $p \neq 0$, we denote the linear functional transformations in (2.17) and (2.18) by $\psi_{p}$ and $\varphi_{p}$ respectively. Invoking [3, Lemma 2.2], we observe that $C_{\lambda z} J$ and $C_{\psi_{p}, \varphi_{p}} J$ are conjugations. Next we will characterize the complex symmetric operators $D_{\psi, \varphi, n}$ with conjugations $C_{\lambda z} J$ and $C_{\psi_{p}, \varphi_{p}} J$.

THEOREM 2.8. Suppose that

$$
\tilde{\varphi}(z)=c+\frac{b z}{1-c z}
$$

and that

$$
\tilde{\psi}(z)=\frac{a z^{n}}{n!(1-c z)^{n+\alpha+2}}
$$

where $a$ and $b$ belong to $\mathbb{C} \backslash\{0\}$ and $c$ belongs to $\mathbb{D}$. Assume that $D_{\tilde{\psi}, \tilde{\varphi}, n}$ is bounded on $A_{\alpha}^{2}$.
(1) For $p \neq 0$, an operator $D_{\psi, \varphi, n}$ on $A_{\alpha}^{2}$ is complex symmetric with conjugation $C_{\psi_{p}, \varphi_{p}} J$ if and only if $\varphi=\tilde{\varphi} \circ \varphi_{p}$ and $\psi=\psi_{p} \cdot\left(\tilde{\psi} \circ \varphi_{p}\right)$ for some $\tilde{\varphi}$ and $\tilde{\psi}$.
(2) For $|\mu|=|\lambda|=1$, an operator $D_{\psi, \varphi, n}$ on $A_{\alpha}^{2}$ is complex symmetric with conjugation $C_{\mu, \lambda z} J$ if and only if $\psi(z)=\mu \tilde{\psi}(\lambda z)$ and $\varphi(z)=\tilde{\varphi}(\lambda z)$ for some $\tilde{\varphi}$ and $\tilde{\psi}$.

Proof. (1) Let $p \neq 0$ and suppose that $D_{\psi, \varphi, n}$ is $C_{\psi_{p}, \varphi_{p}} J$-symmetric. As mentioned in the paragraph preceding the statement of Theorem 2.8, the operator $C_{\psi_{p}, \varphi_{p}}^{*}$ is unitary and $J$-symmetric, so it is not difficult to see that $C_{\psi_{p}, \varphi_{p}}^{*}$ is $C_{\psi_{p}, \varphi_{p}} J$-symmetric. Then [3, Proposition 2.3] implies that $C_{\psi_{p}, \varphi_{p}}^{*} D_{\psi, \varphi, n}$ is $J$-symmetric. It follows from Theorem 2.7 that there is a $J$-symmetric operator $D_{\tilde{\psi}, \tilde{\varphi}, n}$ so that $D_{\psi, \varphi, n}=C_{\psi_{p}, \varphi_{p}} D_{\tilde{\psi}, \tilde{\varphi}, n}$. Hence we observe that $\varphi=\tilde{\varphi} \circ \varphi_{p}$ and $\psi=\psi_{p} \cdot\left(\tilde{\psi} \circ \varphi_{p}\right)$.

Conversely, suppose that $\varphi=\tilde{\varphi} \circ \varphi_{p}$ and $\psi=\psi_{p} \cdot\left(\tilde{\psi} \circ \varphi_{p}\right)$ for some $\tilde{\varphi}$ and $\tilde{\psi}$. Then $D_{\psi, \varphi, n}=C_{\psi_{p}, \varphi_{p}} D_{\tilde{\psi}, \tilde{\varphi}, n}$. Since the weighted composition operator $C_{\psi_{p}, \varphi_{p}}$ is unitary and $J$-symmetric and the operator $D_{\tilde{\psi}, \tilde{\varphi}, n}$ is $J$-symmetric (see Theorem 2.7), the operator $D_{\psi, \varphi, n}$ is $C_{\psi_{p}, \varphi_{p}} J$-symmetric by [3, Proposition 2.3].
(2) The result follows immediately from the technique demonstrated in the proof of part (1).

In the following example, we give some complex symmetric weighted compositiondifferentiation operators.

Example 2. a) Suppose that

$$
\varphi(z)=\frac{1}{4}+\frac{1+2 i z}{4+\frac{i}{2}+(2 i-1) z}
$$

and

$$
\psi(z)=\frac{3^{\alpha+2}\left(z-\frac{i}{2}\right)^{n}}{4^{\alpha+1} n!\left(1+\frac{i}{8}+\left(\frac{i}{2}-\frac{1}{4}\right) z\right)^{n+\alpha+2}}
$$

We can see that $\varphi=\tilde{\varphi} \circ \varphi_{p}$ and $\psi=\psi_{p} \cdot\left(\tilde{\psi} \circ \varphi_{p}\right)$, where $p=i / 2, \varphi_{p}(z)=\frac{z-i / 2}{1+i z / 2}$, $\tilde{\varphi}(z)=\frac{1}{4}+\frac{i z / 2}{1-z / 4}, \psi_{p}(z)=\frac{(3 / 4)^{\alpha+2}}{(1+i z / 2)^{\alpha+2}}$, and $\tilde{\psi}(z)=\frac{4 z^{n}}{n!(1-z / 4)^{n+\alpha+2}}$. It is easy to see that $\|\tilde{\varphi}\|_{\infty}<1$ and so $D_{\tilde{\psi}, \tilde{\varphi}, n}$ is bounded on $A_{\alpha}^{2}$. Invoking Theorem 2.7, the operator $D_{\tilde{\Psi}, \tilde{\varphi}, n}$ is $J$-symmetric and it satisfies the conditions of Proposition 2.4. Theorem 2.8 shows that $D_{\psi, \varphi, n}$ is $C_{\psi_{p}, \varphi_{p}} J$-symmetric.
b) Suppose that $\varphi(z)=\frac{1}{5}+\frac{((i / 10)+1 / 100) i z}{1-i z / 5}$ and $\psi(z)=\frac{e^{i \pi / 4} i^{n} z^{n}}{n!(1-i z / 5)^{n+\alpha+2}}$. We can see that $\varphi(z)=\tilde{\varphi}(\lambda z)$ and $\psi(z)=\mu \tilde{\psi}(\lambda z)$, where $\tilde{\varphi}(z)=\frac{1}{5}+\frac{((i / 10)+1 / 100) z}{1-z / 5}, \tilde{\psi}(z)=$ $\frac{z^{n}}{n!(1-z / 5)^{n+\alpha+2}}, \lambda=i$, and $\mu=e^{i \pi / 4}$. One can see that $D_{\tilde{\Psi}, \tilde{\varphi}, n}$ is bounded and by Theorem 2.7, it is $J$-symmetric. Theorem 2.8 implies that $D_{\psi, \varphi, n}$ is $C_{e^{i \pi / 4}, i z} J$-symmetric.

## 3. Some examples of complex symmetric operators

In this section, we will show that the complex symmetric operators $D_{\psi, \varphi, n}$ we have already identified include all the self-adjoint operators $D_{\psi, \varphi, n}$ and some of the normal operators $D_{\psi, \varphi, n}$. The next proposition provides a characterization of selfadjoint weighted composition-differentiation operators of order $n$ on $A_{\alpha}^{2}$.

Proposition 3.1. A bounded operator $D_{\psi, \varphi, n}$ is self-adjoint on $A_{\alpha}^{2}$ if and only if

$$
\psi(z)=\frac{a z^{n}}{n!(1-\bar{c} z)^{n+\alpha+2}}=\frac{a}{t n!} K_{c}^{[n]}(z)
$$

and

$$
\varphi(z)=c+\frac{b z}{1-\bar{c} z}
$$

where $a=\psi^{(n)}(0)$ and $b=\varphi^{\prime}(0)$ are both nonzero real numbers and $c=\varphi(0)$ belongs to $\mathbb{D}$. Furthermore, for the self-adjoint operator $D_{\psi, \varphi, n}$, one of the following holds:
(i) If $c=0$, then $D_{\psi, \varphi, n}$ is $J$-symmetric.
(ii) If $c \neq 0$, then $D_{\psi, \varphi, n}$ is $C_{e^{-2 i \theta_{z}}} J$-symmetric, where $\theta=\operatorname{Arg}(c)$.

Proof. Suppose that $D_{\psi, \varphi, n}$ is self-adjoint on $A_{\alpha}^{2}$. By (2.4) and Remark 2.5, we have $D_{\psi, \varphi, n}^{*} K_{0}^{[n]}=\overline{\psi^{(n)}(0)} K_{\varphi(0)}^{[n]}$. Moreover, we can see that $D_{\psi, \varphi, n} K_{0}^{[n]}(z)=D_{\psi, \varphi, n}\left(t z^{n}\right)$ $=t n!\psi(z)$. Since $D_{\psi, \varphi, n}$ is self-adjoint, we conclude that

$$
\begin{equation*}
\psi(z)=\frac{\overline{\psi^{(n)}(0)}}{t n!} K_{\varphi(0)}^{[n]}(z)=\frac{\overline{\psi^{(n)}(0)} z^{n}}{n!(1-\overline{\varphi(0)} z)^{n+\alpha+2}} \tag{3.1}
\end{equation*}
$$

Differentiating both sides of (3.1) $n$ times with respect to $z$, we obtain

$$
\begin{align*}
\psi^{(n)}(z) & =\frac{\overline{\psi^{(n)}(0)}}{n!} \sum_{j=0}^{n}\binom{n}{j}\left(z^{n}\right)^{(n-j)}\left(\frac{1}{(1-\overline{\varphi(0)} z)^{n+\alpha+2}}\right)^{(j)} \\
& =\frac{\overline{\psi^{(n)}(0)}}{n!} \sum_{j=0}^{n}\binom{n}{j} \frac{n!}{j!} z^{j}\left(\frac{1}{(1-\overline{\varphi(0)} z)^{n+\alpha+2}}\right)^{(j)} \tag{3.2}
\end{align*}
$$

(note that $\left(z^{n}\right)^{(t)}=\frac{n!}{(n-t)!} z^{n-t}$ for each $t$ with $0 \leqslant t \leqslant n$ ). It follows from (3.2) that $\psi^{(n)}(0)=\overline{\psi^{(n)}(0)}$, and so $\psi^{(n)}(0)$ is real. Moreover, note that $\psi^{(n)}(0) \neq 0$ since $\psi$ is not identically 0 . On the other hand, differentiating both sides of (3.1) $n+1$ times with respect to $z$ yields

$$
\begin{equation*}
\psi^{(n+1)}(0)=(n+1)(n+\alpha+2) \overline{\varphi(0)} \psi^{(n)}(0) \tag{3.3}
\end{equation*}
$$

We can see that

$$
\begin{align*}
D_{\psi, \varphi, n}\left(K_{0}^{[n+1]}\right)(z) & =D_{\psi, \varphi, n}\left(t(n+\alpha+2) z^{n+1}\right) \\
& =\frac{t(n+1)(n+\alpha+2) \psi^{(n)}(0) z^{n}}{(1-\overline{\varphi(0)} z)^{n+\alpha+2}} \varphi(z) \tag{3.4}
\end{align*}
$$

Furthermore, by the idea from (2.4) and the fact that $\psi^{(m)}(0)=0$ for each $m<n$ (see Remark 2.5), we have

$$
\begin{align*}
D_{\psi, \varphi, n}^{*}\left(K_{0}^{[n+1]}\right)(z)= & \overline{\psi^{(n+1)}(0)} K_{\varphi(0)}^{[n]}(z)+(n+1) \overline{\psi^{(n)}(0) \varphi^{\prime}(0)} K_{\varphi(0)}^{[n+1]}(z) \\
= & \frac{t \overline{\psi^{(n+1)}(0)} z^{n}}{(1-\overline{\varphi(0)} z)^{n+\alpha+2}} \\
& \quad+\frac{(n+1) \overline{\psi^{(n)}(0) \varphi^{\prime}(0)} t(n+\alpha+2) z^{n+1}}{(1-\overline{\varphi(0)} z)^{n+\alpha+3}} . \tag{3.5}
\end{align*}
$$

Since $D_{\psi, \varphi, n}$ is self-adjoint, by combining (3.3), (3.4), and (3.5), we see that

$$
\begin{equation*}
\varphi(z)=\varphi(0)+\frac{\overline{\varphi^{\prime}(0)} z}{1-\overline{\varphi(0)} z} \tag{3.6}
\end{equation*}
$$

Differentiating both sides of (3.6) with respect to $z$ and taking $z=0$, we observe that $\varphi^{\prime}(0)$ is real. In addition, because $\varphi$ is not constant, we see that $\varphi^{\prime}(0) \neq 0$.

For the converse, take $\varphi$ and $\psi$ as in the statement of the proposition and suppose that $D_{\psi, \varphi, n}$ is bounded on $A_{\alpha}^{2}$. Proposition 2.6 dictates that

$$
D_{\psi, \varphi, n}^{*}=\frac{\bar{a}}{t n!} D_{K_{\sigma(0)}^{[n]}, \varphi, n}^{*}=\frac{\bar{a}}{t n!} D_{K_{\varphi(0)}^{[n]}, \sigma, n}=D_{\psi, \varphi, n} .
$$

Thus $D_{\psi, \varphi, n}$ is self-adjoint.
We infer from Theorem 2.7 that the operator $D_{\psi, \varphi, n}$ is $J$-symmetric when $c=0$. Now take $c \neq 0$ and set $\tilde{\psi}(z)=\frac{a e^{2 n i \theta} z^{n}}{n!(1-c z)^{n+\alpha+2}}$ and $\tilde{\varphi}(z)=c+\frac{b e^{2 i \theta} z}{1-c z}$. From Theorem 2.7, the operator $D_{\tilde{\psi}, \tilde{\varphi}, n}$ is $J$-symmetric. By [3, Lemma 2.2] and [3, Proposition 2.3], we observe that $C_{e^{-2 i \theta_{z}}} D_{\tilde{\psi}, \tilde{\varphi}, n}$ is $C_{e^{-2 i \theta_{z}}} J$-symmetric. (As stated in the paragraph preceding Theorem 2.8, the composition operator $C_{e^{-2 i \theta_{z}}}$ is unitary and $J$-symmetric.) A direct computation shows that $C_{e^{-2 i \theta_{z}}} D_{\tilde{\psi}, \tilde{\varphi}, n}=D_{\psi, \varphi, n}$, so the result follows.

Now we will characterize those operators $D_{\psi, \varphi, n}$ on $A_{\alpha}^{2}$ that are normal in the case where $\varphi(0)=0$.

Proposition 3.2. Suppose that an operator $D_{\psi, \varphi, n}$ is bounded on $A_{\alpha}^{2}$ and that $\varphi(0)=0$. Then $D_{\psi, \varphi, n}$ is normal if and only if $\psi(z)=a z^{n}$ and $\varphi(z)=b z$, where a belongs to $\mathbb{C} \backslash\{0\}$ and belongs to $\mathbb{D} \backslash\{0\}$. Moreover, in this case $D_{\psi, \varphi, n}$ is $J$-symmetric.

Proof. Assume that $D_{\psi, \varphi, n}$ is normal on $A_{\alpha}^{2}$. We can see that

$$
\begin{equation*}
\left\|D_{\psi, \varphi, n} K_{0}^{[n]}\right\|^{2}=\left\|\left(\frac{n!}{\beta(n)}\right)^{2} \psi\right\|^{2}=\left(\frac{n!}{\beta(n)}\right)^{4} \sum_{j=0}^{\infty}\left(\frac{\beta(j)}{j!}\right)^{2}\left|\psi^{(j)}(0)\right|^{2} \tag{3.7}
\end{equation*}
$$

On the other hand, by (2.4) and Remark 2.5, we observe that

$$
\begin{equation*}
\left\|D_{\psi, \varphi, n}^{*} K_{0}^{[n]}\right\|^{2}=\left\|\overline{\psi^{(n)}(0)} K_{0}^{[n]}\right\|^{2}=\left|\psi^{(n)}(0)\right|^{2}\left(\frac{n!}{\beta(n)}\right)^{2} \tag{3.8}
\end{equation*}
$$

Because $D_{\psi, \varphi, n}$ is normal, by Remark 2.5, (3.7), and (3.8), we conclude that

$$
\begin{equation*}
\left|\psi^{(n)}(0)\right|^{2}\left(\frac{n!}{\beta(n)}\right)^{2}=\left(\frac{n!}{\beta(n)}\right)^{4} \sum_{j=n}^{\infty}\left(\frac{\beta(j)}{j!}\right)^{2}\left|\psi^{(j)}(0)\right|^{2} \tag{3.9}
\end{equation*}
$$

Remark 2.5 implies that $\psi^{(n)}(0) \neq 0$, so (3.9) shows that $\psi^{(j)}(0)=0$ for each $j>n$. Since Remark 2.5 also shows that $\psi^{(j)}(0)=0$ for any $j<n$, the map $\psi$ must have the form $\psi(z)=a z^{n}$ for some $a$ in $\mathbb{C} \backslash\{0\}$. We have

$$
\begin{align*}
D_{\psi, \varphi, n}\left(K_{0}^{[n+1]}\right)(z) & =\left(\frac{(n+1)!}{\beta(n+1)}\right)^{2} \psi(z) \varphi(z)  \tag{3.10}\\
& =\left(\frac{(n+1)!}{\beta(n+1)}\right)^{2} a z^{n} \varphi(z)
\end{align*}
$$

On the other hand, by using (2.4) and the fact that $\psi^{(m)}(0)=0$ for each $m \neq n$, we see that

$$
\begin{align*}
D_{\psi, \varphi, n}^{*}\left(K_{0}^{[n+1]}\right)(z) & =(n+1) \overline{\psi^{(n)}(0) \varphi^{\prime}(0)} K_{0}^{[n+1]}(z) \\
& =\overline{a \varphi^{\prime}(0)}\left(\frac{(n+1)!}{\beta(n+1)}\right)^{2} z^{n+1} \\
& =\overline{a \varphi^{\prime}(0)}(n+1)!K_{0}^{[n+1]}(z) \tag{3.11}
\end{align*}
$$

so $K_{0}^{[n+1]}$ is an eigenvalue for $D_{\psi, \varphi, n}^{*}$ corresponding to eigenvalue $\overline{a \varphi^{\prime}(0)}(n+1)$ !. Therefore

$$
\begin{equation*}
D_{\psi, \varphi, n} K_{0}^{[n+1]}=a \varphi^{\prime}(0)(n+1)!K_{0}^{[n+1]} \tag{3.12}
\end{equation*}
$$

Since $D_{\psi, \varphi, n}$ is normal on $A_{\alpha}^{2}$, by (3.10) and (3.12), we see that

$$
a \varphi^{\prime}(0)(n+1)!K_{0}^{[n+1]}(z)=\left(\frac{(n+1)!}{\beta(n+1)}\right)^{2} a z^{n} \varphi(z)
$$

Thus $\varphi(z)=\varphi^{\prime}(0) z$. Because $\varphi$ is not identically 0 , we conclude that $\varphi(z)=b z$ for some $b$ in $\mathbb{D} \backslash\{0\}$.

For the converse, take $\psi$ and $\varphi$ as in the statement of the proposition and assume that $D_{\psi, \varphi, n}$ is bounded on $A_{\alpha}^{2}$. Proposition 2.6 implies that $D_{a z^{n}, b z, n}^{*}=D_{\bar{a} z^{n}, \bar{b} z, n}$. After some computation, we see that

$$
\begin{align*}
D_{a z^{n}, b z, n} D_{a z^{n}, b z, n}^{*}(f)(z) & =D_{a z^{n}, b z, n} D_{\bar{a} z^{n}, \overline{b z}, n}(f)(z) \\
& =D_{a z^{n}, b z, n}\left(\bar{a} z^{n} f^{(n)}(\bar{b} z)\right) \\
& =|a|^{2} z^{n} \sum_{j=0}^{n}\binom{n}{j} \frac{n!}{j!}|b|^{2 j} z^{j} f^{(n+j)}\left(|b|^{2} z\right) \tag{3.13}
\end{align*}
$$

for each $f$ in $A_{\alpha}^{2}$; similarly,

$$
\begin{equation*}
D_{a z^{n}, b z, n}^{*} D_{a z^{n}, b z, n}(f)(z)=|a|^{2} z^{n} \sum_{j=0}^{n}\binom{n}{j} \frac{n!}{j!}|b|^{2 j} z^{j} f^{(n+j)}\left(|b|^{2} z\right) \tag{3.14}
\end{equation*}
$$

Hence (3.13) and (3.14) show that $D_{\psi, \varphi, n}$ is normal. Furthermore, Theorem 2.7 shows that $D_{\psi, \varphi, n}$ is $J$-symmetric.

Next we describe the conditions under which the analytic functions $\varphi$ and $\psi$ from Proposition 3.1 induce a normal operator $D_{\psi, \varphi, n}$.

Proposition 3.3. Suppose that $D_{\psi, \varphi, n}$ is a bounded operator, with

$$
\psi(z)=\frac{a z^{n}}{n!(1-\bar{c} z)^{n+\alpha+2}}
$$

and

$$
\varphi(z)=c+\frac{b z}{1-\bar{c} z},
$$

where $a=\psi^{(n)}(0)$ and $b=\varphi^{\prime}(0)$ are both nonzero complex numbers and $c=\varphi(0)$ belongs to $\mathbb{D}$. The operator $D_{\psi, \varphi, n}$ is normal on $A_{\alpha}^{2}$ if and only if either $b$ belongs to $\mathbb{R} \backslash\{0\}$ or $c=0$. Moreover, when $D_{\psi, \varphi, n}$ is normal, one of the following holds:
(i) If $c=0$, then $D_{\psi, \varphi, n}$ is $J$-symmetric.
(ii) If $c \neq 0$, then $D_{\psi, \varphi, n}$ is $C_{e^{-2 i \theta_{z}}}$ J-symmetric, where $\theta=\operatorname{Arg}(c)$.

Proof. If $b$ belongs to $\mathbb{R} \backslash\{0\}$ or $c=0$, Propositions 3.1 and 3.2 imply that $D_{\psi, \varphi, n}$ is normal.

For the converse, suppose that $b$ and $c$ belong to $\mathbb{C} \backslash \mathbb{R}$. We have

$$
D_{\psi, \varphi, n}\left(K_{\frac{1}{2}}\right)(z)=\frac{t \psi(z)}{2^{n}\left(1-\frac{1}{2} \varphi(z)\right)^{n+\alpha+2}}=\frac{a}{2^{n} n!(1-c / 2)^{n+\alpha+2}} K_{p_{1}}^{[n]}(z)
$$

where $p_{1}=c+\frac{\bar{b} / 2}{1-\bar{c} / 2}$. On the other hand, by Lemma 2.3, we see that

$$
D_{\psi, \varphi, n}^{*}\left(K_{\frac{1}{2}}\right)(z)=\overline{\psi(1 / 2)} K_{\varphi(1 / 2)}^{[n]}(z)=\frac{\bar{a}}{2^{n} n!(1-c / 2)^{n+\alpha+2}} K_{p_{2}}^{[n]}(z)
$$

where $p_{2}=c+\frac{b / 2}{1-\bar{c} / 2}$.
If $D_{\psi, \varphi, n}$ were normal, then

$$
\begin{aligned}
\left\|D_{\psi, \varphi, n}\left(K_{\frac{1}{2}}\right)\right\|^{2} & =\left|\frac{a}{2^{n} n!(1-c / 2)^{n+\alpha+2}}\right|^{2}\left\|K_{p_{1}}^{[n]}\right\|^{2} \\
& =\left|\frac{a}{2^{n} n!(1-c / 2)^{n+\alpha+2}}\right|^{2} \sum_{j=n}^{\infty} \frac{\left(\left|p_{1}\right|^{2}\right)^{j-n}}{\beta(j)^{2}}\left(\frac{j!}{(j-n)!}\right)^{2}
\end{aligned}
$$

would equal

$$
\begin{aligned}
\left\|D_{\psi, \varphi, n}^{*}\left(K_{\frac{1}{2}}\right)\right\|^{2} & =\left|\frac{a}{2^{n} n!(1-c / 2)^{n+\alpha+2}}\right|^{2}\left\|K_{p_{2}}^{[n]}\right\|^{2} \\
& =\left|\frac{a}{2^{n} n!(1-c / 2)^{n+\alpha+2}}\right|^{2} \sum_{j=n}^{\infty} \frac{\left(\left|p_{2}\right|^{2}\right)^{j-n}}{\beta(j)^{2}}\left(\frac{j!}{(j-n)!}\right)^{2}
\end{aligned}
$$

Therefore $\left|p_{1}\right|^{2}=\left|p_{2}\right|^{2}$. Thus

$$
\left(c+\frac{\bar{b}}{2-\bar{c}}\right)\left(\bar{c}+\frac{b}{2-c}\right)=\left(c+\frac{b}{2-\bar{c}}\right)\left(\bar{c}+\frac{\bar{b}}{2-c}\right)
$$

$$
\begin{aligned}
|c|^{2}+\frac{b c}{2-c}+\frac{\overline{b c}}{2-\bar{c}}+\frac{|b|^{2}}{|2-c|^{2}} & =|c|^{2}+\frac{c \bar{b}}{2-c}+\frac{b \bar{c}}{2-\bar{c}}+\frac{|b|^{2}}{|2-c|^{2}} \\
\frac{c(b-\bar{b})}{2-c} & =\frac{\bar{c}(b-\bar{b})}{2-\bar{c}} \\
\frac{c}{2-c} & =\frac{\bar{c}}{2-\bar{c}} .
\end{aligned}
$$

Then $b=\bar{b}$ or $c=\bar{c}$, which is a contradiction. If $D_{\psi, \varphi, n}$ were normal, with $b$ belonging to $\mathbb{C} \backslash \mathbb{R}$ and $c$ belonging to $\mathbb{R} \backslash\{0\}$, a similar argument would show that $\left\|D_{\psi, \varphi, n}^{*} K_{\frac{i}{2}}\right\| \neq\left\|D_{\psi, \varphi, n} K_{\frac{i}{2}}\right\|$, which is also a contradiction. The rest of the proof is obtained by an argument similar to that of Proposition 3.1.

Now, we give examples for the results of this section.
EXAMPLE 3. Suppose that $\varphi_{1}(z)=\frac{1}{4}+\frac{i z / 2}{1-z / 4}, \varphi_{2}(z)=\frac{e^{\frac{i \pi}{3}}}{4}+\frac{z / 2}{1-\frac{e^{\frac{-i \pi}{3}}}{4}}, \psi_{1}(z)=$ $\frac{3 z^{n}}{n!(1-z / 4)^{n+\alpha+2}}$, and $\psi_{2}(z)=\frac{3 z^{n}}{n!\left(1-\frac{\frac{e-i \pi}{3}}{4} z\right)^{n+\alpha+2}}$. It is easy to see that $\left\|\varphi_{1}\right\|_{\infty}<1$ and $\left\|\varphi_{2}\right\|_{\infty}<1$ and so $D_{\psi_{1}, \varphi_{1}, n}$ and $D_{\psi_{2}, \varphi_{2}, n}$ are bounded on $A_{\alpha}^{2}$. By Theorem 2.7, the operator $D_{\psi_{1}, \varphi_{1}, n}$ is $J$-symmetric. Proposition 3.3 implies that $D_{\psi_{1}, \varphi_{1}, n}$ is not normal. We observe that $D_{\psi_{2}, \varphi_{2}, n}$ is self-adjoint and $C_{e}{ }_{e} \frac{-2 \pi}{3} i_{z} J$-symmetric by Proposition 3.1.

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